

EE243 Advanced Electromagnetic Theory

Lec # 15 Plasmons

Fixed Slides 6, 12

- **(Group and Phase Velocity for Guided Waves)**
- **(Sign Reversals for Propagation Direction Reversal)**
- **Plasmon Fields, B.C. and Dispersion Equation**
- **Physical Interpretation**
- **Typical Values for Silver**
- **ω – β diagram and fields at various points**
- **Assistance from topographical features**

Reading: These PPT notes and Harrington on corrugated surface 4.8.

Midterm Exam

Review Th, Oct 19

- In Class Tuesday October 24th
- Covers material through today (Chapter 7)
- Open Book, Open Notes, Bring Calculator, Paper Provided
- Topics
 - Green's functions free space and use in Theorems and concepts with emphasis on statics
 - Separation of variables in rectangular coordinates using N-1 and N method
 - Time-Harmonic ME, planewaves, boundary conditions, and dispersion

Overview

- (First touch up
 - product of group and phase = c^2 for guided waves
 - Reverse direction reverse sign E_z and H_t).
- Surface Plasmons on metals
 - are TM guided waves
 - are made possible by free electrons oscillating in the propagating field just below their plasma frequency
 - this makes the permittivity both negative and larger in magnitude than the permittivity of air or a coating dielectric.
- Plasmons are described by Maxwell Equations
 - Their dispersion relationship is found by matching B.C.
 - They propagate a distance of up to $30 \mu\text{m}$
 - Near the plasma frequency their wavelength along the surface can be 5X smaller than the free space wavelength
 - They can be used to form high resolution ($\lambda/6$) intensity probes but the probe height is roughly equal to the resolution and thus decreases.

Clarification on Phase and Group Velocity

$$\gamma^2 = \mu\epsilon\omega^2 - k^2$$

$$\omega = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\lambda^2 + k^2}$$

$$v_p = f\lambda = \frac{2\pi f}{2\pi} = \frac{\omega}{k} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\frac{\lambda^2}{k^2} + 1} > c$$

$$v_g = \frac{\partial\omega}{\partial k} = \frac{\frac{1}{2}(2k)}{\sqrt{\mu\epsilon} \sqrt{\lambda^2 + k^2}} = \frac{1}{\sqrt{\mu\epsilon}} \frac{1}{\sqrt{\frac{\lambda^2}{k^2} + 1}} < c$$

$$v_p v_g = c^2$$

- The fact that the product of the phase and group velocity is equal to the square of the speed of light holds for any (propagating) mode in a lossless system due to its eigenfunction and eigenvalue γ^2 .

Sign Reversal with Direction Reversal

$$\nabla_t \times \bar{E}_{t\lambda} = i\omega\mu\bar{H}_{z\lambda}$$

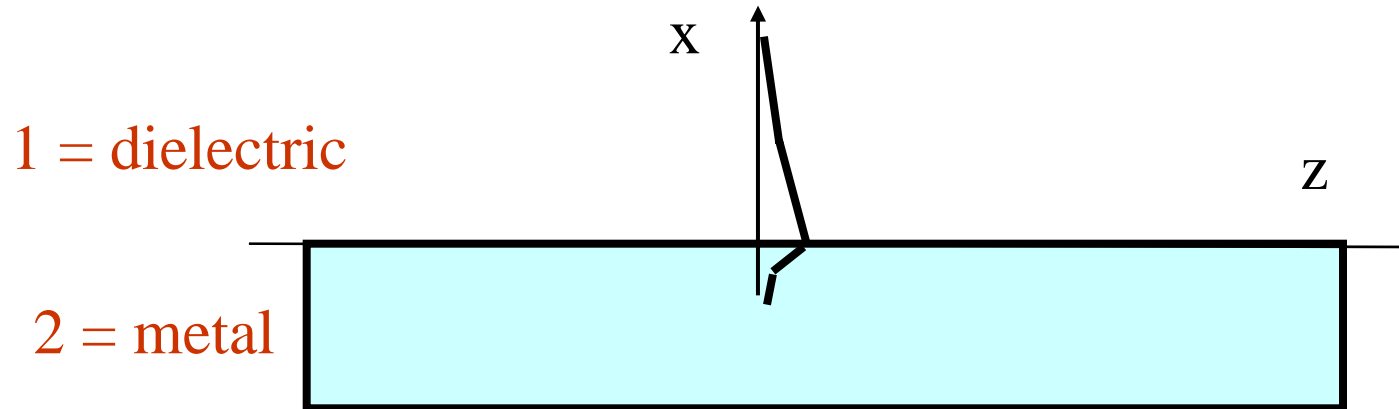
$$\nabla_t \times \bar{H}_{t\lambda} = -i\omega\varepsilon\bar{E}_{z\lambda}$$

