

Lecture 26: High Frequency Circuit Analysis

- Announcements:
- HW#8 online and due Friday via Gradescope
- Lab#5 due Tuesday, Oct. 30, 5 p.m.
- I am on travel (as indicated in the schedule shown on the first day)

↳ This is a video lecture

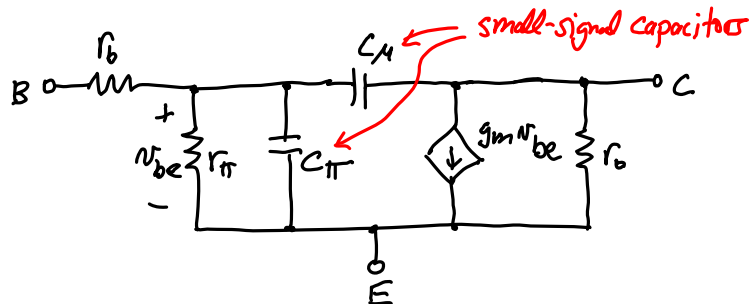
 • **Lecture Topics:**

- ↳ Measuring High Frequency Model for BJT
- ↳ MOS High Frequency Model
- ↳ Brute Force CE HF Analysis
- ↳ Open Circuit Time Constant (OCTC) Analysis

 • **Last Time:**

- Generated the high frequency small-signal model for a BJT
- Examining methods to extract the model experimentally
- Now, continue with this ...

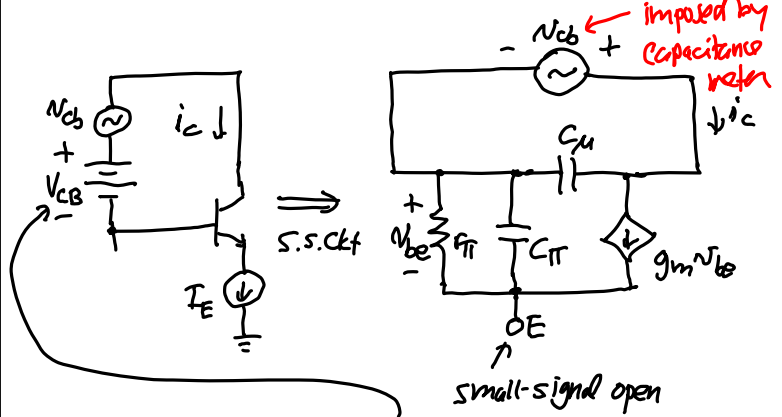
High Frequency Hybrid- π Model



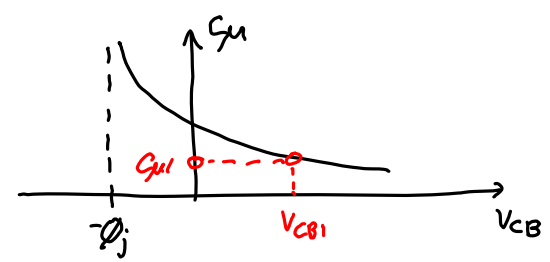
$$C_{\pi} = C_b + C_{je} = I_E g_m + \frac{C_{je0}}{\sqrt{1 + \frac{V_{EB}}{\phi_{je}}}}$$

Determining C_{π} and C_{μ}

- Can experimentally determine C_{μ} by measuring the small-signal capacitance between the base and collector terminals with the emitter incrementally open, i.e., use a current source to set I_E

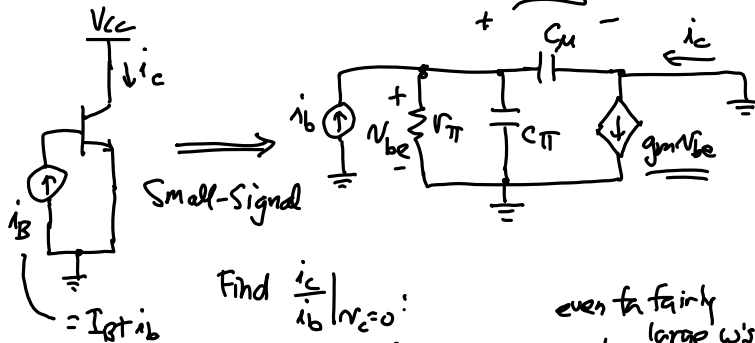


⇒ can measure C_{μ} vs. V_{CB}



- To find C_{π} , find an expression for C_{π} in terms of C_{μ} and known measurable parameters
- One parameter we can conveniently measure is the short-circuit current gain:

