

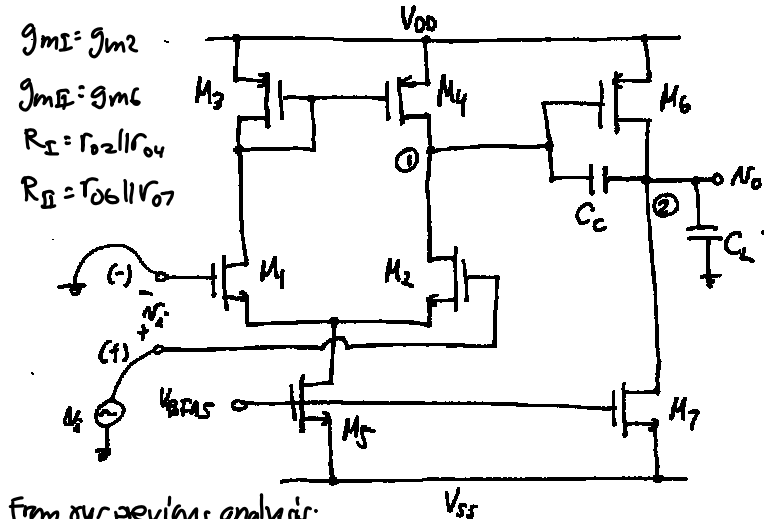
Lecture 22: RHP Zero

• Announcements:

- HW#9 due next Wednesday at 8 a.m.
 - Lab#3 (Design Project) in progress
 - Design Project Checkpoint:
 - ↳ Due Tuesday, Nov. 17, 11:59 p.m.
 - ↳ Send to your TA a spice file for your op amp design that simulates correctly, i.e., that reaches a stable bias point where all transistors are saturated (or linear if an MOS resistor)
 - ↳ It doesn't need to meet the project specs, but it should simulate correctly
 - Thursday (next week) I'll be traveling
 - ↳ Lecture by videotape
 - Handout on "RHP Zero" online
 - Lecture Topics:
 - ↳ CMOS Op Amp Compensation
 - ↳ Nulling the RHP Zero
-
- Last Time: Choosing C_c

over

CMOS 2-Stage Op Amp Compensation (Summary)



from our previous analysis:

$$p_1 = -\frac{1}{g_{mI} R_I R_{II} C_c} \quad [C_c \gg C_I \text{ or } C_{II}] \quad [C_c \gg C_I]$$