$$\nabla_t \times \bar{E}_{z\lambda} + i\beta_\lambda \hat{z} \times \bar{E}_{t\lambda} = i\omega\mu\bar{H}_{t\lambda}$$

$$\nabla_t \times \bar{H}_{z\lambda} + i\beta_\lambda \hat{z} \times \bar{H}_{t\lambda} = -i\omega\varepsilon\bar{E}_{t\lambda}$$

- Take longitudinal and transverse components of curl equations with z phase variation.
- With propagation direction reversal the sign of β_λ changes
- But if sign of E_z and H_t are also reversed the original equations result
- Thus the reverse guided wave solution can be nothing more than the forward guided wave solution with these changes

Are There Waves on Material Surfaces?



$$H_{1y} = H_1 \hat{y} e^{-\nu_1 x} e^{ik_z z}$$

$$H_{2y} = H_2 \hat{y} e^{+\nu_2 x} e^{ik_z z}$$

$$\nu_1 = \sqrt{k_z^2 - \omega^2 \mu_0 \epsilon_1}$$

$$\nu_2 = \sqrt{k_z^2 - \omega^2 \mu_0 \epsilon_2}$$

- Consider TM w/r z case with H_y given and same z phase variation
- Will have H_y , E_z and E_x (but $E_y = H_x = H_z = 0$)

Boundary Conditions

$$E_{1z} = \frac{-v_1}{-i\omega\epsilon_1} H_{1y}$$

$$E_{2z} = \frac{+v_2}{-i\omega\epsilon_2} H_{2y}$$

$$\frac{-iv_1}{\epsilon_1} = \frac{+iv_2}{\epsilon_2}$$

$$\frac{E_{1z}}{H_{1y}} = \frac{E_{2z}}{H_{2y}}$$

$$-Z_{+x} = \frac{-v_1}{-i\omega\epsilon_1} = \frac{+v_2}{-i\omega\epsilon_2} = Z_{-x}$$

- H_y continuous (or) D normal continuous gives $H_{10} = H_{20}$.
- E_z continuous gives final constraint to find k_z .
- This constraint is the same as setting the impedance looking upward equal to the negative of the impedance looking downward.
- Impedance looking upward is capacitive (neg imy).
- Impedance looking downward thus need to be inductive.

