

## Magnitude Bode Plot for the Low Pass Filter

- Transfer function is

$$\frac{V_{out}}{V_{in}} = \frac{1}{1 + j\omega RC}$$

“The magnitude of the ratio is the ratio of the magnitudes:”

$$\left| \frac{V_{out}}{V_{in}} \right|_{dB} = 20 \log \left| \frac{1}{1 + j\omega RC} \right| = 20 \log \left( \frac{1}{\sqrt{1 + (\omega/\omega_o)^2}} \right)$$

- $\omega_o = 1 / RC$  is the “break frequency” or “-3 dB frequency”

$\omega \ll \omega_o$  results in a magnitude of  $20 \log (1/1) = 0$  dB

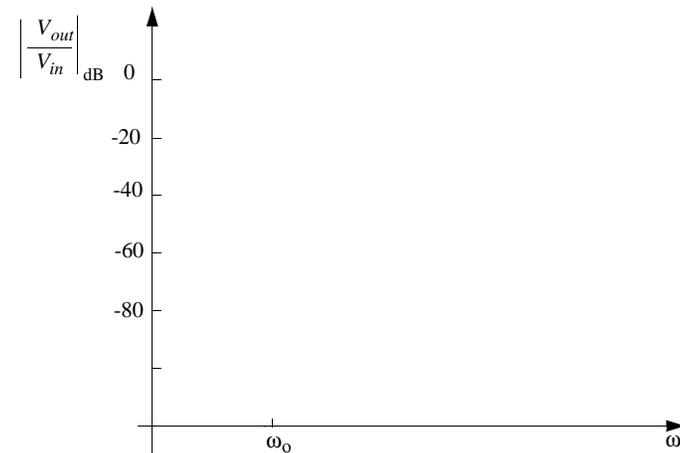
$\omega \gg \omega_o$  results in a magnitude of

$$\left| \frac{V_{out}}{V_{in}} \right|_{dB} = 20 \log \left( \frac{1}{\sqrt{1 + (\omega/\omega_o)^2}} \right) \cong 20 \log \left( \frac{1}{(\omega/\omega_o)} \right) = -20 \log \left( \frac{\omega}{\omega_o} \right)$$

- substitute  $\omega = 10 \omega_o, 100 \omega_o, 1000 \omega_o$

## Approximate Magnitude Bode Plot

- Sketch asymptotes above and below the break frequency



At  $\omega = \omega_o$ , the exact magnitude is:

$$\left| \frac{V_{out}}{V_{in}} \right|_{dB} = 20 \log \left| \frac{1}{1 + j(\omega_o/\omega_o)} \right| = 20 \log \left[ \frac{1}{\sqrt{1 + 1}} \right] = -3 \text{ dB}$$

## Phase Bode Plot for Low Pass Filter

- From the definition of the phase,

$$\angle \frac{V_{out}}{V_{in}} = \angle \left( \frac{1}{1 + j\omega RC} \right) = \angle 1 - \angle(1 + j(\omega/\omega_0)) = -\angle(1 + j(\omega/\omega_0))$$

- Substituting the arctangent,

$$\angle \frac{V_{out}}{V_{in}} = -\text{atan}(\omega/\omega_0)$$

- Look at asymptotes, again:

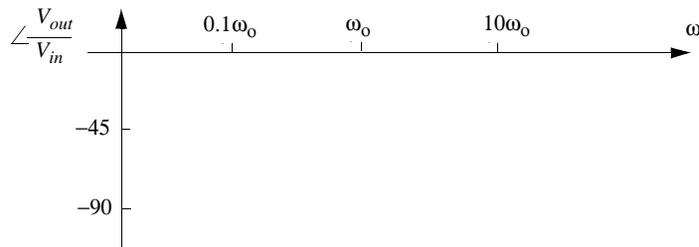
$\omega \ll \omega_0$  results in a phase of  $-\text{atan}(0) = 0$

$\omega \gg \omega_0$  results in a phase of  $-\text{atan}(\text{infinity}) = -90^\circ$

$\omega = (1/10)\omega_0$  results in a phase of  $-\text{atan}(0.1) = -6^\circ$

$\omega = (10)\omega_0$  results in a phase of  $-\text{atan}(10) = -84^\circ$

$\omega = \omega_0$  results in a phase of  $-\text{atan}(1) = -45^\circ$



## Finding the Output Waveform from the Bode Plot

- Suppose that  $v_{in}(t) = 100 \text{ mV} \cos(\omega_0 t + 0^\circ)$

note that the input signal frequency is equal to the break frequency and that the phase is  $0^\circ$  ... the input signal phase is arbitrary and is generally selected to be 0.

