CS 283 Advanced Computer Graphics

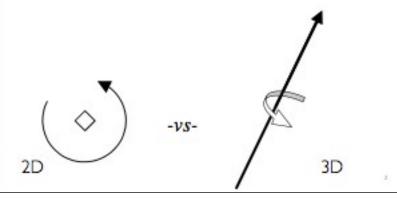
Rotations and Inverse Kinematics

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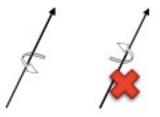
Rotations

- · 3D Rotations fundamentally more complex than in 2D
 - · 2D: amount of rotation
 - . 3D: amount and axis of rotation



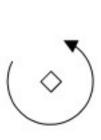
Rotations

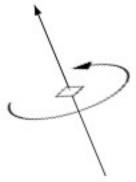
- · Rotations still orthonormal
- $Det(\mathbf{R}) = 1 \neq -1$
- Preserve lengths and distance to origin
- 3D rotations DO NOT COMMUTE!
- Right-hand rule DO NOT COMMUTE!
- Unique matrices



Axis-aligned 3D Rotations

 2D rotations implicitly rotate about a third out of plane axis





Axis-aligned 3D Rotations

 2D rotations implicitly rotate about a third out of plane axis

$$\mathbf{R} = \begin{bmatrix} \cos(\theta) & -\sin(\theta) \\ \sin(\theta) & \cos(\theta) \end{bmatrix}$$

$$\mathbf{R} = \begin{bmatrix} \cos(\theta) & -\sin(\theta) & 0 \\ \sin(\theta) & \cos(\theta) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$



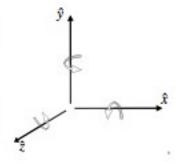


Axis-aligned 3D Rotations

$$\mathbf{R}_{i} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta) & -\sin(\theta) \\ 0 & \sin(\theta) & \cos(\theta) \end{bmatrix}$$

$$\mathbf{R}_{s} = \begin{bmatrix} \cos(\theta) & 0 & \sin(\theta) \\ 0 & 1 & 0 \\ -\sin(\theta) & 0 & \cos(\theta) \end{bmatrix}$$

$$\mathbf{R} = \begin{bmatrix} \cos(\theta) & -\sin(\theta) & 0\\ \sin(\theta) & \cos(\theta) & 0\\ 0 & 0 & 1 \end{bmatrix}$$



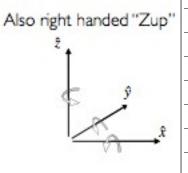
Axis-aligned 3D Rotations

$$\begin{split} \mathbf{R}_{s} &= \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta) & -\sin(\theta) \\ 0 & \sin(\theta) & \cos(\theta) \end{bmatrix} \\ \mathbf{R}_{s} &= \begin{bmatrix} \cos(\theta) & 0 & \sin(\theta) \\ 0 & 1 & 0 \\ -\sin(\theta) & 0 & \cos(\theta) \end{bmatrix} \\ \mathbf{R}_{t} &= \begin{bmatrix} \cos(\theta) & -\sin(\theta) & 0 \\ \sin(\theta) & \cos(\theta) & 0 \\ 0 & 0 & 1 \end{bmatrix} \end{split}$$

Axis-aligned 3D Rotatio

$$\mathbf{R}_{s} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta) & -\sin(\theta) \\ 0 & \sin(\theta) & \cos(\theta) \end{bmatrix}$$
$$\mathbf{R}_{g} = \begin{bmatrix} \cos(\theta) & 0 & \sin(\theta) \\ 0 & 1 & 0 \\ -\sin(\theta) & 0 & \cos(\theta) \end{bmatrix}$$

$$\mathbf{R}_{i} = \begin{bmatrix} \cos(\theta) & -\sin(\theta) & 0\\ \sin(\theta) & \cos(\theta) & 0\\ 0 & 0 & 1 \end{bmatrix}$$