$$h_{fe} = \left. \frac{i_c}{i_b} \right|_{N_c=0}$$



Find $\frac{i_c}{i_b} \Big|_{N_c=0}$:

$$v_{be} = i_b \left(r_{\pi} \parallel \frac{1}{sC_{\pi}} \parallel \frac{1}{sC_{\mu}} \right)$$

$$i_c = g_m v_{be} - sC_{\mu} v_{be} = (g_m - sC_{\mu}) v_{be} \approx g_m v_{be}$$

$$= g_m \left(r_{\pi} \parallel \frac{1}{sC_{\pi}} \parallel \frac{1}{sC_{\mu}} \right) i_b$$

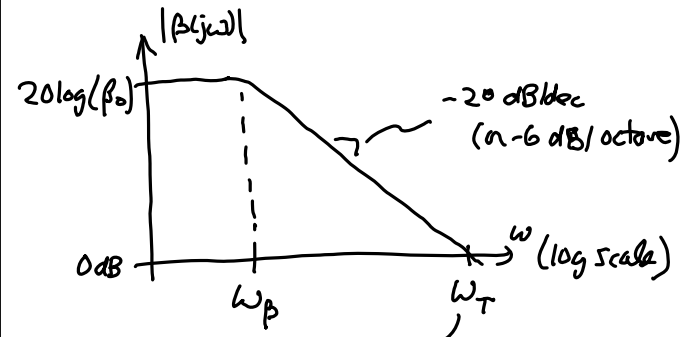
$$\frac{i_c}{i_b} = \frac{g_m}{\frac{1}{r_{\pi}} + s(C_{\pi} + C_{\mu})} = \frac{g_m r_{\pi}}{1 + s r_{\pi} (C_{\pi} + C_{\mu})} = \frac{\beta_0}{1 + s r_{\pi} (C_{\pi} + C_{\mu})}$$

[$\beta_0 = g_m r_{\pi}$]
[low freq β]

$$\beta(j\omega) = \frac{\beta_0}{1 + \frac{j\omega}{\omega_{\beta}}}$$

$$\omega_{\beta} = \frac{1}{r_{\pi} (C_{\pi} + C_{\mu})}$$

Plot $|\beta(j\omega)|$: (Bode plot)



ω_T is defined as the radian frequency where $|\beta(j\omega)| = 1$

$$|\beta(j\omega_T)| \equiv \frac{\beta_0}{\omega_T r_{\pi} (C_{\pi} + C_{\mu})} = 1$$

$$\omega_T = \frac{g_m}{C_{\pi} + C_{\mu}} \Rightarrow f_T = \frac{\omega_T}{2\pi}$$

Also, note:

