Micromachined Inertial Sensors

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Invited Paper

This paper presents a review of silicon micromachined accelerometers and gyroscopes. Following a brief introduction to their operating principles and specifications, various device structures, fabrication technologies, device designs, packaging, and interface electronics issues, along with the present status in the commercialization of micromachined inertial sensors, are discussed. Inertial sensors have seen a steady improvement in their performance, and today, microaccelerometers can resolve accelerations in the micro-g range, while the performance of gyroscopes has improved by a factor of 10^x every two years during the past eight years. This impressive drive to higher performance, lower cost, greater functionality, higher levels of integration, and higher volume will continue as new fabrication, circuit, and packaging techniques are developed to meet the ever increasing demand for inertial sensors.

Keywords—Accelerometer, gyroscope, inertial sensors, microfabrication technologies, micromachined sensors, micromachining, rate sensor, silicon sensors.

I. INTRODUCTION

Micromachined inertial sensors, consisting of accelerometers and gyroscopes, are one of the most important types of silicon-based sensors. Microaccelerometers alone have the second largest sales volume after pressure sensors, and it is believed that gyroscopes will soon be mass produced at similar volumes. The large volume demand for accelerometers is due to their automotive applications, where they are used to activate safety systems, including air bags, to implement vehicle stability systems and electronic suspension. However, the application of accelerometers covers a much broader spectrum where their small size and low cost have even a larger impact. They are used in biomedical applications for activity monitoring; in numerous consumer applications, such as active stabilization of picture in camcorders, head-mounted displays and virtual reality, three-dimensional mouse, and sport equipment; in industrial applications such as robotics and machine and vibration monitoring; in many other applications, such as tracking and monitoring mechanical shock and vibration during transportation and handling of a variety of equipment and goods; and in several military applications, including impact and void detection and safing and arming in missiles and other ordnance. High-sensitivity accelerometers are crucial components in self-contained navigation and guidance systems, seismometry for oil exploration and earthquake prediction, and microgravity measurements and platform stabilization in space. The impact of low-cost, small, high-performance, micromachined accelerometers in these applications is not just limited to reducing overall size, cost, and weight. It opens up new market opportunities such as personal navigators for consumer applications, or it enhances the overall accuracy and performance of the systems by making formation of large arrays of devices feasible.

Micromachined gyroscopes for measuring rate or angle of rotation have also attracted a lot of attention during the past few years for several applications. They can be used either as a low-cost miniature companion with micromachined accelerometers to provide heading information for inertial navigation purposes or in other areas [1], including automotive applications for ride stabilization and rollover detection; some consumer electronic applications, such as video-camera stabilization, virtual reality, and inertial mouse for computers; robotics applications; and a wide range of military applications. Conventional rotating wheel as well as precision fiber-optic and ring laser gyroscopes are all too expensive and too large for use in most emerging applications. Micromachining can shrink the sensor size by orders of magnitude, reduce the fabrication cost significantly, and allow the electronics to be integrated on the same silicon chip.

This paper will review both micromachined accelerometers and gyroscopes by providing an introduction to their basic operation, reviewing different types of devices reported in the literature and their performance, discussing important issues in their design and operation, and presenting a summary and conclusion along with a discussion on the future trends in this important category of micromachined
sensors. Due to space constraints, interface electronics and packaging issues for inertial sensors are only briefly described. Also, other topics such as temperature compensation, calibration, and self-testing are not discussed in this paper.

II. MICROMACHINED ACCELEROMETERS

A. Structure, Operation, and Specifications

An accelerometer generally consists of a proof mass suspended by compliant beams anchored to a fixed frame. The proof mass has a mass of $M$, the suspension beams have an effective spring constant of $K$, and there is a damping factor ($D$) affecting the dynamic movement of the mass. The accelerometer can be modeled by a second-order mass-damper-spring system, as shown in Fig. 1. External acceleration displaces the support frame relative to the proof mass, which in turn changes the internal stress in the suspension spring. Both this relative displacement and the suspension-beam stress can be used as a measure of the external acceleration.

