

$$\frac{1}{g(z)} = \frac{1}{f(z)}$$

NOTE ON "GAUSSIAN BEAMS"

- what are we really talking about here?

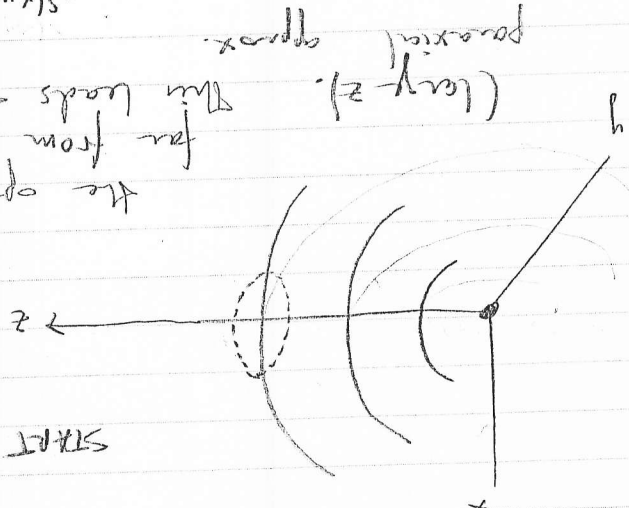
mathematically

START: We wish to describe

the spatial phase evolution of a spherical wave in a special location close to

the optic axis (small x^2+y^2) and far from the origin of the source

(large z). This leads to the validity of the paraxial approx.



$\int S(x,y,z) e^{ikr-i\omega t}$

$$E_{\text{spherical}}(x,y,z,t) = A_0 e^{\frac{r}{\lambda}}$$

look at r , assume $z \gg \sqrt{x^2+y^2}$

$$r = \sqrt{x^2+y^2+z^2} = z \sqrt{1 + \frac{x^2+y^2}{z^2}} = z \left(1 + \frac{1}{2} \frac{x^2+y^2}{z^2} + \dots \right)$$

TAYLOR EXPANSION

$$\text{for } x^2+y^2 \ll z^2, \quad \frac{x^2+y^2}{z^2} \ll 1 \quad \therefore \quad \frac{1}{2} \left[\frac{x^2+y^2}{z^2} \right] \ll \frac{1}{2} \left[\frac{x^2+y^2}{z^2} \right]^2$$

so neglect this term & all other H.O.T. in the PHASE and just let $r \approx z$ in the amplitude.

$$E_{\text{spherical}}(x,y,z,t) \approx A_0 e^{ikz + i\frac{x^2+y^2}{2z} - i\omega t}$$

$$\approx A_0 e^{ikz - i\omega t} \exp \left\{ i\frac{k}{2z} (x^2+y^2) \right\}$$

QUADRATIC PHASE APPROX TO 1-DIMENSIONAL PARAXIAL

Mathematically

$$\therefore R_r(z) = \frac{z^2 + z_0^2}{z}$$

$$\frac{\lambda}{\pi n W^2(z)} = \frac{z_0}{z^2 + z_0^2}$$

$$W^2(z) = \frac{\lambda}{\pi n} \frac{(z^2 + z_0^2)}{z_0}$$

$$W_0^2 = W^2(z=0)$$

$$W^2(z) =$$

Why? adding an imaginary part to z allows us to physically create a field whose amplitude can change ~~the~~ ~~value~~ in z and in x^2+y^2 (which we know is true in reality)

So, make this "weird" substitution & try to see if it leads to any consequences that make sense.
 { obey reality.

At this point I know this makes no sense!

$$\underbrace{\exp\left\{\frac{ik}{2} \frac{z}{z^2 + z_0^2} (x^2 + y^2)\right\}}_{\text{spatial phase evolution}} \underbrace{\exp\left\{-\frac{k}{2} \frac{z_0}{z^2 + z_0^2} (x^2 + y^2)\right\}}_{\text{spatial field amplitude evolution}}$$