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Axis-aligned 3D Rotations

. Also known as "direction-cosine" matrices

$$\mathbf{R}_{s} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta) & -\sin(\theta) \\ 0 & \sin(\theta) & \cos(\theta) \end{bmatrix} \qquad \mathbf{R}_{g} = \begin{bmatrix} \cos(\theta) & 0 & \sin(\theta) \\ 0 & 1 & 0 \\ -\sin(\theta) & 0 & \cos(\theta) \end{bmatrix}$$

$$\mathbf{R}_{i} = \begin{bmatrix} \cos(\theta) & -\sin(\theta) & 0 \\ \sin(\theta) & \cos(\theta) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Arbitrary Rotations

· Can be built from axis-aligned matrices:

$$\mathbf{R} = \mathbf{R}_{\hat{z}} \cdot \mathbf{R}_{\hat{y}} \cdot \mathbf{R}_{\hat{x}}$$

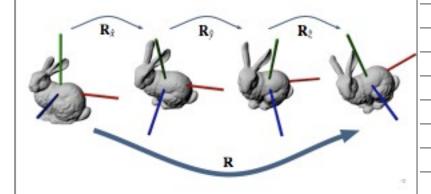
- Result due to Euler... hence called Euler Angles
- · Easy to store in vector
- · But NOT a vector.

$$\mathbf{R} = \operatorname{rot}(x, y, z)$$



Arbitrary Rotations

$$\mathbf{R} = \mathbf{R}_{\hat{z}} \cdot \mathbf{R}_{\hat{y}} \cdot \mathbf{R}_{\hat{x}}$$

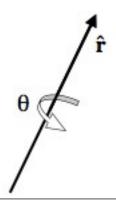


Arbitrary Rotations

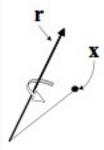
- Allows tumbling
- Euler angles are non-unique
- Gimbal-lock
- Moving -vs- fixed axes
 - · Reverse of each other

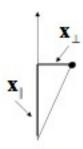
- Direct representation of arbitrary rotation
- AKA: axis-angle, angular displacement vector
- Rotate θ degrees about some axis
- Encode θ by length of vector

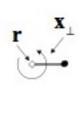
$$\theta = |\mathbf{r}|$$



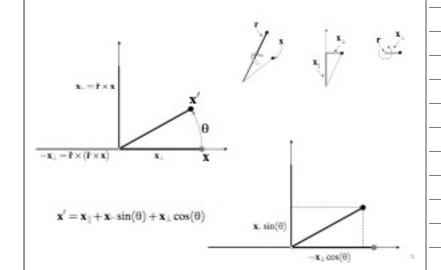
- ullet Given vector ${f r}$, how to get matrix ${f R}$
- Method from text:
 - I. rotate about x axis to put r into the x-y plane
 - 2. rotate about z axis align r with the x axis
 - 3. rotate θ degrees about x axis
 - 4. undo #2 and then #1
 - 5. composite together





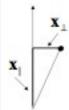


- Vector expressing a point has two parts
 - . X does not change
 - X_rotates like a 2D point



Rodriguez Formula

$$\mathbf{x}' = \mathbf{\hat{r}}(\mathbf{\hat{r}} \cdot \mathbf{x}) \\ + \sin(\theta)(\mathbf{\hat{r}} \times \mathbf{x}) \\ - \cos(\theta)(\mathbf{\hat{r}} \times (\mathbf{\hat{r}} \times \mathbf{x}))$$



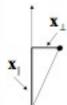


Actually a minor variation

Exponential Maps

Rodriguez Formula

$$\mathbf{x}' = \mathbf{\hat{r}}(\mathbf{\hat{r}} \cdot \mathbf{x}) \\ + \sin(\theta)(\mathbf{\hat{r}} \times \mathbf{x}) \\ - \cos(\theta)(\mathbf{\hat{r}} \times (\mathbf{\hat{r}} \times \mathbf{x}))$$



Linear in x

Actually a minor variation ... ,

· Building the matrix

$$\mathbf{x}' = ((\mathbf{\hat{r}}\mathbf{\hat{r}}^t) + \sin(\theta)(\mathbf{\hat{r}}\times) - \cos(\theta)(\mathbf{\hat{r}}\times)(\mathbf{\hat{r}}\times))\,\mathbf{x}$$

$$\begin{pmatrix} \hat{\mathbf{r}} \times \end{pmatrix} = \begin{bmatrix} 0 & -\hat{r}_z & \hat{r}_y \\ \hat{r}_z & 0 & -\hat{r}_x \\ -\hat{r}_y & \hat{r}_x & 0 \end{bmatrix}$$