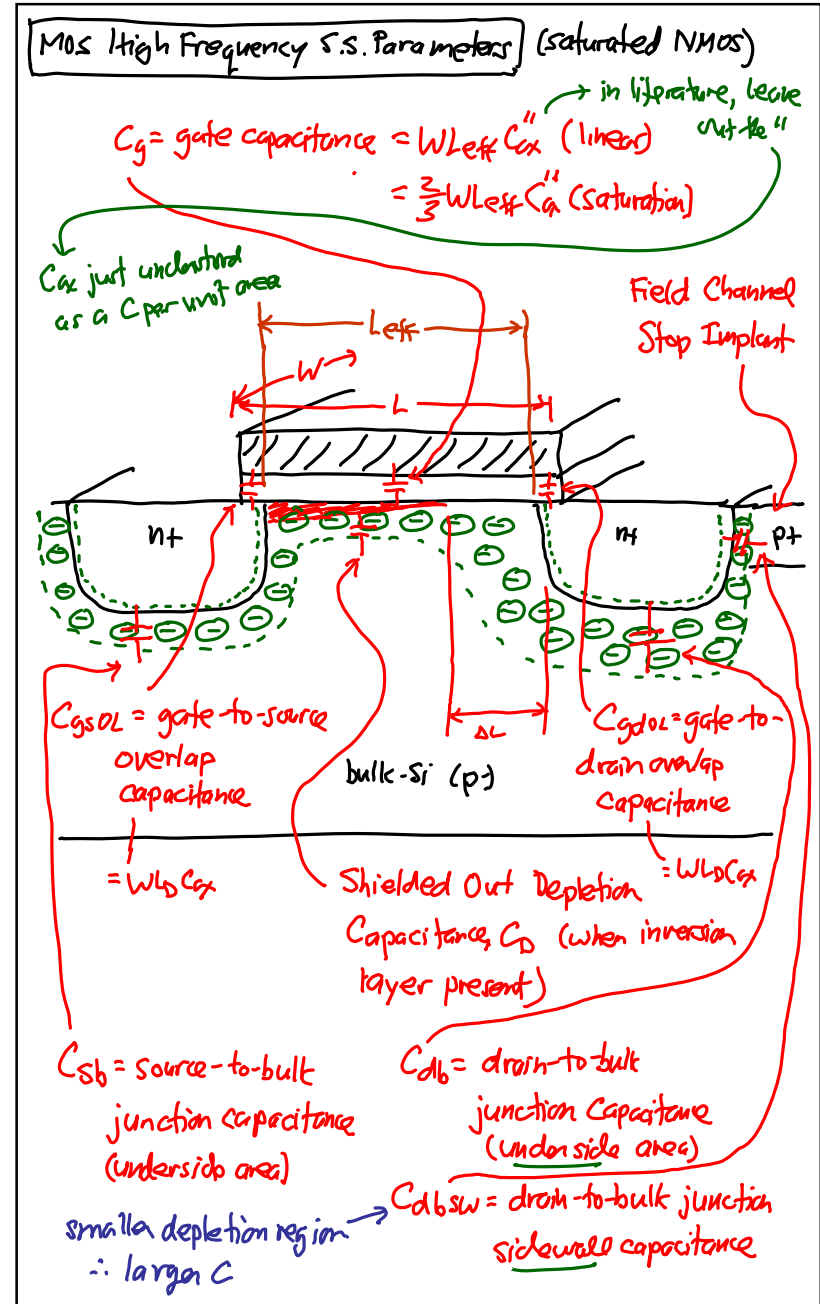
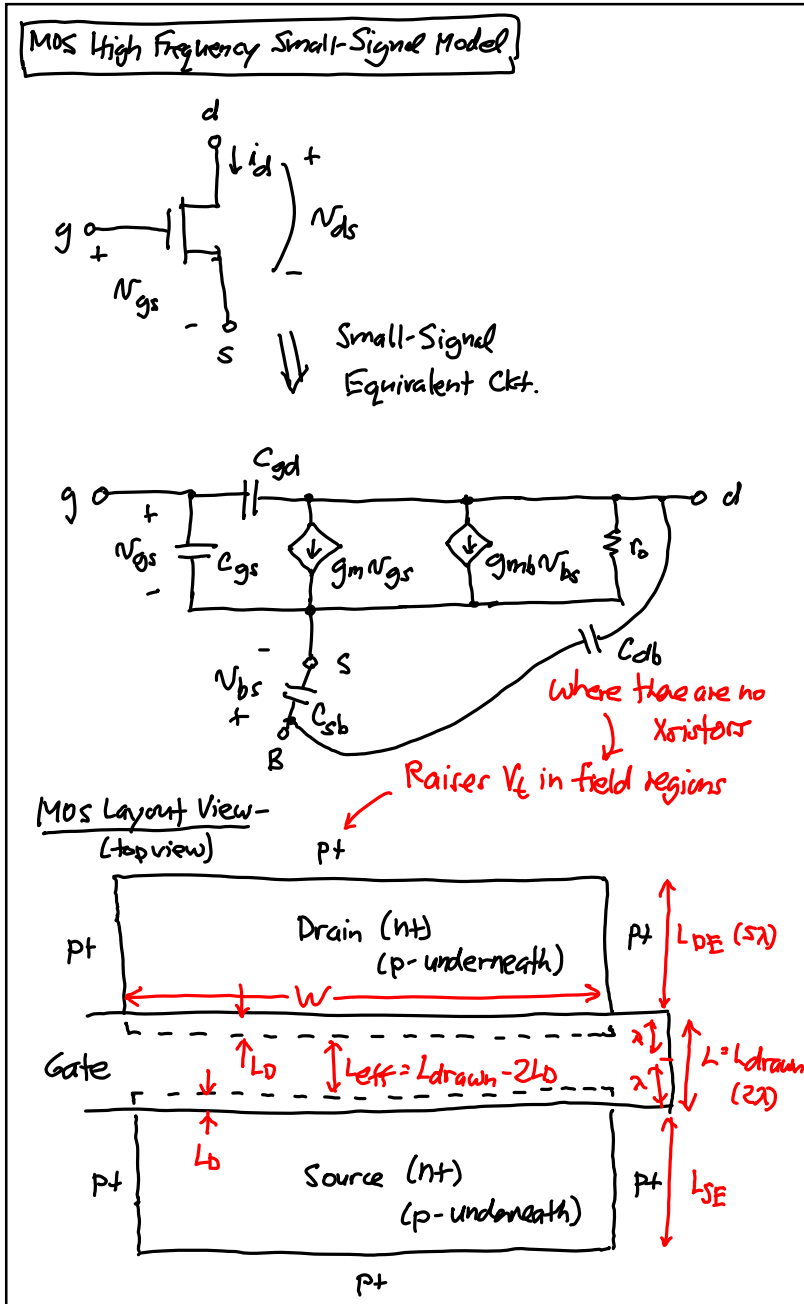
$$\omega_T = \beta_0 \omega_{\beta}$$