MAKE IT LOOK "PRETTIER" } more like a spherical wave
 } a gaussian beam

$$\underbrace{\exp\left\{\frac{ik(x^2+y^2)}{2R(z)}\right\}}_{\text{spherical-like}} \underbrace{\exp\left\{-\frac{(x^2+y^2)}{W^2(z)}\right\}}_{\text{gaussian-like}}$$

$$\frac{z_0}{z^2 + z_0^2} = \frac{2}{kW^2(z)}$$

→ NOTE: when arg of exp = -1, the field amplitude is $\frac{1}{e}$ of its peak at $x=y=0$.

$$\text{DEF: } \frac{kz_0}{2(z^2 + z_0^2)} = \frac{1}{W^2(z)} \text{ so that when } r^2 = W^2(z)$$

field is $\frac{1}{e}$ of its original value.

$$W(z) = \frac{1}{e} \text{ radius of beam}$$

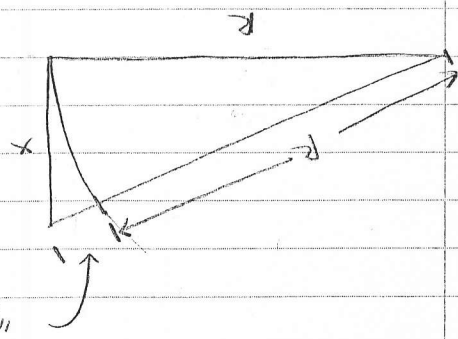
↳ The ray in class all the time.

WHAT DOES THIS PICTURE TELL ABOUT? WHY IS z LIKE THE RADIUS OF CURVATURE?

RECALL: (Go to RECALL SECOND)

Consider 1-D example spherical (cylindrical) wave

"Extra" optical path length for wavefront to reach the plane of interest @ x w.r.t. the dist R travelled for the axial ray. ($x=0$)



Assume $x \ll R$.

$$OPL(x) = \sqrt{x^2 + R^2} - R = R \sqrt{1 + \frac{x^2}{R^2}} - R$$

$$= R \left(1 + \frac{x^2}{2R^2} + \frac{1}{2} \left(\frac{x^2}{R^2} \right)^2 + \dots \right) - R$$

ALL TERMS OF ORDER EXP

$$= \frac{x^2}{2R} + \frac{x^4}{8R^3} + \dots$$

So, when we drop all H.O.T. we get

$$OPL(x) \approx \frac{x^2}{2R}$$

which says that the spherical wavefront

must travel $\frac{x^2}{2R}$ ~~more than the distance~~

to reach the plane of interest at a dist x off the optical axis. EXTRA to 3-D

$$\text{Ray must travel } R + \frac{x^2 + y^2}{2R}$$

$R =$ dist from apparent sphere point source.

"effective radius of curvature dz "

so: $R(z) = \frac{z^2 + z_0^2}{z}$

$$w^2(z) = \frac{(z^2 + z_0^2)z}{kz_0} = \frac{2\lambda z_0^2}{2\pi n z_0} \left(1 + \left(\frac{z}{z_0} \right)^2 \right)$$

look at $w^2(z=0) = w_0^2$

$$w^2(z) = w_0^2 \left(1 + \left(\frac{z}{z_0} \right)^2 \right) = \frac{1}{e^2} \text{ width of radius as a function of } z.$$

From math, we see that no matter what, w is smallest at $z=0$
 where $w_0^2 = w^2(z=0) = \frac{2\lambda z_0}{2\pi n} = \frac{2z_0}{k}$

(PHYSICALLY, w_0 is the $\frac{1}{e}$ radius of the beam at $z=0$)

OK, what the hell does this all mean?

$w_0 = w(z=0)$ is the $\frac{1}{e}$ radius of the beam at $z=0$. OK, so what? Why is $z=0$ significant?

$$R(z=0) = \frac{z^2 + z_0^2}{z} = \infty !!!$$

The radius of curvature is INFINITE at $z=0$
 AND the waist is THE SMALLEST IT CAN BE (at least the math says so) at $z=0$

LET'S TALK REALITY. What can we measure? How about w_0 ?