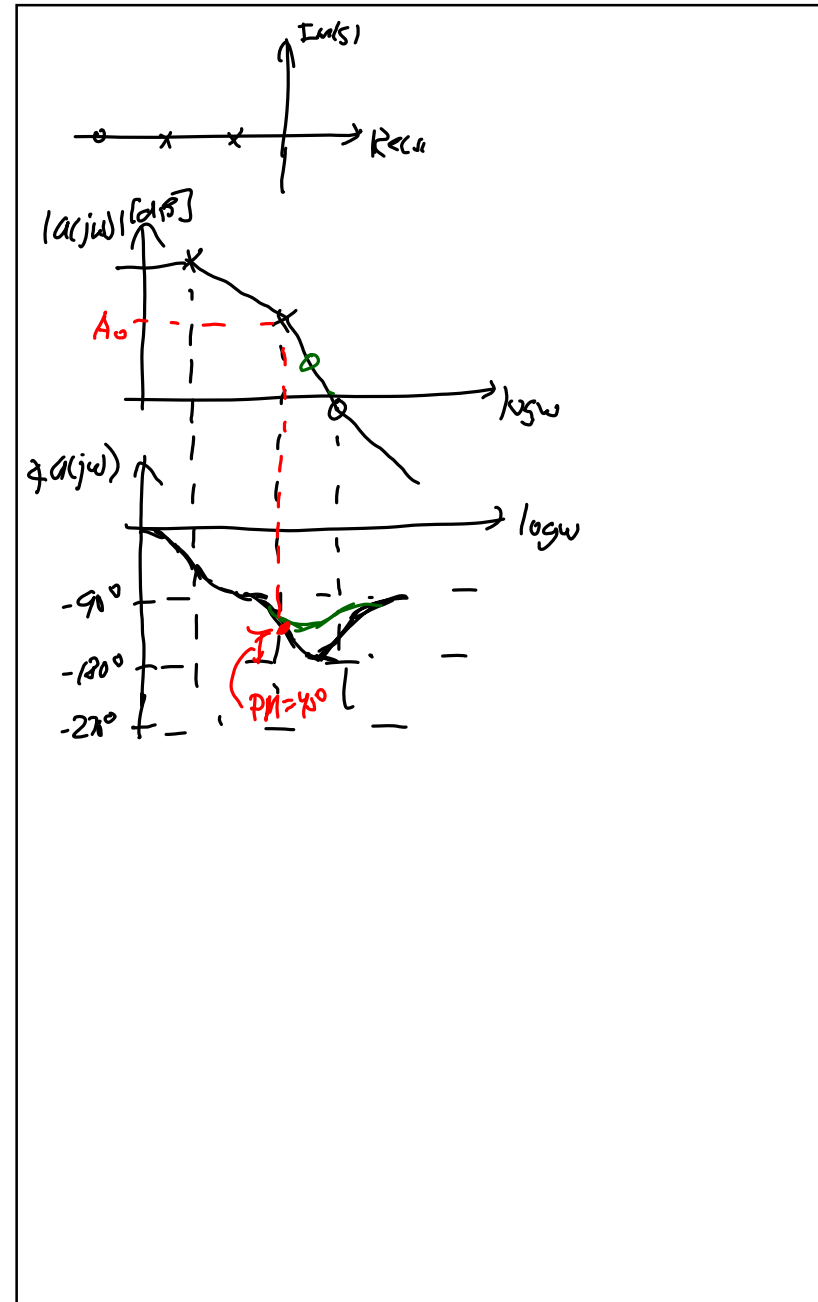
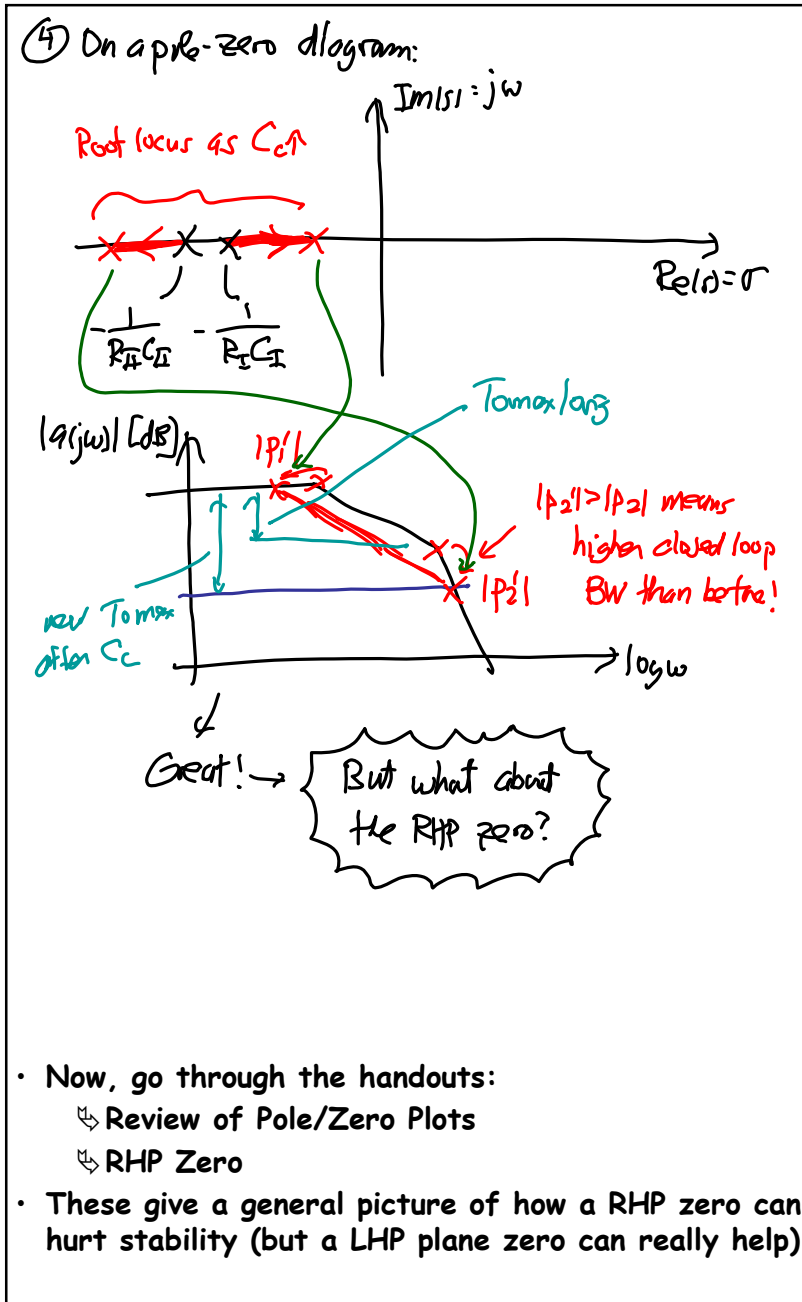
$$p_2 = -\frac{g_{mII} C_c}{C_I C_{II} + C_c (C_I + C_{II})} \approx -\frac{g_{mII}}{C_I + C_{II}} \approx -\frac{g_{m6}}{C_L}$$