By using Newton’s second law and the accelerometer model, the mechanical transfer function can be obtained

$$H(s) = \frac{x(s)}{a(s)} = \frac{1}{s^2 + \frac{D}{M} s + \frac{K}{M}} = \frac{1}{s^2 + \frac{\omega_r}{Q} s + \omega_r^2}$$

where $a$ is the external acceleration, $x$ is the proof mass displacement, $\omega_r = \sqrt{K/M}$ is the natural resonance frequency, and $Q = \sqrt{KM/D}$ is the quality factor. The static sensitivity of the accelerometer is shown to be

$$\frac{x_{\text{static}}}{a} = \frac{M}{K} = \frac{1}{\omega_r^2}.$$  

As evident, the resonance frequency of the structure can be increased by increasing the spring constant and decreasing the proof mass, while the quality factor of the device can be increased by reducing damping and by increasing proof mass and spring constant. Last, the static response of the device can be improved by reducing its resonant frequency.

The primary mechanical noise source for the device is due to Brownian motion of the gas molecules surrounding the proof mass and the Brownian motion of the proof-mass suspension or anchors. The total noise equivalent acceleration (TNEA) $[\text{m/s}^2/\sqrt{\text{Hz}}]$ is [2]

$$\text{TNEA} = \frac{\sqrt{4K_B T D}}{M} = \sqrt{\frac{4K_B T \omega_r}{Q M}}$$

where $K_B$ is the Boltzmann constant and $T$ is the temperature in kelvin. Equation (3) clearly shows that to reduce mechanical noise, the quality factor and proof mass have to be increased.

In the most general case, the proof-mass motion can have six degrees of freedom. But typically in a unidirectional accelerometer, the geometrical design of the suspension is such that one of these is dominant [3] and the device has low off-axis sensitivity. The cantilever support has been one of the early popular suspension support designs [4], [5] due to its simplicity, lower spring constant, and internal stress relief of the beams. However, this configuration results in a larger off-axis sensitivity unless the device is fully symmetric. Also, symmetric full-bridge supports result in a very low off-axis sensitivity [6]–[8], and by using a crab-leg or folded-beam configuration [3] in a full-bridge support, the residual stress of the beams can also be relieved. The general design of accelerometers can be performed using the above equations, as well as mechanical relations describing the spring constant [9], [10] and damping factor as a function of device geometry and ambient pressure [11], [12]. Further, the device first-order design optimization can be obtained using the same equations [4], [5], while the final accelerometer design can be simulated and optimized using commercially available finite element method or dedicated microelectromechanical systems (MEMS) software packages [13].

Accelerometers are typically specified by their sensitivity, maximum operation range, frequency response, resolution, full-scale nonlinearity, offset, off-axis sensitivity, and shock survival. Since micromachined accelerometers are used in a wide range of applications, their required specifications are also application dependent and cover a rather broad spectrum. For instance, for microgravity measurements devices with a range of operation greater than $\pm 0.1 \text{ g}$, a resolution of less than $1 \mu\text{g}$ in a frequency range of zero frequency to 1 Hz are desired, while in ballistic and impact sensing applications, a range of over 10,000 g with a resolution of less than 1 g in a 50 kHz bandwidth is required. Table 1 summarizes typical performance parameters of accelerometers with medium resolution for automotive applications and high performance for inertial navigation applications.

B. Device Types

A variety of transduction mechanisms have been used in microaccelerometers. Some of the more relevant and useful approaches will be reviewed here.

1) Piezoresistive Devices: The first micromachined [4], and one of the first commercialized, microaccelerometers [14] were piezoresistive. These accelerometers incorporate silicon piezoresistors in their suspension beam. As the
support frame moves relative to the proof mass, the suspension beams will elongate or shorten, which changes their stress profile and hence the resistivity of their embedded piezoresistors. These piezoresistors are generally placed at the edge of the support rim and proof mass, where the stress variation is maximum. Therefore, a resistive half-bridge or full bridge can be formed by employing two or four piezoresistors.

The main advantage of piezoresistive accelerometers is the simplicity of their structure and fabrication process, as well as their readout circuitry, since the resistive bridge generates a low output-impedance voltage. However, piezoresistive accelerometers have larger temperature sensitivity, and smaller overall sensitivity compared to capacitive devices, and hence a larger proof mass is preferred for them. The early development of bulk-micromachining technology, and the experience gained by several companies in the development and commercialization of piezoresistive pressure sensors, helped the initial development of piezoresistive microaccelerometers using bulk-micromachining and wafer-bonding technology [4], [14]–[18]. The device reported in [4] uses a silicon middle wafer to form the proof mass and the beams, while two glass wafer caps are bonded on top and bottom to cover the structure and provide shock stop and damping. Other devices [14], [15] use a lower bonded glass base and a silicon overhang on the top for shock stop, which is formed by silicon fusion bonding and etch back. Recently, monolithic implementation of a piezoresistive microaccelerometer with its interface complementary metal–oxide–semiconductor (CMOS) circuitry was demonstrated [17], [18], which is based on using a slight modification of a standard CMOS process to implement the accelerometer readout and temperature compensation circuitry and bulk etching of the silicon wafer from backside to form the device structure. Self-testing of the piezoresistive accelerometers is possible by using thermal actuation [15] or electrostatic forces [16]. The sensitivity of all these devices is typically around 1–2 mV/g in a 20–50-g range with an uncompensated temperature coefficient of sensitivity (TCS) of $<0.2\%/\text{C}$.