Antisymmetric matrix $(\mathbf{a} \times) \mathbf{b} = \mathbf{a} \times \mathbf{b}$ Easy to verify by expansion

- Allows tumbling
- · No gimbal-lockl
- Orientations are space within π-radius ball
- · Nearly unique representation
- Singularities on shells at 2π
- Nice for interpolation

- Why exponential?
- \cdot Recall series expansion of e^x

$$e^x = 1 + \frac{x}{1!} + \frac{x^2}{2!} + \frac{x^3}{3!} + \cdots$$

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- Why exponential?
- Recall series expansion of e^x
- Euler: what happens if you put in $i\theta$ for x

$$e^{i\theta} = 1 + \frac{i\theta}{1!} + \frac{-\theta^2}{2!} + \frac{-i\theta^3}{3!} + \frac{\theta^4}{4!} + \cdots$$

$$= \left(1 + \frac{-\theta^2}{2!} + \frac{\theta^4}{4!} + \cdots\right) + i\left(\frac{\theta}{1!} + \frac{-\theta^3}{3!} + \cdots\right)$$

$$= \cos(\theta) + i\sin(\theta)$$

· Why exponential?

$$e^{(\hat{\mathbf{r}}\times)\theta} = \mathbf{I} + \frac{(\hat{\mathbf{r}}\times)\theta}{1!} + \frac{(\hat{\mathbf{r}}\times)^2\theta^2}{2!} + \frac{(\hat{\mathbf{r}}\times)^3\theta^3}{3!} + \frac{(\hat{\mathbf{r}}\times)^4\theta^4}{4!} + \cdots$$

But notice that: $(\hat{\mathbf{r}} \times)^3 = -(\hat{\mathbf{r}} \times)$

$$e^{(\hat{\mathbf{r}}\times)\theta} = \mathbf{I} + \frac{(\hat{\mathbf{r}}\times)\theta}{1!} + \frac{(\hat{\mathbf{r}}\times)^2\theta^2}{2!} + \frac{-(\hat{\mathbf{r}}\times)\theta^3}{3!} + \frac{-(\hat{\mathbf{r}}\times)^2\theta^4}{4!} + \cdots$$

$$e^{(\hat{\mathbf{r}}\times)\theta} = \mathbf{I} + \frac{(\hat{\mathbf{r}}\times)\theta}{1!} + \frac{(\hat{\mathbf{r}}\times)^2\theta^2}{2!} + \frac{-(\hat{\mathbf{r}}\times)\theta^3}{3!} + \frac{-(\hat{\mathbf{r}}\times)^2\theta^4}{4!} + \cdots$$

$$e^{(\hat{\mathbf{r}}\times)\theta} = (\hat{\mathbf{r}}\times)\left(\frac{\theta}{1!} - \frac{\theta^3}{3!} + \cdots\right) + \mathbf{I} + (\hat{\mathbf{r}}\times)^2\left(+\frac{\theta^2}{2!} - \frac{\theta^4}{4!} + \cdots\right)$$

$$e^{(\hat{\mathbf{r}}\times)\theta} = (\hat{\mathbf{r}}\times)\sin(\theta) + \mathbf{I} + (\hat{\mathbf{r}}\times)^2(1-\cos(\theta))$$

Quaternions

- More popular than exponential maps
- Natural extension of $e^{i\theta} = \cos(\theta) + i\sin(\theta)$
- Due to Hamilton (1843)
 - · Interesting history
 - · Involves "hermaphroditic monsters"

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Quaternions

Uber-Complex Numbers

$$q = (z_1, z_2, z_3, s) = (\mathbf{z}, s)$$

 $q = iz_1 + jz_2 + kz_3 + s$

$$i^{2} = j^{2} = k^{2} = -1$$
 $ij = k \quad ji = -k$
 $jk = i \quad kj = -i$
 $ki = j \quad ik = -j$

Quaternions

· Multiplication natural consequence of defn.

$$\mathbf{q} \cdot \mathbf{p} = (\mathbf{z}_q s_p + \mathbf{z}_p s_q + \mathbf{z}_p \times \mathbf{z}_q \ , \ s_p s_q - \mathbf{z}_p \cdot \mathbf{z}_q)$$