Solving for Surface Wave Conditions

$$\frac{v_1}{\epsilon_1} = \frac{-v_2}{\epsilon_2}$$

$$v_1 = \sqrt{k_z^2 - \omega^2 \mu_0 \epsilon_1}$$

$$v_2 = \sqrt{k_z^2 - \omega^2 \mu_0 \epsilon_2}$$

$$k_z = k_1 \sqrt{\frac{\epsilon_2}{(\epsilon_2 + \epsilon_1)}}$$

$$v_1 = k_1 \sqrt{\frac{-\epsilon_1}{(\epsilon_2 + \epsilon_1)}}$$

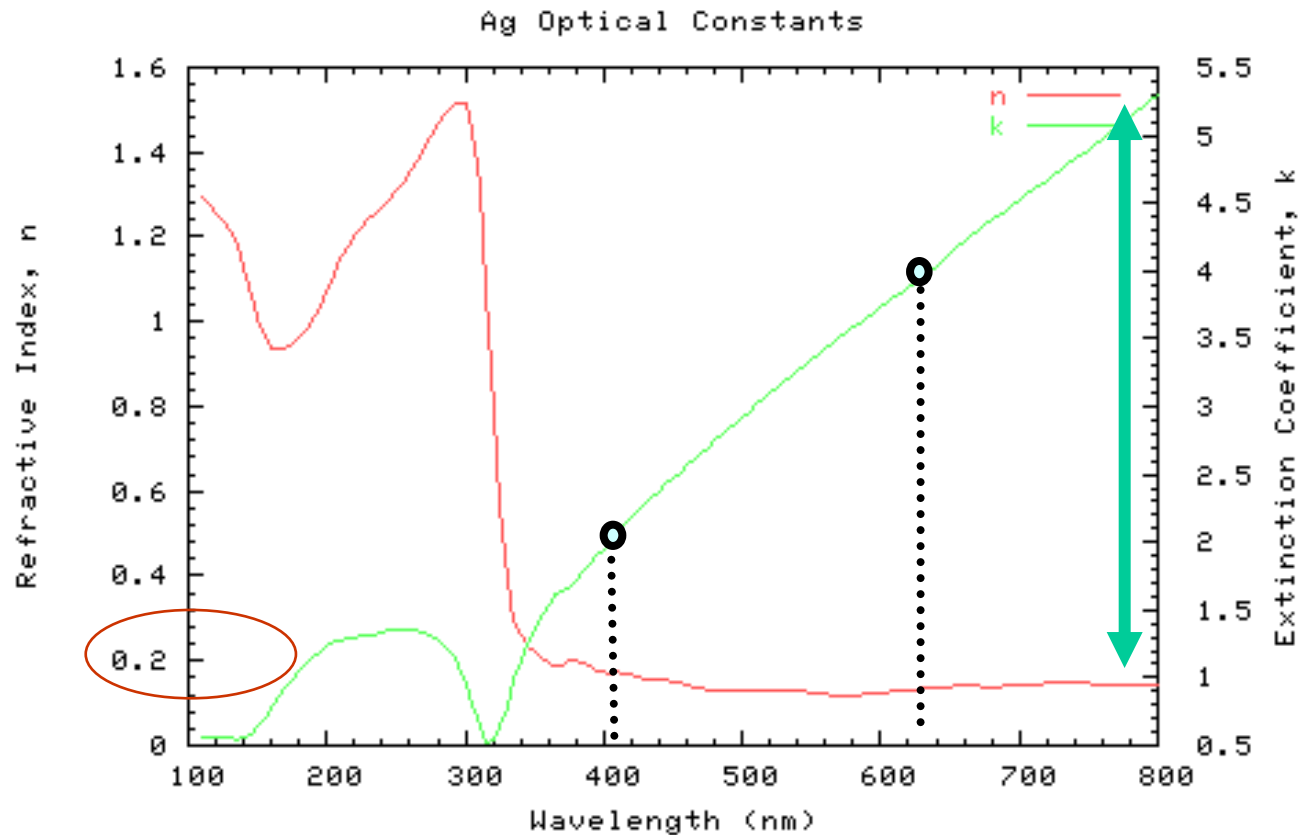
- Constraint
- Substitute definition of v_1 and v_2 to solve for k_z .
- Substitute solution for k_z to find other properties
 - v_1 and v_2 (localization in x)
 - Resolution in z with large k_z
 - Probe height in x

Physical Interpretation

$$v_1 = k_1 \sqrt{\frac{-\varepsilon_1}{(\varepsilon_2 + \varepsilon_1)}} \quad k_z = k_1 \sqrt{\frac{\varepsilon_2}{(\varepsilon_2 + \varepsilon_1)}} \quad \frac{v_2}{v_1} = \frac{-\varepsilon_2}{\varepsilon_1}$$