2) Capacitive Devices: In the presence of external acceleration, the support frame of an accelerometer moves from its rest position, thus changing the capacitance between the proof mass and a fixed conductive electrode separated from it with a narrow gap. This capacitance can be measured using electronic circuitry. Silicon capacitive accelerometers have several advantages that make them very attractive for numerous applications ranging from low-cost, large-volume automotive accelerometers [19], [20] to high-precision inertial-grade microgravity devices [21]–[26]. They have high sensitivity, good dc response and noise performance, low drift, low temperature sensitivity, low-power dissipation, and a simple structure. However, capacitive accelerometers can be susceptible to electromagnetic interference (EMI), as their sense node has high impedance. This issue can be addressed by proper packaging and by shielding the accelerometer and its interface circuit.

Some of the most widely used structures for capacitive accelerometers are vertical and lateral structures, as shown in Fig. 2. Many capacitive accelerometers utilize the vertical structure, where the proof mass is separated by a narrow air gap from a fixed plate, forming a parallel plate sense capacitance [7], [8], [20]–[27]. In these devices, the proof mass moves in the direction perpendicular to its plane ($z$-axis) and changes the air gap. In a lateral accelerometer, a number of moving sense fingers are attached to the proof mass, and the sense capacitance is formed between these and the fixed fingers parallel to them. The sense direction in lateral accelerometers is in the proof-mass plane ($x$–$y$ directions) [19], [29]–[31]. Some designs use a “see-saw” structure, shown in Fig. 3, where a proof mass is suspended by torsional beams so that one side is heavier than the other side and in response to acceleration in the $z$-axis, the proof mass moves out of its plane [32]–[34]. The advantages of this structure over conventional parallel-plate $z$-axis devices are built-in overrange protection, larger sensitivity, and higher pull-in voltage [34].

The open-loop sensitivity of a capacitive accelerometer is proportional to the proof-mass size and capacitance overlap area and inversely proportional to the spring constant and air gap squared. Early micromachined capacitive accelerometers [7], [21]–[23] utilized bulk silicon micromachining and wafer bonding to achieve a thick, large proof mass and high sensitivity. One of the first reported devices [21] used a silicon middle wafer anodically bonded to two glass wafers on top and bottom to form a $z$-axis accelerometer. The device had two differential sense capacitors, with the proof mass forming the middle electrode and metal on the glass wafers forming the top/bottom fixed electrodes. The air gap was formed by recessing the silicon or glass wafers. This device with a proof-mass size of 4.6 mg and air gap of 2 μm provided μg-level performance. The second generation of this device [22] had a resolution of better than 1 μg/√Hz in a bandwidth of zero frequency to 100 Hz, with a temperature coefficient of offset (TCO) of 30 μg/°C and TCS of 150 ppm/°C. To

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Table 1 Typical Specifications of Accelerometers for Automotive and Inertial Navigation Applications