Figure of Merit for the frequency performance of a transistor

For Si bipolar:

$$f_T \sim 100 \text{ MHz} \rightarrow 50 \text{ GHz}$$

$$C_{\pi} = \frac{g_m}{\omega_T} - C_{\mu}$$



Gate-to-Source Capacitor, C_{gs} : (saturated MOS)

$$C_{gs} = C_{ox} + \frac{2}{3} W L_{eff} C_{ox}$$

$W L_{eff} C_{ox}$ ← accounts for the fact that the inversion charge is only under a portion of the gate

Gate-to-Drain Capacitance, C_{gd} :

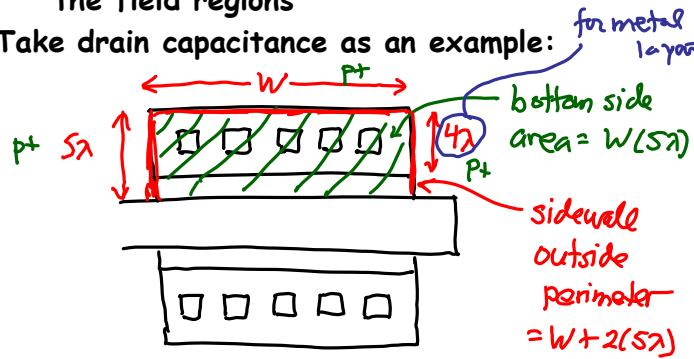
$$C_{gd} = C_{ox} \quad (\text{no inversion charge near the drain in the saturation region})$$

\uparrow
 $W L_{eff} C_{ox}$

- Source/Drain Junction Capacitance, C_{sb} & C_{db} :
- These are depletion capacitors associated with bulk-to-drain and bulk-to-source pn junctions
- Bottom capacitance per unit area differs from sidewall capacitance due to higher p+ bulk doping at the sidewalls
- The higher doping is near the silicon surface and designed to raise the threshold voltage in field areas, i.e., areas between transistors

↳ This way unwanted inversion does not occur in the field regions

- Take drain capacitance as an example: for metal layout



$$C_{db} = \frac{C_{db0}}{\sqrt{1 + \frac{V_{DS}}{\phi_j}}}$$

, where:

$$C_{db0}^A = \text{depletion capacitance w/ } V_{DS} = 0V$$

$$= (\text{junction bottom-side area}) \times C_{j0}$$

$$+ (\text{junction outside perimeter}) \times C_{jsw0}$$

$$= W(5\lambda) \times C_{j0} + (W + 2(5\lambda)) \times C_{jsw0}$$

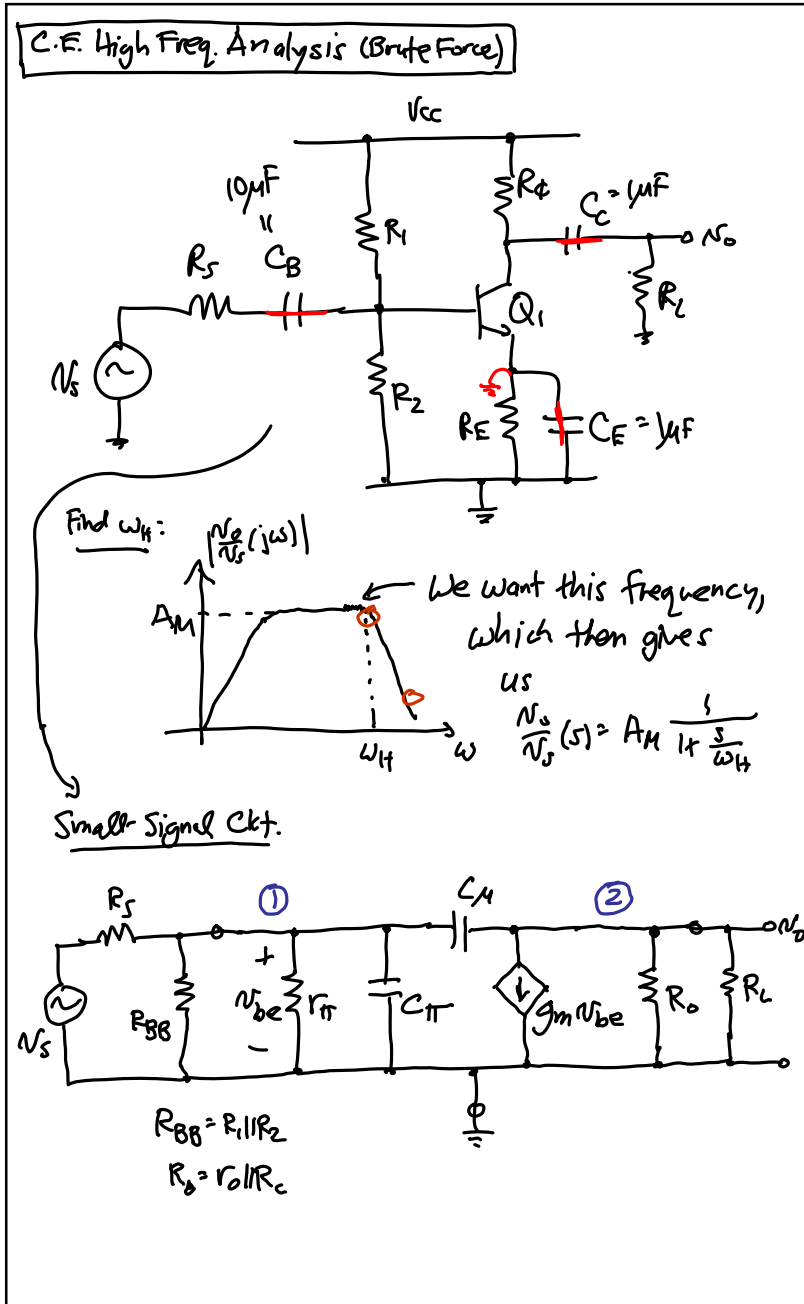
and where:

$$C_{j0} = \sqrt{\frac{q\epsilon_s N_B}{2\phi_j}} \quad \text{and} \quad C_{jsw0} = \sqrt{\frac{q\epsilon_s N_{CS}}{2\phi_j}} \times x_j$$

bulk doping concentration

channel-stop implant doping concentration

S/D junction depth



KCL ①: $\frac{N_s - V_{be}}{R_s} = \frac{V_{be}}{r_{\pi} || R_{BB}} + N_{be}(sC_{\pi}) + (V_{be} - N_o)(sC_{\mu})$

KCL ②: $(V_{be} - N_o)(sC_{\mu}) = g_m V_{be} + \frac{N_o}{R_o || R_L} = g_m V_{be} + R''$
 $[R'' = R_o || R_L]$

Rearrange:

$$\frac{N_s}{R_s} = V_{be} \left[\frac{1}{R_s} + \frac{1}{r_{\pi} || R_{BB}} + s(C_{\pi} + C_{\mu}) \right] - N_o(sC_{\mu})$$

$$\frac{1}{r_{\pi} || R_s || R_{BB}} = \frac{1}{R'}$$

...math...

$$\frac{N_o}{N_s}(s) = - \frac{g_m R' R''}{R_s} \left\{ \frac{(1 - s \frac{C_{\mu}}{g_m})}{(1 + s[R'(C_{\pi} + C_{\mu}) + R''C_{\mu} + g_m s R' R''])} \right\}$$

A_M
 $F_H(s)$
 models the freq. response

constant $\hat{=}$ midband gain

*

*
↓

$$F_H(s) = \frac{1 - s \frac{C_u}{g_m}}{1 + s [R'(C_{\pi} + C_u) + R'_e C_u + g_m R'_e R'''] + s^2 R' R'' C_{\pi} C_u}$$

$$= \frac{1 - \frac{s}{z_1}}{(1 - \frac{s}{p_1})(1 - \frac{s}{p_2})} = \frac{1 + \frac{s}{\omega_z}}{(1 + \frac{s}{\omega_{p1}})(1 + \frac{s}{\omega_{p2}})} = \frac{N(s)}{D(s)}$$