$$z = +\frac{g_{mII}}{C_c} \leftarrow \text{RHP zero (this will cause problems)}$$

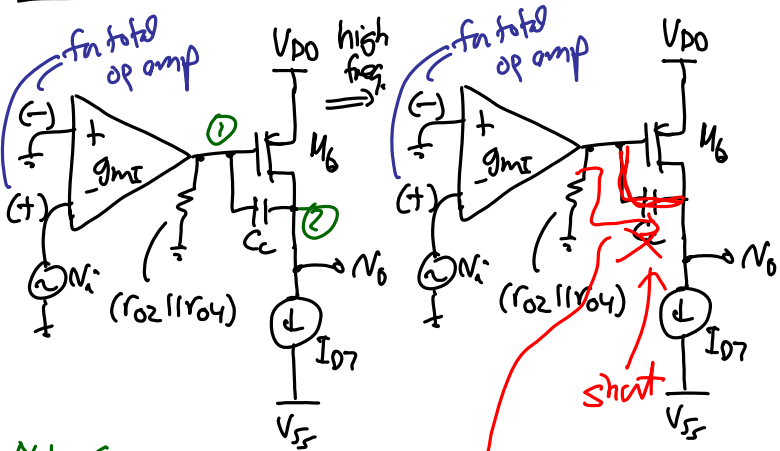
Remarks:

- ① Note that as $C_c \uparrow \rightarrow |p_1| \downarrow$
- ② As $C_c \uparrow \rightarrow |p_2| \uparrow \rightarrow |p_2| = \frac{g_{mII}}{C_I + C_{II}}$ } pole-splitting
- ③ With $C_c = 0$ (i.e., before compensation)

$$p_1 = -\frac{1}{R_I C_I}, \quad p_2 = -\frac{1}{R_{II} C_{II}}$$



Where Does the RHP Zero Come From?



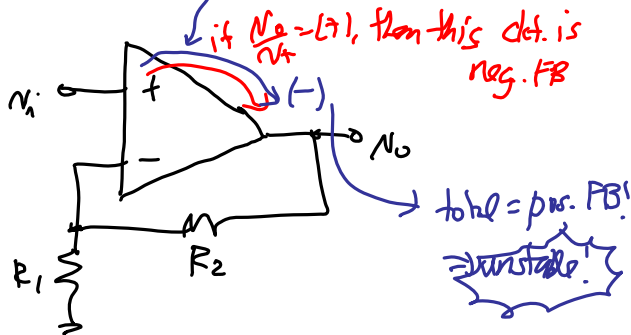
At low freq.:

$$\frac{V_o}{V_i} = -g_{m6} (r_{o6} || r_{o4})$$

$$\text{total gain} = \frac{V_o}{V_i} = \frac{V_o}{V_+} \cdot \frac{V_+}{V_i} = (+)$$