Conjugate

$$q^* = (-\mathbf{z}, s)$$

Magnitude

$$||\mathbf{q}||^2 = \mathbf{z} \cdot \mathbf{z} + s^2 = \mathbf{q} \cdot \mathbf{q}^*$$

2

Quaternions

Vectors as quaternions

$$v = (v, 0)$$

Rotations as quaternions

$$r = (\hat{\mathbf{r}}\sin\frac{\theta}{2}, \cos\frac{\theta}{2})$$

· Rotating a vector

$$x' = r \cdot x \cdot r^*$$

Composing rotations

$$r = r_1 \cdot r_2$$

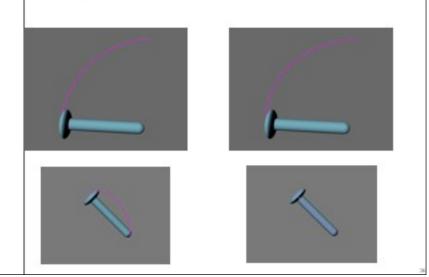
Compare to Exp. Map

Quaternions

- No tumbling
- No gimbal-lock
- Orientations are "double unique"
- Surface of a 3-sphere in 4D ||r|| = 1
- Nice for interpolation

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Interpolation



Rotation Matrices

- Eigen system
 - · One real eigenvalue
 - . Real axis is axis of rotation
 - · Imaginary values are 2D rotation as complex number
- · Logarithmic formula

$$(\hat{\mathbf{r}} \times) = \ln(\mathbf{R}) = \frac{\theta}{2\sin\theta} (\mathbf{R} - \mathbf{R}^{\mathsf{T}})$$

 $\theta = \cos^{-1} \left(\frac{\operatorname{Tr}(\mathbf{R}) - 1}{2} \right)$

Similar formulae as for exponential... ,

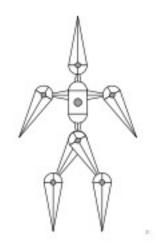
Rotation Matrices

Consider:

$$\mathbf{RI} = \begin{bmatrix} r_{xx} & r_{xy} & r_{xz} \\ r_{yx} & r_{yy} & r_{yz} \\ r_{zx} & r_{zy} & r_{zz} \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

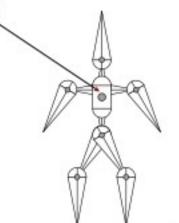
- Columns are coordinate axes after transformation (true for general matrices)
- Rows are original axes in original system (not true for general matrices)

- Articulated skeleton
 - · Topology (what's connected to what)
 - · Geometric relations from joints
 - · Independent of display geometry
 - Tree structure
 - · Loop joints break "tree-ness"

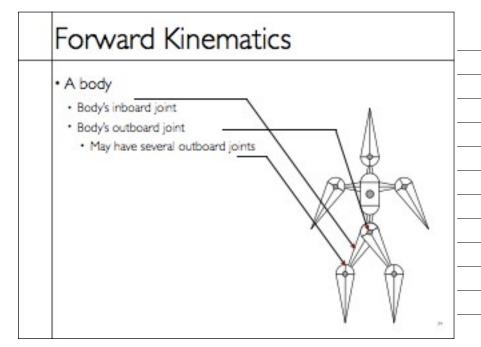


Forward Kinematics

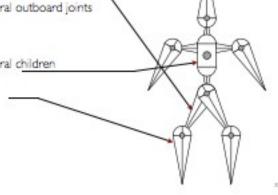
- Root body
 - · Position set by "global" transformation
 - · Root joint
 - Position
 - Rotation
 - · Other bodies relative to root
 - · Inboard toward the root
 - · Outboard away from root



Forward Kinematics • A joint • Joint's inboard body • Joint's outboard body

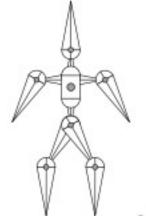


- · A body
 - · Body's inboard joint
 - . Body's outboard joint
 - · May have several outboard joints
 - · Body's parent
 - · Body's child
 - · May have several children