Zeros:
 ⇒ identify a RHP zero: $z_1 = + \frac{g_m}{C_u} \rightarrow \omega_z = \frac{g_m}{C_u}$

Note that $\frac{g_m}{C_u} \gg \omega_T$ (since $\frac{g_m}{C_u} \gg \frac{g_m}{C_{\pi} + C_u}$ and $C_{\pi} \gg C_u$)
 ↓
 ∴ ω_z is a very high freq.!
 ↳ can ignore relative to the lower freq. poles

Poles:
 ⇒ often, $\omega_{p1} \ll \omega_{p2}$ (i.e., one pole frequency is much lower than other)
 (ω_{p1} dominates over ω_{p2})

$$D(s) = (1 + \frac{s}{\omega_{p1}})(1 + \frac{s}{\omega_{p2}})$$

$$= 1 + s(\frac{1}{\omega_{p1}} + \frac{1}{\omega_{p2}}) + \frac{s^2}{\omega_{p1}\omega_{p2}}$$

$$\approx 1 + \frac{s}{\omega_{p1}} + \frac{s^2}{\omega_{p1}\omega_{p2}} \quad \left. \vphantom{\frac{s^2}{\omega_{p1}\omega_{p2}}} \right\} \text{Dominant pole approximation}$$

Comparing ω_1 the denominator of $F_H(s)$:

$$\omega_{p1} = -p_1 \approx \frac{1}{R'(C_{\pi} + C_u) + R'_e C_u + g_m R'_e R''}$$

↓

$$\omega_{p1} \approx \frac{1}{(C_{\pi} + C_u)(1 + g_m R'' + \frac{R''}{R'}) R'}$$

$$\omega_{p1}\omega_{p2} \approx \frac{1}{R' R'' C_{\pi} C_u} \Rightarrow \omega_{p2} \approx \frac{[C_{\pi} + C_u (1 + g_m R'' + \frac{R''}{R'})] R'}{R' R'' C_{\pi} C_u}$$

$$\therefore \omega_{p2} = \frac{g_m}{C_{\pi}} + \frac{1}{R'_e C_u} + \frac{1}{C_{\pi}} \left(\frac{1}{R'} + \frac{1}{R''} \right)$$

↑

$$\frac{g_m}{C_{\pi}} = \omega_T (= \frac{g_m}{C_{\pi} + C_u})$$

∴ ω_{p2} is indeed at a very high freq. $\gg \omega_{p1}$
 ∴ satisfies our dominant pole approx. ✓

Take a closer look @ the form of ω_{PI} :

$$\omega_{PI} \cong \frac{1}{\underbrace{R_1 C_{T1}}_{\tau_{T10}} + \underbrace{(R_1 + R_1' + g_m R_1 R_1') C_M}_{\tau_{T10}}}$$

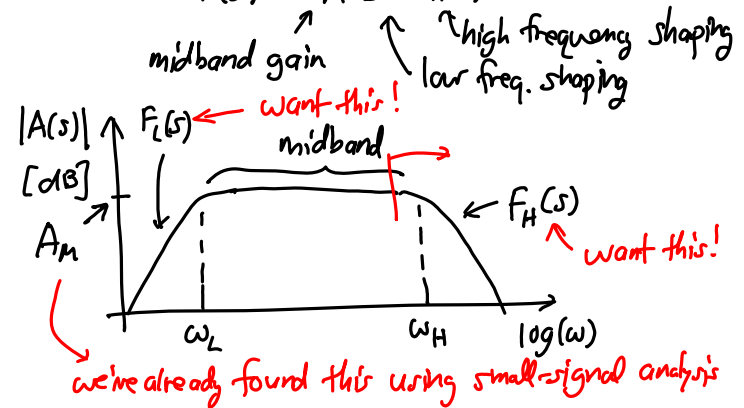
$$\tau_{T10} = R_{T10} C_{T1} \quad \tau_{T10} = R_{T10} C_M$$

Can think of the denominator as the sum of time constants associated w/ each of the capacitor in the ckt!