INTERPRETATION:

~~W(z) = W_0 \sqrt{1 + \left(\frac{z}{z_0}\right)^2}~~

$$W^2(z) = W_0^2 \left(1 + \left(\frac{z}{z_0}\right)^2 \right)$$

$$W^2(z) = W_0^2 \left[1 + \left(\frac{z\lambda}{W_0^2 \pi n} \right)^2 \right]$$

z_0 is the distance from the ~~min-waist~~ ~~position of the waist~~ ~~longitudinal~~ position where the waist has expanded to $\sqrt{2}$ times its original value.

$z_0 = \frac{W_0^2 \pi n}{\lambda}$ $\therefore$ the smaller W_0 , the quicker the beam ^{waist} expands to $\sqrt{2}$ of its min value.

WHAT ABOUT $R(z)$ the "effective radius of curvature"

$$R(z) = z \left[1 + \left(\frac{z_0}{z}\right)^2 \right] = \text{Radius of curvature changes as a function of } z$$

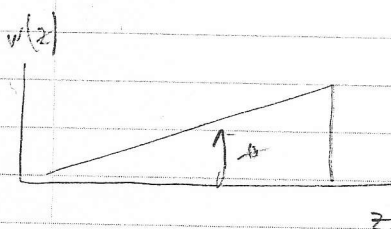
$$R(z) = z \left[1 + \left(\frac{\pi n W_0^2}{\lambda_0 z} \right)^2 \right]$$

CHECK LIMIT $z \gg z_0$

$$R(z) \approx z \quad \left[1 \gg \left(\frac{z_0}{z}\right)^2 \right]$$

$$W(z) \approx \frac{W_0 z}{z_0} \quad \left[1 \ll \left(\frac{z}{z_0}\right)^2 \right]$$

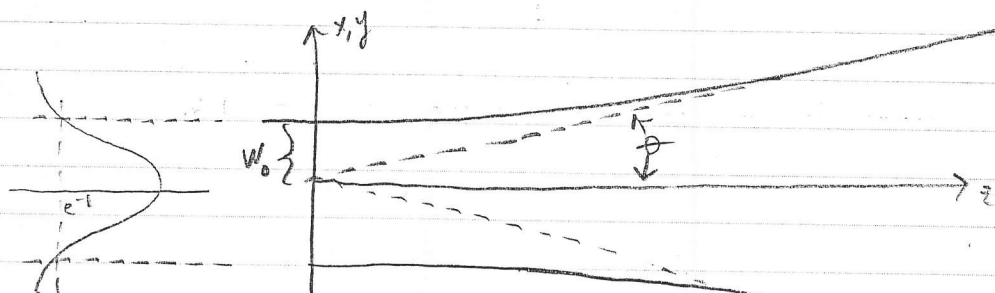
and note: $\frac{d[W(z)]}{dz} = \text{slope @ } z \approx 0$



$$\theta = \frac{W_0}{z_0} = \frac{W_0 \lambda}{W_0^2 \pi n} = \frac{\lambda}{W_0 \pi n}$$

So, at $z=0$ it's as if you're observing the phase from a point source a dist $R(z=0) = \infty$ away!

As z increases, $R(z)$ remain very large but quickly approaches the value z so ~~at~~ for $z \gg z_0$ it always looks like the effective point source radiator is at $z=0$!



EX: (A) A commercial He-Ne laser is advertised to have a far field divergence angle of 1 mrad. $\lambda_0 = 632 \text{ nm}$. What is spot size w_0 ?

Farfield $z \gg z_0$ half angle $\theta = 1 \text{ mrad}$.

$$1 \text{ mrad} = \frac{\lambda}{w_0 \pi}$$

$$w_0 = \frac{\lambda}{10^{-3} \pi}$$

$$= \frac{506 \cdot 10^{-9}}{10^{-3} \pi (1)} = \frac{506 \cdot 10^{-6}}{\pi} \approx 150 \mu\text{m}$$

(3) Power = 5 mW what is intensity $\left(\frac{W}{m^2}\right)$ @

waist ($z=0$)? Spot size area = πw_0^2

$$\therefore I = \frac{\text{POWER}}{\text{AREA}} = \frac{5 \text{ mW}}{\pi w_0^2 \text{ m}^2} = \frac{5 \cdot 10^{-3}}{\pi (150)^2 (10^{-6})^2} =$$