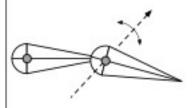


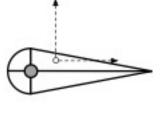
Forward Kinematics

- Interior joints
 - · Typically not 6 DOF joints
 - . Pin rotate about one axis
 - · Ball arbitrary rotation
 - · Prism translation along one axis



- · Pin Joints
 - · Translate inboard joint to local origin
 - . Apply rotation about axis
 - . Translate origin to location of joint on outboard body

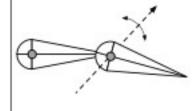


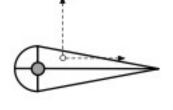


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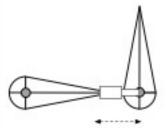
Forward Kinematics

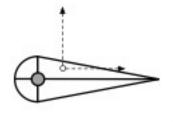
- Ball Joints
 - · Translate inboard joint to local origin
- · Apply rotation about arbitrary axis
- . Translate origin to location of joint on outboard body





- Prismatic Joints
 - · Translate inboard joint to local origin
 - . Translate along axis
 - . Translate origin to location of joint on outboard body

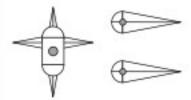




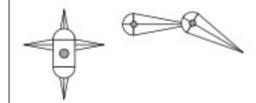
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Forward Kinematics

· Composite transformations up the hierarchy



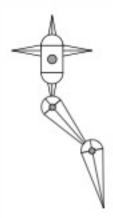
Composite transformations up the hierarchy



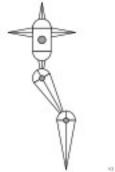
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Forward Kinematics

Composite transformations up the hierarchy

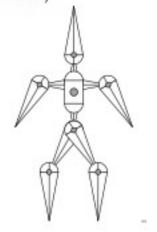


Composite transformations up the hierarchy

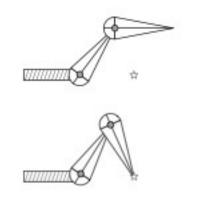


Forward Kinematics

Composite transformations up the hierarchy



- Given
 - · Root transformation
 - · Initial configuration
 - · Desired end point location
- Find
- · Interior parameter settings



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Inverse Kinematics



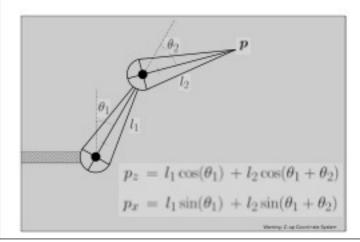






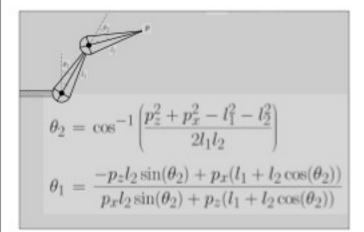
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A simple two segment arm in 2D

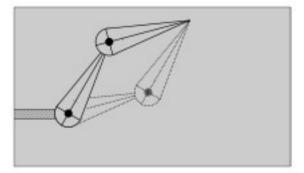


Inverse Kinematics

· Direct IK: solve for the parameters



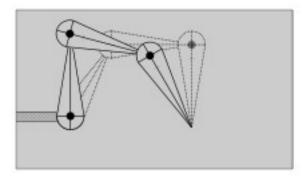
- Why is the problem hard?
 - · Multiple solutions separated in configuration space



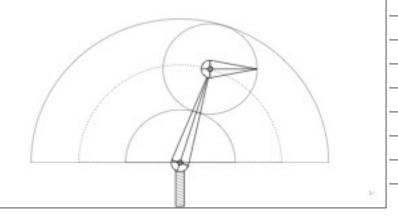
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Inverse Kinematics

- · Why is the problem hard?
 - · Multiple solutions connected in configuration space



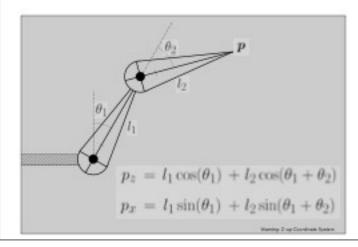
- Why is the problem hard?
 - · Solutions may not always exist



Inverse Kinematics

- Numerical Solution
 - · Start in some initial configuration
 - . Define an error metric (e.g. goal pos current pos)
 - . Compute Jacobian of error w.r.t. inputs
 - · Apply Newton's method (or other procedure)
 - · Iterate...