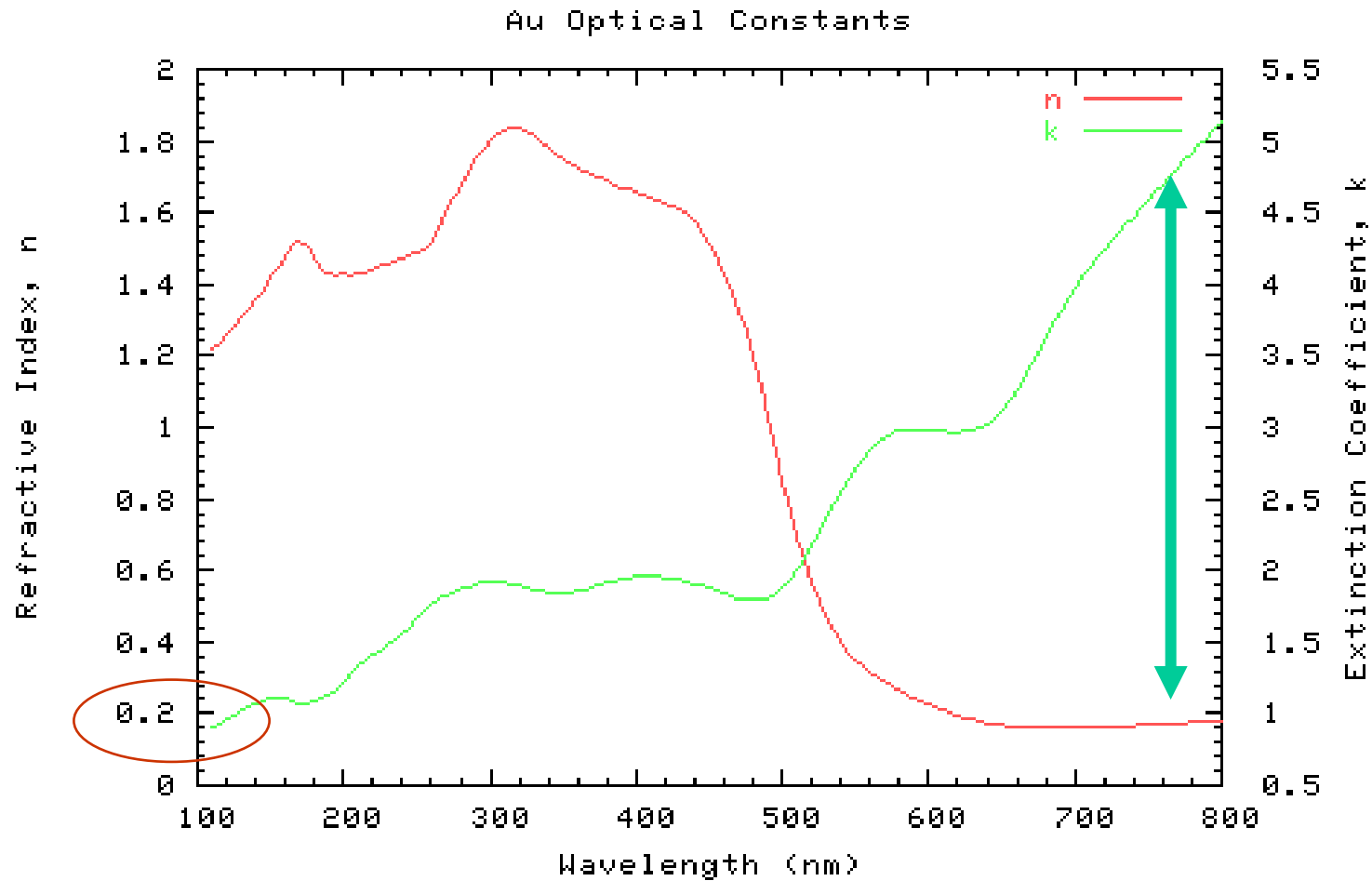
- Want v_1 real so exponential decay away from boundary
- Since ε_1 is real and positive this means that the denominator must go negative.
- This occurs when $\varepsilon_2 < -\varepsilon_1$. (plasmon condition for a single surface to support a guided wave)
- In metals just below their plasma resonance frequency a negative permittivity can occur (Ag, Au, Cu, Al)
- In general k_z is complex and the waves die out in a few microns

