

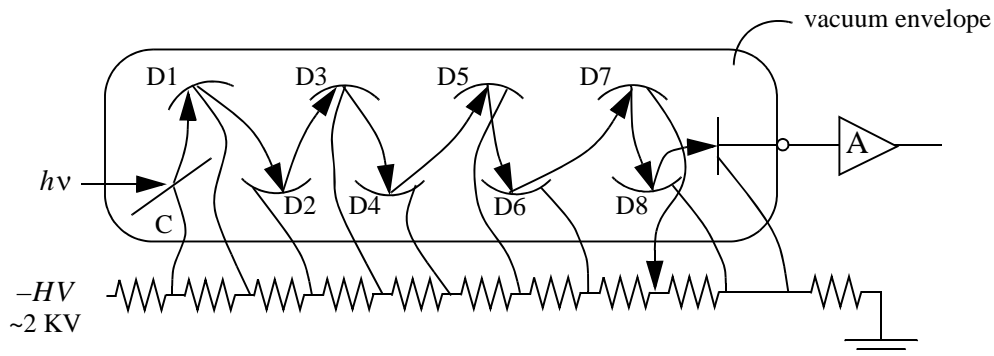
Lecture 12

OPTICAL DETECTORS

(Reference: *Optical Electronics in Modern Communications*, A. Yariv, Oxford, 1977, Ch. 11.)

Photomultiplier Tube (PMT)

- Highly sensitive detector for light from near infrared → ultraviolet
- Can detect as little as 10^{-19} Watt!



Voltage divider chain to bias dynodes

R chosen to have $\sim 100\text{V}$ drop per dynode

Photocathode C:– absorbs photon \rightarrow ejects electron

– work function ϕ is the minimum energy needed to eject an electron

– the photon energy must exceed the work function $h\nu > \phi$ to get photoelectrons

Dynodes D1-D8:Secondary electron emission. Electron from cathode accelerated by ~ 100 eV.

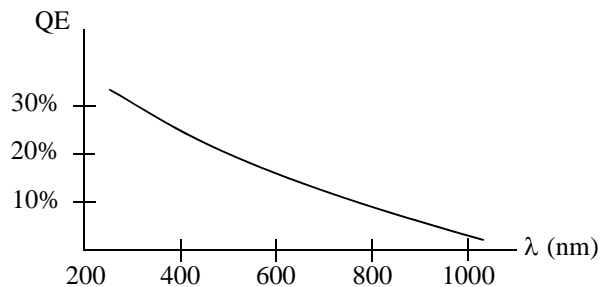
Impact into dynode surface causes ejection of multiple electrons, $\delta \cong 5$.

For N dynodes, the total gain is then δ^N .

Photocathode quantum efficiency: $QE \equiv$ probability a photon will eject one electron

$$0 \leq QE \leq 1$$

Typical photocathode response



Sensitivity:

For 10 dynodes, $\delta = 5$, $G = 5^{10} \cong 10^7$.

Take 2eV photons (620 nm), 1 picoW = $10^{-12}\text{W} = 10^{-12}\text{J/s}$

$$1 \text{ eV} = 1.6 \times 10^{-19}\text{J}, \text{ so, } 10^{-12}\text{W} = 3 \times 10^6 \text{ photons/s}$$

With $QE = 30\%$, Anode current is

$$(3 \times 10^6)(0.3)10^7 = 10^{13} \text{ electrons/s} = 1.6 \times 10^{-6} \text{ A}$$

Phototube **dark current**: 1) random thermal excitation of electrons from photocathode

2) cosmic rays, ambient radioactivity

- Thermal excitation rate is proportional to $e^{-\Phi/kT}$, where Φ represents the cathode work function

so lower work function \rightarrow IR sensitivity, **but** larger dark current

- For room temperature, typical cathode dark current, I_{cd} , is $\approx 10^4$ electrons/sec. Anode dark current is then $I_{ad} = I_{cd} \cdot G$

- Dark current sets a lower limit to phototube sensitivity to low light levels. To distinguish a light signal above the background dark current, the photoelectric cathode current must exceed the dark current. If I_{cd} is $\approx 10^4$ e/sec, then the sensitivity to light can be $\sim 3 \times 10^4$ photons/sec (assuming $QE = 30\%$). 3×10^4 red photons/s $\cong 10^{-14}$ W!