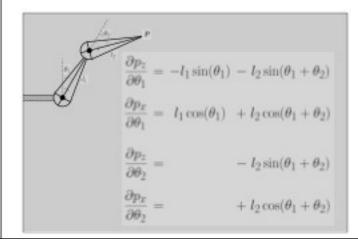
• Recall simple two segment arm:

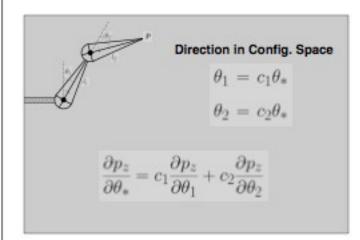


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Inverse Kinematics

. We can write of the derivatives





Inverse Kinematics

The Jacobian (of p w.r.t. θ)

$$J_{ij} = \frac{\partial p_i}{\partial \theta_j}$$

Example for two segment arm

$$J = \begin{bmatrix} \frac{\partial p_z}{\partial \theta_1} & \frac{\partial p_z}{\partial \theta_2} \\ \frac{\partial p_x}{\partial \theta_1} & \frac{\partial p_x}{\partial \theta_2} \end{bmatrix}$$

The Jacobian (of p w.r.t. θ)

$$J = \begin{bmatrix} \frac{\partial p_z}{\partial \theta_1} & \frac{\partial p_z}{\partial \theta_2} \\ \frac{\partial p_x}{\partial \theta_1} & \frac{\partial p_x}{\partial \theta_2} \end{bmatrix}$$

$$\frac{\partial \boldsymbol{p}}{\partial \theta_*} = J \cdot \begin{bmatrix} \frac{\partial \theta_1}{\partial \theta_*} \\ \frac{\partial \theta_2}{\partial \theta_*} \end{bmatrix} = J \cdot \begin{bmatrix} c_1 \\ c_2 \end{bmatrix}$$

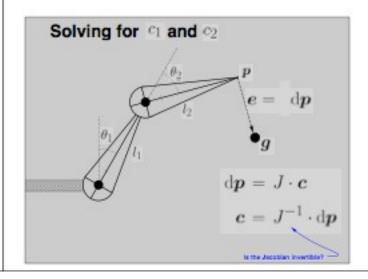
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Inverse Kinematics

Solving for cand co

$$\boldsymbol{c} = \begin{bmatrix} c_1 \\ c_2 \end{bmatrix}$$
 $d\boldsymbol{p} = \begin{bmatrix} dp_z \\ dp_x \end{bmatrix}$

$$d\mathbf{p} = J \cdot \mathbf{c}$$
$$\mathbf{c} = J^{-1} \cdot d\mathbf{p}$$

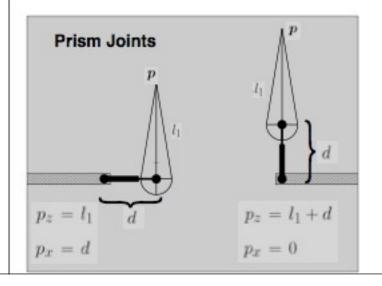


Inverse Kinematics

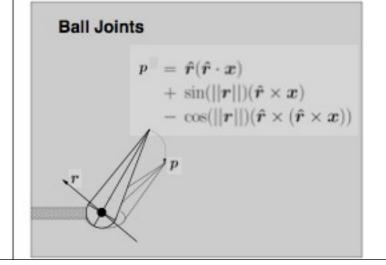
- Problems
 - · Jacobian may (will) not always be invertible
 - · Use pseudo inverse (SVD)
 - · Robust iterative method
 - · Jacobian is not constant

• Nonlinear
$$J = \begin{bmatrix} \frac{\partial p_z}{\partial \theta_1} & \frac{\partial p_z}{\partial \theta_2} \\ \frac{\partial p_x}{\partial \theta_1} & \frac{\partial p_x}{\partial \theta_2} \end{bmatrix} = J(\theta)$$
 by well behaved