the output phasor is:

$$V_{out} = V_{in} \left[ \frac{1}{1 + j(\omega_0/\omega_0)} \right] = V_{in} \left[ \frac{1}{1 + j} \right]$$

magnitude:

$$\left| \frac{V_{out}}{V_{in}} \right|_{dB} = -3 \text{ dB} \quad \dots \quad |V_{out}| = \frac{|V_{in}|}{\sqrt{2}} = \frac{100 \text{ mV}}{\sqrt{2}} = 71 \text{ mV}$$

phase:

$$\angle \frac{V_{out}}{V_{in}} = \angle 1 - \angle(1 + j) = 0 - 45^\circ \quad \angle V_{out} = -45^\circ$$

$$V_{out} = (71 \text{ mV}) e^{-j45^\circ}$$

output waveform  $v_{out}(t)$  is given by:

$$v_{out}(t) = \text{Re} \left( V_{out} e^{j\omega_0 t} \right) = \text{Re} \left( 71 \text{ mV} e^{-j45^\circ} e^{j\omega_0 t} \right)$$

$$v_{out}(t) = 71 \text{ mV} \cos(\omega_0 t - 45^\circ)$$

## Bode Plots of General Transfer Functions

- Procedure is to identify standard forms in the transfer functions, apply asymptotic techniques to sketch each form, and then combine the sketches graphically

$$H(j\omega) = \frac{Aj\omega(1+j\omega\tau_2)(1+j\omega\tau_4)\dots(1+j\omega\tau_n)}{(1+j\omega\tau_1)(1+j\omega\tau_3)\dots(1+j\omega\tau_{n-1})}$$

where the  $\tau_i$  are time constants --  $(1/\tau_i)$  are the break frequencies, which are called *poles* when in the denominator and *zeroes* when in the numerator

- From complex algebra, the factors can be dealt with *separately* in the magnitude and in the phase and the results *added up* to find  $|H(j\omega)|$  and phase  $(H(j\omega))$

Three types of factors:

1. poles (binomial factors in the denominator)
2. zeroes (binomial factors in the numerator)
3.  $j\omega$  in the numerator (or denominator)

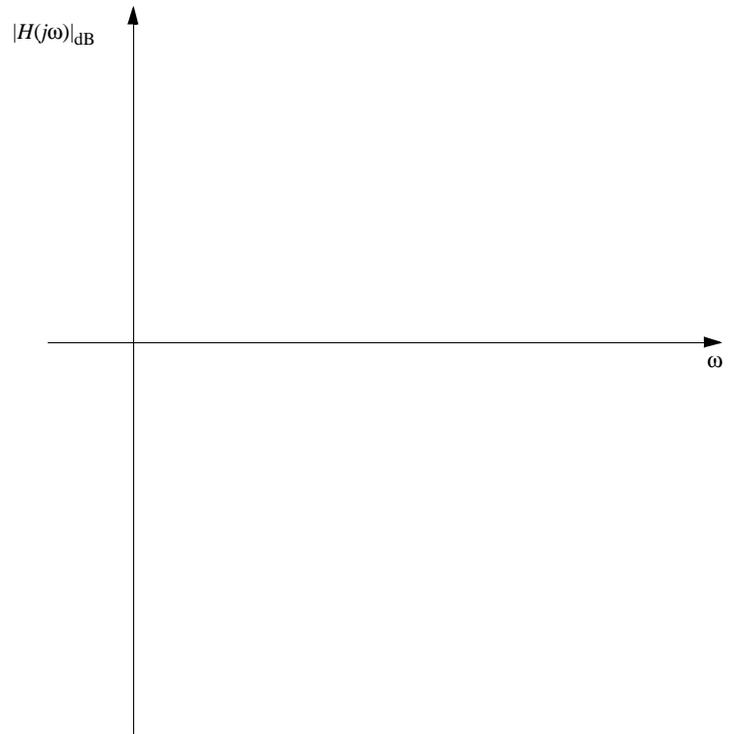
(note: we aren't going to consider complex poles)