At high freq., get feed-forward path

$$\frac{V_o}{V_i} = \frac{V_o}{V_+} = -\frac{g_{m1}}{g_{m6}} = (-)$$

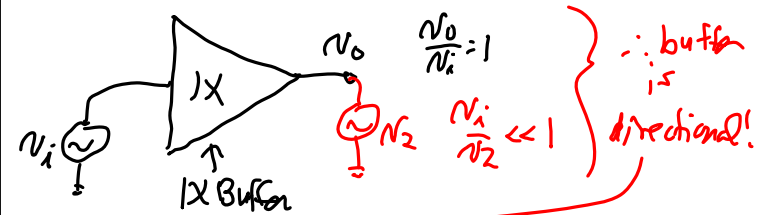


Observation.

Miller effect compensation requires neg. FB path.

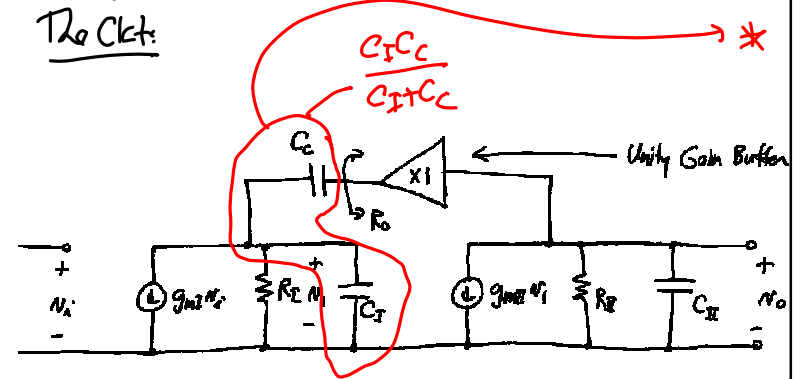
↳ BUT: The feed-forward path (that creates the zero) is not needed!

- Solution:
- ① Kill the feedforward path.
 - ② Keep the feedback path.



↳ Solution: Put a 1x buffer in series w/ Cc to prevent feedforward, but allow FB!

The Ckt:



Apply KCL:

$$P_1 \approx -\frac{1}{g_{mII} R_T R_{II} C_c} \quad (\text{same as before})$$

$$P_2 \approx -\frac{g_{mII} C_c}{C_{II}(C_I + C_c)} \approx -\frac{g_{mII}}{C_{II}} \quad [C_c \gg C_I]$$

$$* P_3 \approx -\frac{1}{R_o(C_I C_c / (C_I + C_c))} \approx -\frac{1}{R_o C_I}$$

series comb of C_I & C_c

$$z_1 \approx -\frac{1}{R_o C_c} \leftarrow \text{LHP zero!} \rightarrow \text{Good!}$$

Remarks:

- An additional pole $P_3 = -\frac{1}{R_o C_I}$ has been created! But since R_o is small (for a buffer) and C_I is small, P_3 is at a very high freq. \rightarrow contributes very little phase @ ω_{ul} , where $|T(j\omega)|=1$.
- A LHP zero now emerges, $z_1 = -\frac{1}{R_o C_c}$.
 This helps stability as discussed before.
 (by contributing (+) phase shift \rightarrow PMM)

Actual Implementation of Buffer-Based Cancellation

$R_{op} = \frac{1}{g_{m3}} = \frac{1}{\sqrt{2\mu_n C_{ox}(W/L)_3 I_{D3}}} \rightarrow$ want this sufficiently small to drive IP_3 up
 \rightarrow increase I_{D3} or increase $(W/L)_3$
 \downarrow
 Problem { more power X more area \rightarrow cost! X

Solution: Nulling Resistor in Series w/ C_c

Doing KCL:

$$P_1 = -\frac{1}{g_{mII} R_2 R_1 C_C} \quad \left. \begin{array}{l} \\ \\ \end{array} \right\} \text{Same as before}$$

$$P_2 = -\frac{g_{mII} C_C}{C_I C_{II} + C_C(C_I + C_{II})} \approx -\frac{g_{mII}}{C_{II}}$$

$$P_3 = -\frac{1}{R_2 C_I} \quad \leftarrow \text{pole due to } R_2$$

$$z = \frac{1}{C_C \left(\frac{1}{g_{mII}} - R_2 \right)} \quad \leftarrow \text{relocated zero (function of } R_2 \text{)}$$

Note: The position of the zero depends upon the value of "nulling resistor" R_2 .