φ,



Inverse Kinematics



Ball Joints (moving axis)

$$\mathrm{d}m{p} = [\mathrm{d}m{r}] \cdot e^{[m{r}]} \cdot m{x} = [\mathrm{d}m{r}] \cdot m{p} = -[m{p}] \cdot \mathrm{d}m{r}$$

$$[r] = \begin{bmatrix} 0 & -r_3 & r_2 \\ r_3 & 0 & -r_1 \\ -r_2 & r_1 & 0 \end{bmatrix}$$

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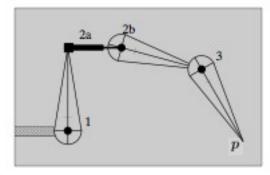
Inverse Kinematics

Ball Joints (fixed axis)

$$\mathrm{d}\boldsymbol{p} = (\mathrm{d}\boldsymbol{\theta})[\boldsymbol{\hat{r}}] \cdot \boldsymbol{x} = -[\boldsymbol{x}] \cdot \boldsymbol{\hat{r}} \mathrm{d}\boldsymbol{\theta}$$

That is the Jacobian for this joint

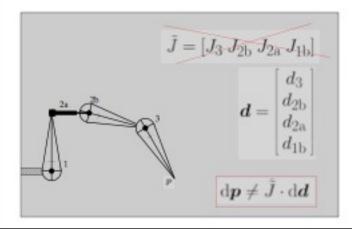
- · Many links / joints
 - · Need a generic method for building Jacobian



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Inverse Kinematics

Can't just concatenate individual matrices



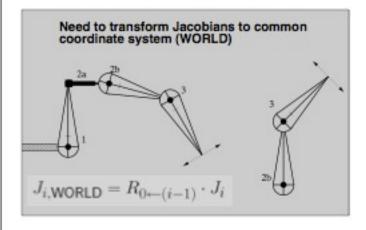
Transformation from body to world

$$X_{0 \leftarrow i} = \prod_{j=1}^{i} X_{(j-1) \leftarrow j} = X_{0 \leftarrow 1} \cdot X_{1 \leftarrow 2} \cdots$$

Rotation from body to world

$$R_{0 \leftarrow i} = \prod_{j=1}^{i} R_{(j-1) \leftarrow j} = R_{0 \leftarrow 1} \cdot R_{1 \leftarrow 2} \cdots$$

Inverse Kinematics



uil.

$$J = \begin{bmatrix} R_{0 \leftarrow 2\mathrm{b}} \cdot J_3(\theta_3, \boldsymbol{p_3}) \\ R_{0 \leftarrow 2\mathrm{a}} \cdot J_{2\mathrm{b}}(\theta_{2\mathrm{b}}, X_{2\mathrm{b} \leftarrow 3} \cdot \boldsymbol{p_3}) \\ R_{0 \leftarrow 1} \cdot J_{2\mathrm{a}}(\theta_{2\mathrm{a}}, X_{2\mathrm{a} \leftarrow 3} \cdot \boldsymbol{p_3}) \\ J_1(\theta_1, X_{1 \leftarrow 3} \cdot \boldsymbol{p_3}) \end{bmatrix}^\mathsf{T}$$

$$\boldsymbol{d} = \begin{bmatrix} d_3 \\ d_{2\mathrm{b}} \\ d_{2\mathrm{a}} \\ d_{1\mathrm{b}} \end{bmatrix}$$

$$| \mathbf{d}\boldsymbol{p} = J \cdot d\boldsymbol{d}$$

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A Cheap Alternative

- · Estimate Jacobian (or parts of it) using finite differences
- Cyclic Coordinate Descent
 - · Solve for each DOF one at a time
- · Iterate till good enough / run out of time

- More complex systems
 - · More complex joints (prism and ball)
 - · More links
 - · Other criteria (COM or height)
 - · Hard constraints (eg foot plants)
 - · Unilateral constraints (eg joint limits)
 - . Multiple criteria and multiple chains
- Smoothness over time
- . DOF are determined by control points of a curve (chain rule)

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Inverse Kinematics

- Some issues
- · How to pick from multiple solutions?
- · Robustness when no solutions
- · Contradictory solutions
- · Smooth interpolation
 - · Interpolation aware of constraints

Grochow, Hartin, Hertzmann, Popović