## Rapid Sketching of Bode Plots

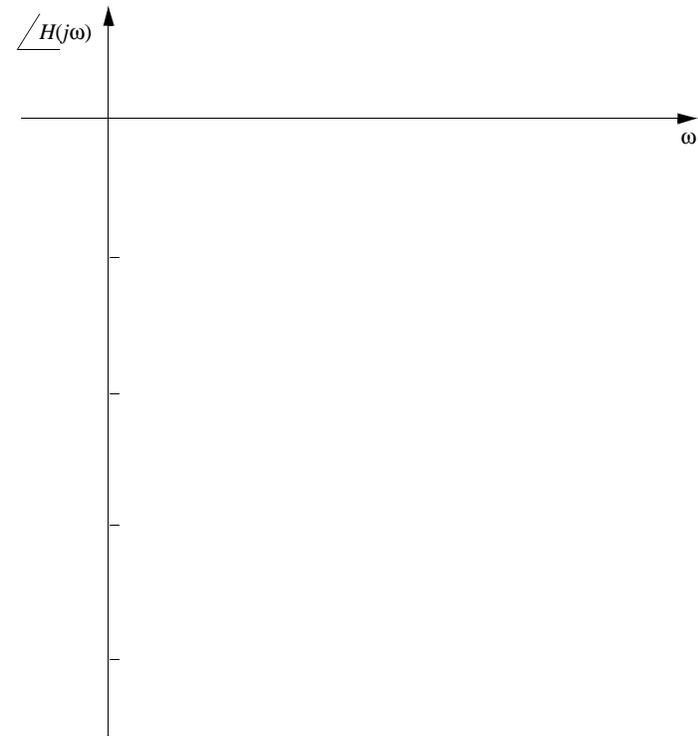
- Poles: -3 dB and  $-45^\circ$  at break frequency  
0 dB below and -20 dB/decade above  
 $0^\circ$  for low frequencies and  $-90^\circ$  for high frequencies; width of transition is between 10 and (1/10) break frequency
- Zeros: +3 dB and  $+45^\circ$  at break frequency  
0 dB below and +20 dB/decade above  
 $0^\circ$  for low frequencies and  $+90^\circ$  for high frequencies; width of transition is between 10 and (1/10) break frequency
- $j\omega$ : +20 dB/decade (0 dB at  $\omega = 1$  rad/s) and  $+90^\circ$  contribution to phase

Example I:

### Magnitude and Phase Bode Plot Sketches



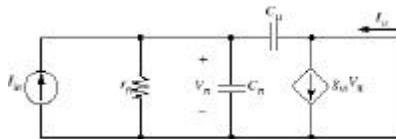
### Phase Bode Plot Sketches



## Device Models and Frequency Response

- “Classic” cases: give insight into the connection between device model and frequency response

1. CE short-circuit current gain  $A_f(j\omega)$  as a function of frequency



Kirchhoff's current law at the output node:

$$I_o = g_m V_\pi - V_\pi j\omega C_\mu$$

Kirchhoff's current law at the input node:

$$I_s = \frac{V_\pi}{Z_\pi} + V_\pi j\omega C_\mu \quad \text{where} \quad Z_\pi = r_\pi \parallel \left( \frac{1}{j\omega C_\pi} \right)$$

- Solving for  $V_\pi$  at the input node:

$$V_\pi = \frac{I_s}{(1/Z_\pi) + j\omega C_\mu}$$

## Short-Circuit Gain Frequency Response

- Substituting  $V_\pi$  into the output node equation--

$$\frac{I_o}{I_s} = \frac{g_m Z_\pi \left( 1 - \frac{j\omega C_\mu}{g_m} \right)}{1 + j\omega C_\mu Z_\pi}$$

- Substituting for  $Z_\pi$  and simplifying --

$$\frac{I_o}{I_s} = \frac{g_m r_\pi \left( 1 - \frac{j\omega C_\mu}{g_m} \right)}{1 + j\omega r_\pi (C_\pi + C_\mu)} = \frac{\beta_o \left( 1 - \frac{j\omega C_\mu}{g_m} \right)}{1 + j\omega r_\pi (C_\pi + C_\mu)} = \beta_o \left[ \frac{1 - j\frac{\omega}{\omega_z}}{1 + j\frac{\omega}{\omega_p}} \right]$$

Current gain has one pole:

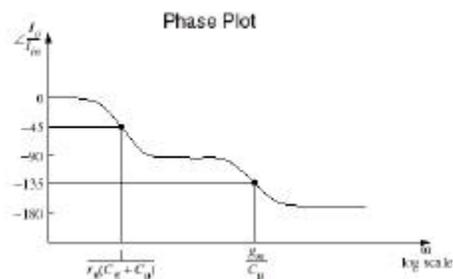
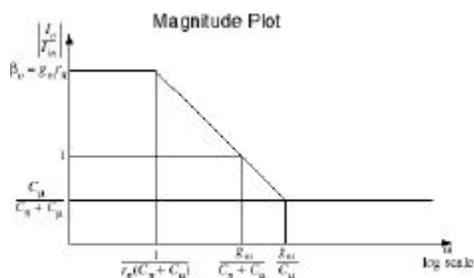
$$\omega_p = (r_\pi (C_\pi + C_\mu))^{-1}$$

and one zero

$$\omega_z = (g_m^{-1} C_\mu)^{-1} \gg \omega_p$$

## Bode Plot of Short-Circuit Current Gain

- Note low frequency magnitude of gain is  $\beta_o$



- Frequency at which current gain is reduced to 0 dB is defined as the **transition frequency**  $\omega_T$ . Neglecting the zero,

$$\omega_T = \frac{g_m}{(C_\pi + C_\mu)}$$

## Transition Frequency of the Bipolar Transistor

- Dependence of transition time  $\tau_T = \omega_T^{-1}$  on the bias collector current  $I_C$ :

$$\tau_T = \frac{1}{\omega_T} = \frac{C_\pi + C_\mu}{g_m} = \frac{g_m \tau_F + C_{jE} + C_\mu}{g_m}$$