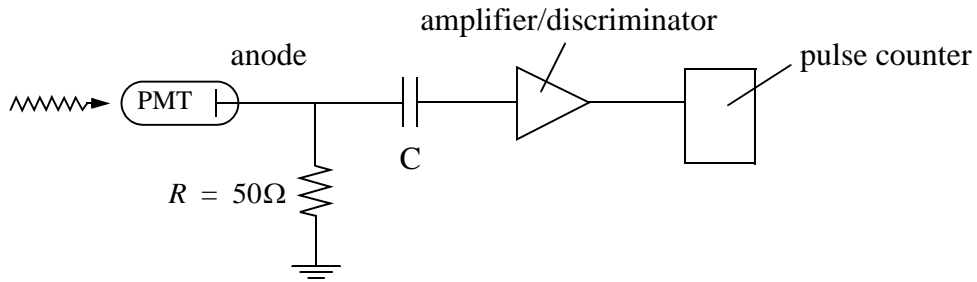
- Dark current can be reduced by *cooling*. Using thermoelectric cooling $T \cong -40\text{C}$ is easily obtained. Assume a work function of $\Phi = 1.5 \text{ eV}$

$$\frac{I_{cd}(260\text{K})}{I_{cd}(300\text{K})} = \frac{e^{-1.5/0.0225}}{e^{-1.5/0.026}} = e^{-1.5(44.4 - 38.5)} = e^{-8.8} \cong 1.4 \times 10^{-4}$$

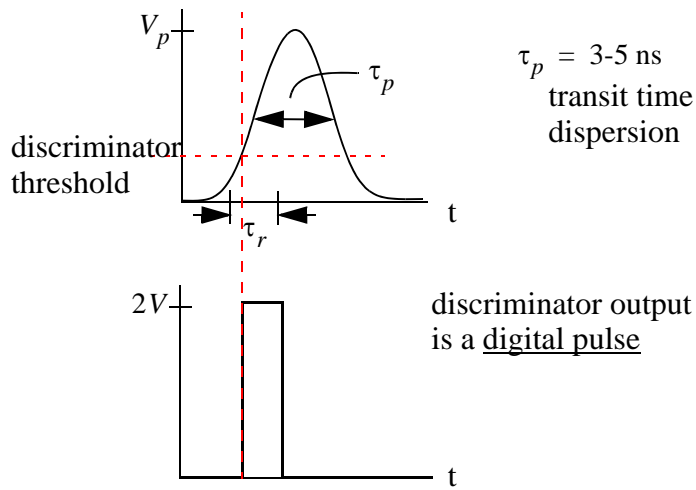
Dark current is reduced by this amount! \rightarrow down to ~ 1 e/sec. Minimum detectable power become $< 10^{-18}$ W!

Photon counting: PMT is so sensitive, we are really counting photons. Often, PMT circuits are specifically optimized to do this.

Photon counting system:

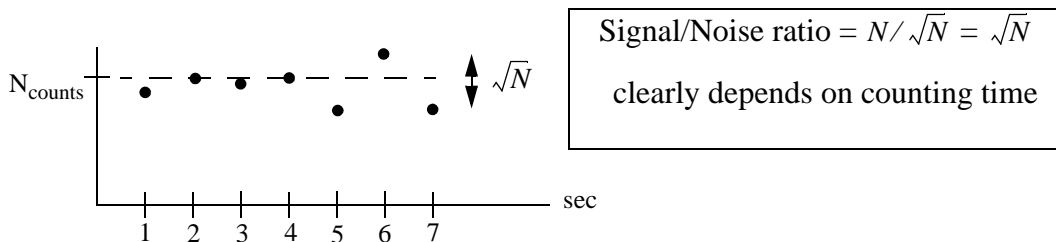


PMT output pulse



- How big is the PMT output pulse from one photon? For $G \sim 10^7$, we get 10^7 electrons $\cong 10^{-12}\text{C}$. For $\tau_p = 10^{-8}$ sec, $I_{apk} \sim 10^{-4}\text{A}$. For $R = 50\Omega$, $V_p \sim 5 \text{ mV}$.
- Discriminator eliminates electrical noise in $< 1\text{mV}$ range. V_p has a variation due to statistical nature of gain process. Discriminator also eliminates this.