Refractive Index of Silver



- <http://www.microe.rit.edu/research/lithography/utilities.htm>

Refractive Index of Gold



- <http://www.microe.rit.edu/research/lithography/utilities.htm>

Physical Estimates for Air and Silver

$$\frac{v_2}{v_1} = \frac{-\varepsilon_2}{\varepsilon_1} \approx \frac{K^2}{1}$$

$$n_2 = n_r + iK$$

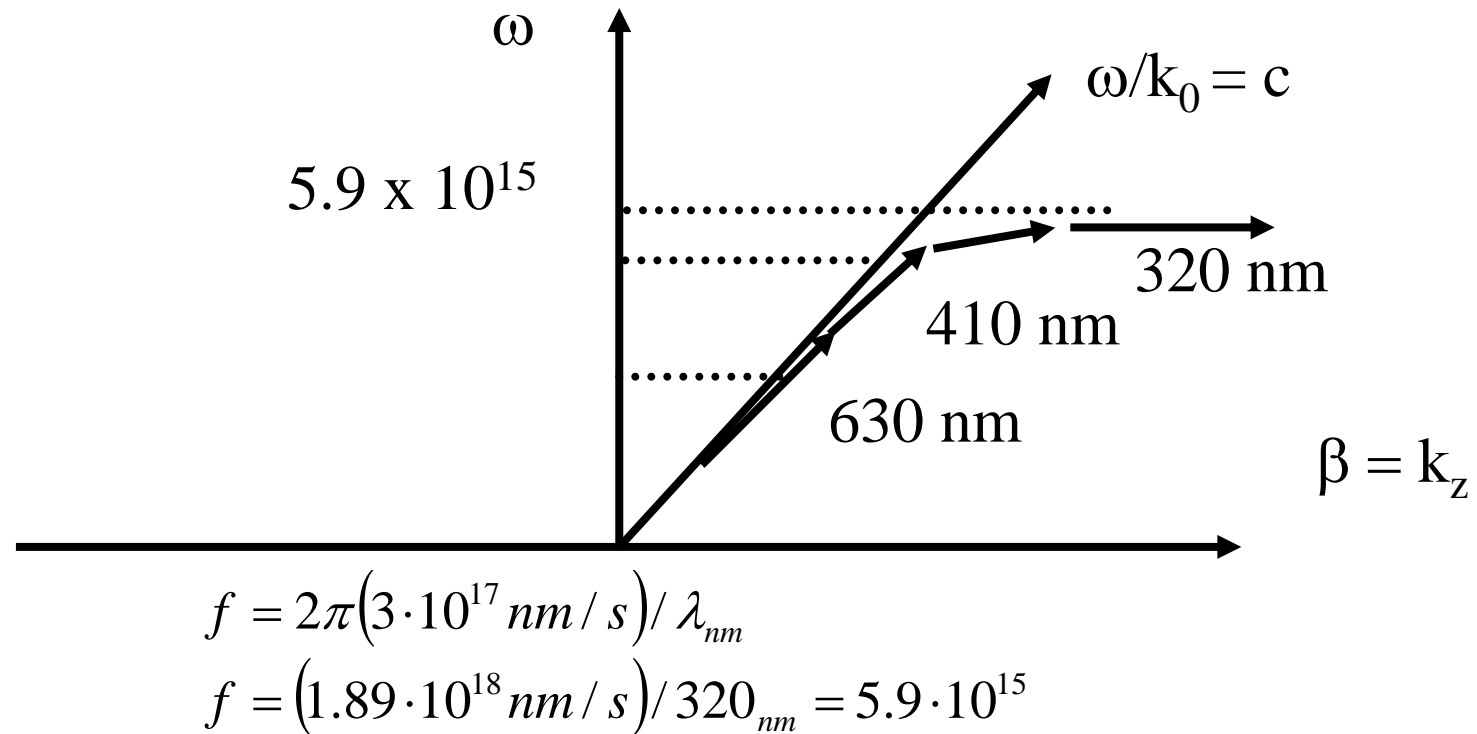
$$n_1 = 1.0$$

$$k_z = k_0 \sqrt{\frac{\varepsilon_2}{(\varepsilon_2 + \varepsilon_1)}} \approx k_0 \left[\frac{K}{\sqrt{K^2 - 1}} + \frac{in_r}{2(K^2 - 1)^{3/2}} \right]$$