If $R_2 < \frac{1}{g_{mII}}$ then z is in the RHP

If $R_2 > \frac{1}{g_{mII}}$ then " " " LHP!

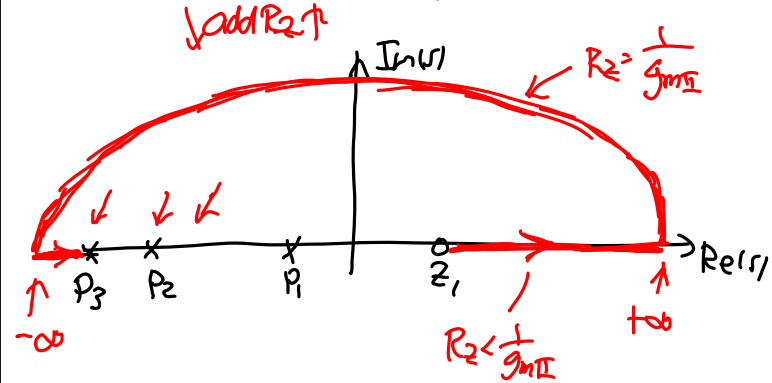
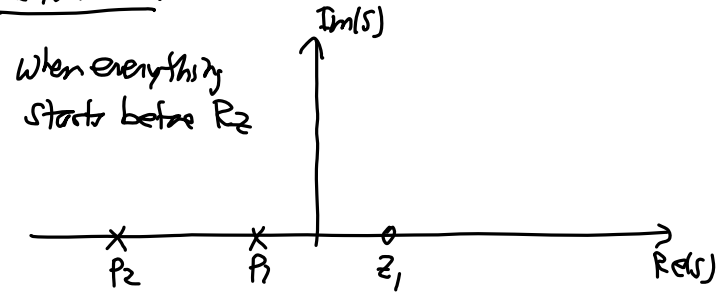
→ This is great! → can convert the RHP zero to a LHP one!

can even stick the zero on top of a pole to eliminate it!

$$H(s) = \dots \frac{(s - z_1)}{(s - p_1)}$$

↑
if $z_1 = p_1$

The Root Locus:



Zero Placement Strategies

① Eliminate z_1 → move it to ∞ :

$$z_1 = \frac{1}{C_C \left(\frac{1}{g_{mII}} - R_2 \right)} \rightarrow \infty \quad \text{when} \quad \boxed{R_2 = \frac{1}{g_{mII}}}$$

→ After doing this: $P_3 \approx -\frac{g_{mII}}{C_I}$
 $P_2 \approx -\frac{g_{mII}}{C_{II}}$

Low C_C .
↓
Usually $C_{II} \gg C_I$
so these poles are far apart...
but be careful

This is good, but we can do better:

② Eliminate p_3 by placing z_1 on top of it:

$$z_1 = p_3 \Rightarrow \frac{1}{C_c \left(\frac{1}{g_{mII}} - R_2 \right)} = - \frac{1}{R_2 C_I}$$

$$R_2 = \frac{1}{g_{mII} \left(1 - \frac{C_I}{C_c} \right)}$$

After this: ① p_3 gone; p_1 & p_2 left

② Now, can place w_{u1} @ $|p_2|$ and really get $PM = 45^\circ \rightarrow$ don't have to worry about phase influence of $p_3!$

But can still do better

③ Eliminate p_2 by placing z_1 on top of it:

$\Rightarrow p_3$ becomes the 'new' p_2 ! (higher freq., higher T_{max})

$$z_1 = p_2 \Rightarrow \frac{1}{C_c \left(\frac{1}{g_{mII}} - R_2 \right)} = - \frac{g_{mII}}{C_I}$$

$$R_2 = \left(\frac{C_c + C_I}{C_c} \right) \left(\frac{1}{g_{mII}} \right) = \frac{1}{g_{mII}} \left(1 + \frac{C_I}{C_c} \right)$$