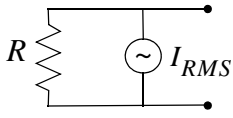
Shot noise: Photon arrival is always statistical. Generally it follows Poisson statistics. Then if the photon arrival rate is N ph/sec, and we count for 1 sec, we get N on average. The standard deviation will be found to be \sqrt{N} . This means we have noise.



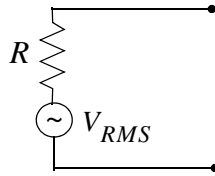
Shot noise is universal for light detection. Even if photons are not explicitly counted, the shot noise is a fundamental limit. It is most significant at low light levels, though, due to \sqrt{N} dependence.

Johnson noise: Random thermal noise in any resistor, R

$$I_{RMS} = \sqrt{\frac{4kT}{R}B}$$



$$V_{RMS} = \sqrt{4kTRB}$$

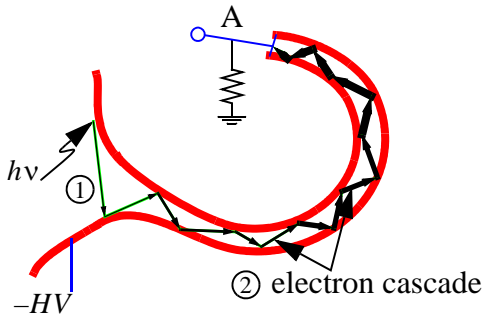


B: bandwidth (Hz)

Equivalent model of noise as either current or voltage source.

Channel Electron Multiplier (Channeltron)

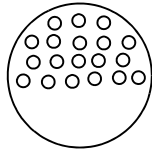
Single monolithic device functions as a PMT:



Glass tube bent around curve. One end open as a funnel shape. Coating acts as photocathode **and** secondary electron emitter. Also, coating has high, but not infinite electrical resistance.

1. Photon hits funnel portion
2. Electrons are accelerated into the bent tube by bias field
3. Secondary electron emission gives gain at each electron collision with wall
4. Must be operated in vacuum
5. Typical gain $\sim 10^4$
6. More compact and rugged than PMT

Microchannel plate MCP – array of channeltrons

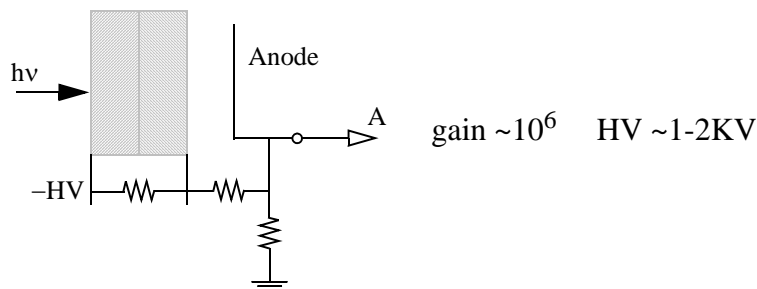


10-20 μ m separation

each hole \sim 5-10 μ m diameter

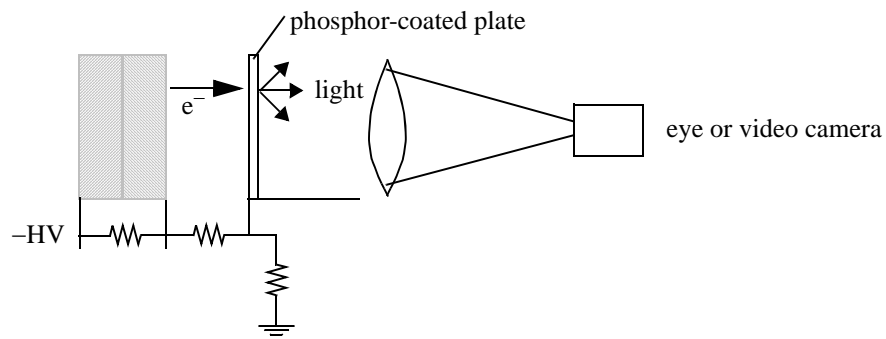
– each channel is a miniature channeltron

– gain $\sim 10^3$



Dual Plate MCP Detector

MCP Image Intensifier



Single or Dual Plate MCP

- Electrons accelerated out of back of MCP into phosphor
- Phosphor QE ~50% photons/electrons

Image intensification $QE_{MCP} \cdot G_{MCP} \cdot QE_{phosphor} \approx 10^5$

→ Night vision goggles