$$l_z/\lambda = \frac{2(K^2 - 1)^{3/2}}{n_r}$$

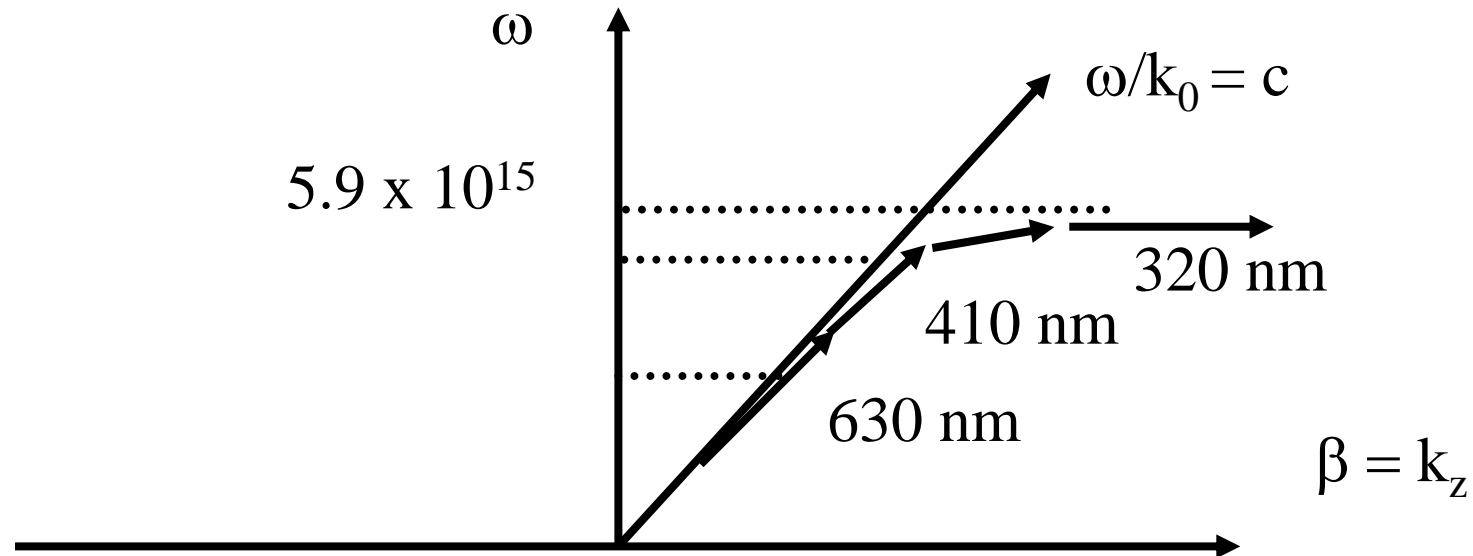
- Approximately for silver in the visible
 - $n_r = 0.2$, and $K = 2$ (410 nm) and 4 (630 nm)
 - the velocity is $0.87c$ and $0.97c$,
 - $\lambda_{\text{surface}} = (2\pi/k_z)$ can be smaller than λ_{air} by over 3X \Rightarrow 3X resolution improve
 - the 1/e length is 4λ (1.8 μm) and 295λ (2.9 μm)
 - The 1/e penetration into the metal is about $\lambda_{\text{surface}}/5$
 - The 1/e penetration into air is 4 to 16 times larger

ω - β Diagram for Plasmon



- The plasmons start as frequency is increased
 - close to the speed of light line,
 - become slightly slower, and
 - turn into a very slow wave (horizontal line) at the plasma frequency.

Plasmon Characteristics



- 630 nm: $0.97c$, $1/e$ 46.3λ ($29 \mu\text{m}$), 16X air
- 410 nm: $0.86c$, $1/e$ 4.1λ ($1.8 \mu\text{m}$), 4X air
- 320 nm:
 - $k_z = 10k_0 \Rightarrow 10\text{X}$ resolution
 - But evanescent near fields are 10X faster decay

Near Field Probe Limits

$$R = \frac{\lambda_{surface}}{2\pi} = \frac{1}{k_z} = \frac{1}{k_1 \sqrt{\frac{\epsilon_2}{(\epsilon_1 + \epsilon_2)}}}$$

$$h_{1/e} = \frac{1}{\nu_1} = \frac{1}{k_1 \sqrt{\frac{-\epsilon_1}{(\epsilon_1 + \epsilon_2)}}}$$

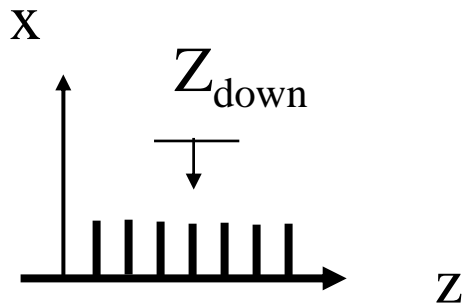
$$h_{1/e} = \sqrt{\frac{-\epsilon_1}{\epsilon_2}} R \approx R$$

- Work near the plasmon frequency to make k_z large and resolution smaller than normal by 3 to 10X.
- But the 1/e height reduces by the same amount that the resolution is reduced

Surface Topography Can Aid Guided Waves

Harrington 4.8

Corrugated Surface



Impedance looking down into the corrugations is inductive.

$$Z_{down} = -i \sqrt{\frac{\mu_1}{\epsilon_1}} \tan k_0 d = -i377 \tan k_0 d$$

$$Z_{up} = \frac{-v_1}{i\omega\epsilon_1}$$

$$k_z = k_0 \sqrt{1 + \tan^2 k_0 d}$$

- Impedance looking into slot is that of a parallel plate waveguide terminated in a short.
- Slot must be narrow compared to a wavelength
- Depth must be $>5\%$ of wavelength to contribute.
- For plasmons
 - effects might add
 - which wavelength should be used