Freq. Response

Recall that the transfer function of a general amplifier can be expressed as a function of frequency via:

$$A(s) = A_M F_L(s) F_H(s) \quad s = j\omega$$

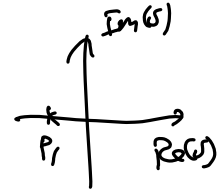


High Freq. Response Determination Using Open Ckt. Time Constant (OCTC) Analysis

In general:

$$F_H(s) = \frac{1 + a_1 s + a_2 s^2 + \dots + a_{n_z} s^{n_z}}{1 + b_1 s + b_2 s^2 + \dots + b_{n_p} s^{n_p}} \quad , n_p > n_z$$

$$= \frac{\prod_{j=1}^{n_z} (1 - \frac{s}{z_j})}{\prod_{i=1}^{n_p} (1 - \frac{s}{p_i})} = \frac{\prod_{j=1}^{n_z} (1 + \frac{s}{\omega_{zj}})}{\prod_{i=1}^{n_p} (1 + \frac{s}{\omega_{pi}})}$$



from which:

$$b_1 = \frac{1}{\omega_{p1}} + \frac{1}{\omega_{p2}} + \dots + \frac{1}{\omega_{pnp}} = \sum_{i=1}^{n_p} \frac{1}{\omega_{pi}} = \sum_{k=1}^{n_p} \tau_{pk}$$

↑
coeff. of the 1st order term

Through network theory, one can prove that: (see Gray & Meyer, Chpt. 7)

$$\sum_{i=1}^{np} \tau_{pi} = \sum_j C_j R_{j0} = \sum_j \tau_{j0} = b_1$$

Where C_j are capacitors in the H.F. ckt., i.e., small ones
 $R_{j0} \hat{=}$ driving pt. resistance seen between the terminals of C_j determined with

- ① all small (< 1nF) capacitors open-circuited
- ② all independent sources eliminated (i.e., short voltage sources, open current sources)
- ③ short all large (coupling/bypass) capacitors (i.e., > 1μF or > 1nF)

In calculating the H.F. response, we use the dominant pole approximation:

(i) $\omega_{p1} \ll \omega_{p2}, \dots, \omega_{pnp}$

(ii) $F_H(s) \cong \frac{1}{1 + b_1 s} = \frac{1}{1 + \frac{s}{\omega_H}}$

(ii) $\omega_H \cong \frac{1}{b_1} = \frac{1}{\sum_j C_j R_{j0}} = \frac{1}{\sum_j \tau_{j0}}$

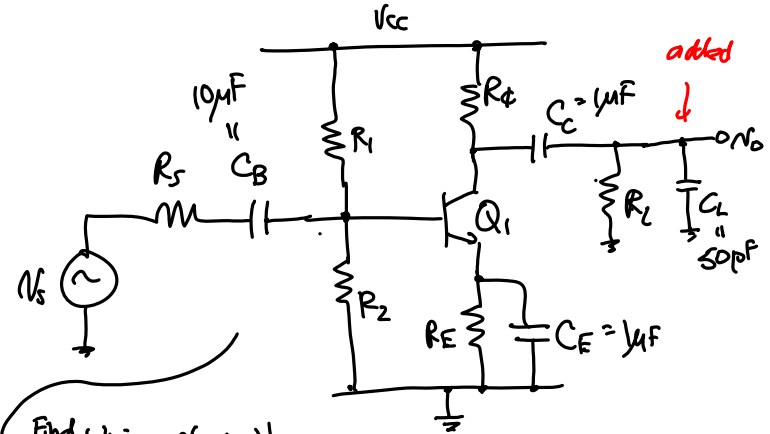
When there is no dominant pole, an approximate expression for ω_H is:

$$\omega_H \approx \frac{1}{\sqrt{\frac{1}{\omega_{p1}^2} + \frac{1}{\omega_{p2}^2} + \dots - \frac{1}{\omega_{z1}^2} - \frac{1}{\omega_{z2}^2} - \dots}}$$

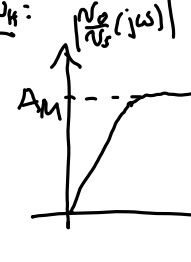
(just FYI)

- Now, use OCTC analysis on the CE amplifier to find the upper cut-off frequency, ω_H

C.E. High Freq. Analysis Using OCTC



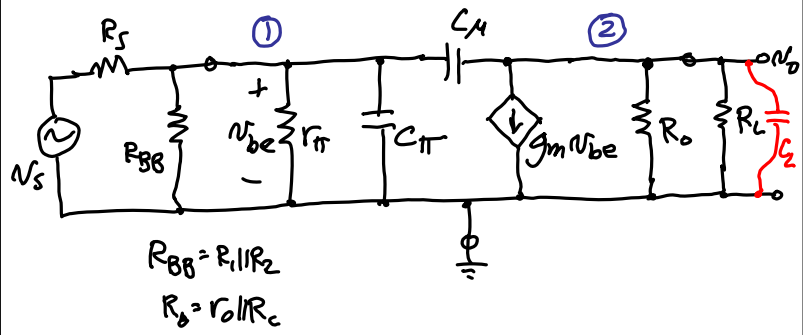
Find ω_H :