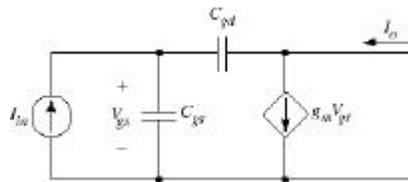
$$\tau_T = \tau_F + \left( \frac{C_{jE} + C_\mu}{g_m} \right) = \tau_F + \frac{V_{th}}{I_C} (C_{jE} + C_\mu)$$

- If the collector current is increased enough to make the second term negligible, then the minimum  $\tau_T$  is the base transit time,  $\tau_F$ . In practice, the  $\omega_T$  decreases at very high values of  $I_C$  due to other effects and the minimum  $\tau_T$  may not be achieved.
- Numerical values of  $f_T = (1/2\pi)\omega_T$  range from 10 MHz for lateral pnp's to 75 GHz for submicron, oxide-isolated, SiGe heterojunction npn's

Note that the small-signal model is not valid above  $f_T$  (due to distributed effects in the base) and the zero in the current gain is not actually observed

## Common-Source Current Gain

- CS amplifier has a non-infinite input *impedance* for  $\omega > 0$  and we can measure its small-signal current gain.



- Analysis is similar to CE case; result is

$$\frac{I_o}{I_{in}} = \frac{g_m \left( 1 - \frac{j\omega C_{gd}}{g_m} \right)}{j\omega(C_{gs} + C_{gd})} \approx \frac{g_m}{j\omega(C_{gs} + C_{gd})}$$

- Transition frequency for the MOSFET is

$$\omega_T \approx \frac{g_m}{C_{gs} + C_{gd}}$$

## Transition Frequency of the MOSFET

- Substitution of gate-source capacitance and transconductance:

$$C_{gs} = \frac{2}{3}WLC_{ox} \gg C_{gd} \text{ and } g_m = \frac{W}{L}\mu_n C_{ox}(V_{GS} - V_{Tn})$$

$$\omega_T \approx \frac{g_m}{C_{gs}} = \frac{\frac{W}{L}\mu_n C_{ox}(V_{GS} - V_{Tn})}{\frac{2}{3}WLC_{ox}} = \frac{3}{2}\mu_n \left[ \frac{(V_{GS} - V_{Tn})}{L} \right] \left( \frac{1}{L} \right)$$

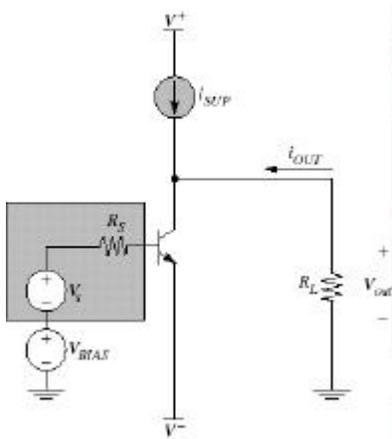
- The transition time is the inverse of  $\omega_T$  and can be written as the average time for electrons to drift from source to drain

$$\tau_T = \frac{L}{\mu_n \left[ \frac{3(V_{GS} - V_{Tn})}{L} \right]} = \left| \frac{L}{v_{dr}} \right|$$

- velocity saturation causes  $\tau_T$  to decrease linearly with  $L$ ; however, submicron MOSFETs have transition frequencies that are approaching those for oxide-isolated BJTs

## Frequency Response of Voltage Amplifiers

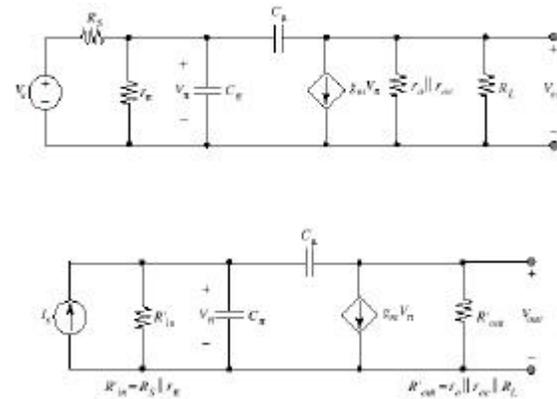
- Common-emitter amplifier:



Procedure: substitute small-signal model at the operating point and perform phasor analysis

## “Brute Force” Phasor Analysis

- “Exact” analysis: transform into Norton form at input to facilitate nodal analysis



Note that  $C_{cs}$  is omitted, along with  $r_b$