Bode Plot Tutorial

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1 Introduction

Although you should have learned about Bode plots in previous courses (such as EE40), this tutorial will give you a brief review of the material in case your memory is rusty.

2 Bode Plots Basics

Making the Bode plots for a transfer function involve drawing both the magnitude and phase plots. The magnitude is plotted in decibels (dB) and the phase is plotted in degrees. For both plots, the horizontal axis is either frequency ($f$) or angular frequency ($\omega$), measured in Hz and rad/s, respectively. The horizontal axis should be logarithmic (i.e. increasing by powers of 10).

Most of the transfer functions we’ll deal with in this class can be separated into a general that resembles the following:

$$H(j\omega) = A \cdot \frac{j\omega/\omega_{z1}}{j\omega/\omega_{p1}} \frac{(1 + j\omega/\omega_{z2})(1 + j\omega/\omega_{z3})}{(1 + j\omega/\omega_{p2})(1 + j\omega/\omega_{p3})} \ldots$$

(1)

$A$ is an arbitrary constant and $j$ is $\sqrt{-1}$. As you can see, the basic component of this transfer function appears to be $1 + j\omega/\omega_c$, where $\omega_c$ is some constant (with the slight variation $j\omega/\omega_c$). Let’s analyze this basic component first before we analyze the entire transfer function.

2.1 Magnitude

Recall that the definition of magnitude (measured in dB) is as follows:

$$20 \log |H(j\omega)| = 20 \log \sqrt{\Re[H(j\omega)]^2 + \Im[H(j\omega)]^2}$$

Let’s apply this definition to our basic transfer function component (this is called a zero when it appears in the numerator of a transfer function):

$$20 \log |1 + j\omega/\omega_c| = 20 \log \sqrt{1 + (\omega/\omega_c)^2}$$

For small $\omega$, we have $20 \log |1 + j\omega/\omega_c| \approx 0$ dB. For large $\omega$, $20 \log |1 + j\omega/\omega_c| \to \infty$. When $\omega = \omega_c$, the magnitude of the transfer function is approximately 3 dB. Since there’s so little change from $\omega = 0$ to $\omega = \omega_c$, we approximate the magnitude in this region as a constant 0 dB.
For $\omega > \omega_c$, the $(\omega/\omega_c)^2$ dominates the magnitude expression, allowing us to approximate the magnitude as $20 \log \omega/\omega_c$. From this expression it’s clear that if we increase $\omega$ by a factor of 10, we increase the magnitude by 20 dB. Thus, our Bode plot approximation for the zero is a constant 0 dB for $\omega < \omega_c$ and a line constantly increase by 20 dB/decade for $\omega > \omega_c$, illustrated in Figure 1.

Figure 1 also illustrates the Bode plot for a DC zero of the form $j\omega/\omega_c$. This differs only slightly from the normal zero in that it lacks the additional 1. Thus, instead of having the constant magnitude region for $\omega < \omega_c$, it simply always increases at 20 dB/decade. We draw its intersection with the frequency axis where $\omega = \omega_c$, since that’s where the magnitude is 0 dB.

![Figure 1: Bode plots (magnitude) for a normal zero and a DC zero for $\omega_c = 10^5$ rad/s (the plots overlap for $\omega > \omega_c$)](image)

The basic transfer function component $1 + j\omega/\omega_c$ can also appear in the denominator (in which case it is called a pole). Although this may seem like an entirely different problem, recall that we take the logarithm of our transfer function because our result is expressed in decibels. Taking the logarithm of the inverse of a function simply gives the negative logarithm of the function, meaning we simply have to negate the results of our zero analysis to get the appropriate expressions for poles. The same argument applies with DC poles of the form $j\omega/\omega_c$, so we can negate our DC zero analysis to get the DC pole results.

A normal pole will have a constant 0 dB value for $\omega < \omega_c$ and will drop by 20 dB/decade for $\omega > \omega_c$. A DC pole will drop by 20 dB/decade for any $\omega$ and will intersect the frequency axis (0 dB) at $\omega = \omega_c$. The results are shown in Figure 2.

2.2 Phase

Let’s take a look at the phase of a zero, DC zero, pole, and DC pole. Recall the definition of phase:

$$\text{Arg}(H(j\omega)) = \tan^{-1} \frac{\Re[H(j\omega)]}{\Im[H(j\omega)]}$$
Let’s apply this to the normal zero first.

\[ \text{Arg}(1 + j\omega/\omega_c) = \tan^{-1} \frac{\omega}{\omega_c} \]

For \( \omega = 0 \), \( \text{Arg}(1 + j\omega/\omega_c) = 0 \). For \( \omega \to \infty \), \( \text{Arg}(1 + j\omega/\omega_c) \to 90^\circ \). For \( \omega = \omega_c \), \( \text{Arg}(1 + j\omega/\omega_c) \to 45^\circ \).

Thus, our approximation for the phase of a zero is \( 0^\circ \) for \( \omega < 0.1\omega_c \), \( 45^\circ \) for \( \omega = \omega_c \), and \( 90^\circ \) for \( \omega > 10\omega_c \) with a straight line connecting these points. We can also look at the phase of a DC zero, which is always \( 90^\circ \). These results are shown in Figure 3.

Similar to our analysis of the magnitude, we can also consider poles and DC poles in our phase plots. It can be shown that \( \tan^{-1} -\theta = -\tan^{-1} \theta \), meaning our phase plots for poles and DC poles will simply be negated versions of the zero plots. These are shown in Figure 4.

## 3 Combining Poles and Zeroes

Generally, a transfer function may involve many poles and zeroes (and their DC counterparts). In order to make it easier to draw Bode plots, your first step should be to factor the transfer function into the canonical form shown in Equation 1. This makes it easy to identify all of the poles and zeroes.

First, you’ll have to handle the constant \( A \) in front (if present). The magnitude of \( A \) will affect your magnitude plot, and the sign of \( A \) will affect your phase plot. Your magnitude plot must be shifted up by \( 20 \log |A| \). For example, if \( A = 10 \), then your magnitude plot must be shifted up by 20 dB. Similarly, if \( A = 1/10 \), then your magnitude plot must be shifted down by 20 dB. If \( A < 0 \), then your phase plot must be shifted up (or down—it’s the same in this case) by 180°.

Second, you need to draw each pole and zero individually on the same set of axes (whether you’re making a magnitude or phase plot).
Third, you simply add the curves that you’ve drawn at each point to obtain the final Bode plot. Remember to shift your plots accordingly based on the constant $A$ as mentioned previously. This superposition principle is possible because of the decomposition of the transfer function into zeroes and poles.

When adding the poles and zeroes in the final plot, remember that in areas where two curves are constant, the result will just be the sum of the constant values. When one is a constant and one is linear, then the result will start at the constant value and have the slope of the linear curve. Finally, when both are linear, the sum will have a slope equal to the sum of the slopes of the individual curves.
Figure 4: Bode plots (phase) for a normal pole and a DC pole for $\omega_c = 10^5 \text{ rad/s}$ (the plots overlap for $\omega > 10\omega_c$)