We want this frequency, which then gives us

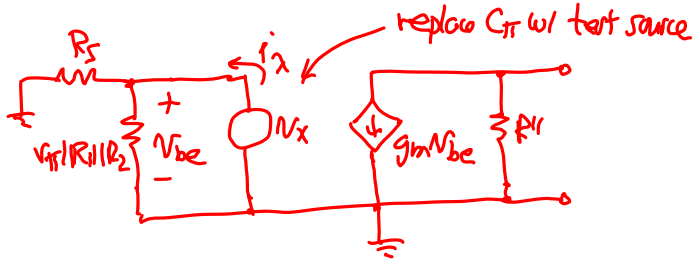
$$\frac{V_o}{V_s}(s) = A_M \frac{1}{1 + \frac{s}{\omega_H}}$$

Small-Signal Ckt.



(a) Determine $T_{\pi 0} = C_{\pi} R_{\pi 0}$

⇒ open ckt. all C's, zero out all sources, and determine the driving point impedance, $R_{\pi 0}$:

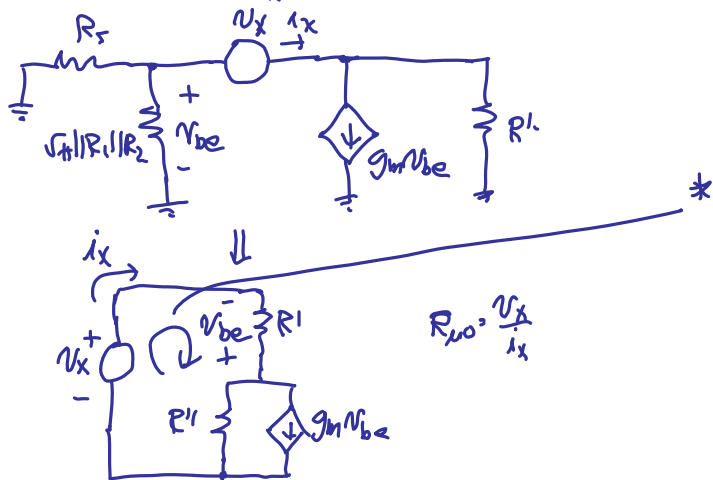


$$R_{\pi 0} = \frac{V_x}{i_x} = R_1 \parallel R_2 \parallel R_5 = R' \quad (\text{by inspection})$$

$$\therefore T_{\pi 0} = C_{\pi} R'$$

(b) Determine $T_{\mu 0} = C_{\mu} R_{\mu 0}$:

⇒ need $R_{\mu 0}$: replace C_{μ} w/ test source



$$R_{\mu 0} = \frac{V_x}{i_x}$$

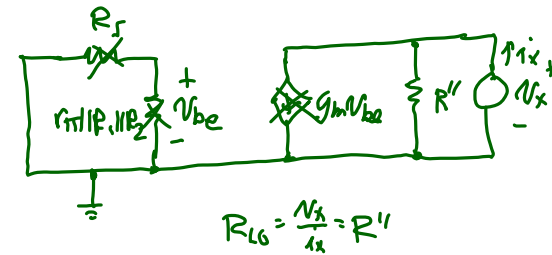
* KVL: $V_x = i_x R' + R''(i_x - g_m V_{be})$

$[V_{be} = -i_x R'] \rightarrow V_x = i_x R' + R''(i_x + i_x R' g_m)$

$$\therefore R_{\mu 0} = \frac{V_x}{i_x} = R' + R'' + g_m R' R''$$

$$\therefore T_{\mu 0} = C_{\mu} (R' + R'' + g_m R' R'')$$

(c) Determine $T_{L 0} = C_L R_{L 0}$:



$$R_{L 0} = \frac{V_x}{i_x} = R''$$

$$\therefore T_{L 0} = C_L R''$$

Thus:

$$\omega_H = \frac{1}{\sum_j T_{j 0}} = \frac{1}{C_{\pi} R' + C_{\mu} (R' + R'' + g_m R' R'') + C_L R''}$$

=