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Automotive</th>
<th>Navigation</th>
</tr>
</thead>
<tbody>
<tr>
<td>Range</td>
<td>±50 g (airbag)</td>
<td>±1 g</td>
</tr>
<tr>
<td></td>
<td>±2 g (vehicle stability system)</td>
<td></td>
</tr>
<tr>
<td>Frequency Range</td>
<td>DC–400 Hz</td>
<td>DC–100 Hz</td>
</tr>
<tr>
<td>Resolution</td>
<td>&lt;100 mg (airbag)</td>
<td>&lt;4 μg</td>
</tr>
<tr>
<td></td>
<td>&lt;10 mg (vehicle stability system)</td>
<td></td>
</tr>
<tr>
<td>Off-axis Sensitivity</td>
<td>&lt;5%</td>
<td>&lt;0.1%</td>
</tr>
<tr>
<td>Nonlinearity</td>
<td>&lt;2%</td>
<td>&lt;0.1%</td>
</tr>
<tr>
<td>Max. Shock in 1 msec</td>
<td>&gt;2000 g</td>
<td>&gt;10 g</td>
</tr>
<tr>
<td>Temperature Range</td>
<td>-40°C to 85°C</td>
<td>-40°C to 80°C</td>
</tr>
<tr>
<td>TC of Offset</td>
<td>&lt;60 mg/°C</td>
<td>&lt;50 μg/°C</td>
</tr>
<tr>
<td>TC of Sensitivity</td>
<td>&lt;900 ppm/°C</td>
<td>±50 ppm/°C</td>
</tr>
</tbody>
</table>
reduce its temperature sensitivity and long-term drift, the later generation of this device was fabricated using three silicon wafers [24], [35]. Another significant early design with $\mu g$ performance was fabricated using glass–silicon bonding and bulk micromachining and utilized a closed-loop $\Sigma \Delta$ readout and control circuit to achieve a 120 dB dynamic range [23].

A number of capacitive microaccelerometers with medium resolution have been fabricated using the bulk silicon dissolved wafer process [27], [33], [34]. Among these, the device reported in [33] is in high-volume production by Ford. The $z$-axis accelerometer presented in [27] incorporates damping holes in the proof mass to control damping and uses a second silicon wafer, bonded on top, to provide overrange protection. The torsional accelerometer shown in [34, Fig. 3] is fabricated using a three-mask dissolved wafer process. It uses multiple interdigitated fingers and a varying overlap area method of capacitance change to achieve higher linearity, larger pull-in voltage, and lower damping. The device consists of a 12-$\mu$m-thick boron-doped silicon inertial mass suspended 7.5 $\mu$m above the glass substrate by two narrow, high-aspect-ratio, $12 \times 3 \mu$m torsion beams. A large number of 300-$\mu$m-long capacitive sense fingers are attached to the end of the proof mass and are separated by a 2 $\mu$m air gap from fixed fingers anchored to the substrate. The accelerometer provides a sensitivity of 20 fF/g over a range of $\pm 4$ g and a bandwidth of 30 Hz in atmosphere.

Surface-micromachined accelerometers [19], [20], [29], [30], [36]–[40] offer the opportunity to integrate the sensor and interface circuitry on a single chip. These devices utilize deposited polysilicon layers to form the sense element and are well suited for both vertical and lateral capacitive accelerometers. Integration of the interface circuitry with the sensor enables detection of very small capacitance variation ($<1\mu$F [38]), while the accelerometer and all its interface electronics can be implemented in a small area. Fig. 4 shows a scanning electron microscope (SEM) view of the ADXL05, a polysilicon accelerometer developed by Analog Devices, which has 0.5 mg/$\sqrt{\text{Hz}}$ noise floor with a range of $\pm 5$ g and shock survivability of 1000 g [39]. Also, by employing a vertical and two lateral accelerometers, an integrated three-axis accelerometer system has been designed by researchers at Berkeley and fabricated through Sandia National Laboratory’s process [38]. The same group has also developed a three-axis accelerometer with a single sense element, as shown in [40], (Fig. 5). Typically, surface microaccelerometers achieve several 100 $\mu$g resolution in $\sim 100 \text{ Hz}$ bandwidth [37]–[40]. All of these devices have small proof mass and hence high mechanical noise (unless the device is vacuum packaged).

While bulk-micromachined devices can attain higher resolution due to their large proof mass, they generally require wafer bonding. These devices, if not formed using only silicon wafers, could have large temperature coefficients. Moreover, forming damping holes in their thick structural layers is not easy, and they typically require packaging at a specified pressure to control the damping [35]. To address these shortcomings, an all-silicon, fully symmetric,
high-precision accelerometer has been developed [26], as shown in Fig. 6. This device uses a combined surface and bulk micromachining process to obtain a large proof mass, controllable/small damping, and a small air gap for large capacitance variation—all by using a single silicon wafer. Sense electrodes are created by depositing polysilicon on the wafer. These electrodes, while thin, are made very stiff by embedding thick vertical stiffeners (using a trench refill technology) in them so that force rebalancing of the proof mass becomes possible. There are eight suspension beams, which are symmetric with respect to the proof-mass centerline and result in low cross-axis sensitivity. An SEM view is shown in Fig. 6(b). This device is operated closed loop using an oversampled \( \Sigma/\Delta \) modulator. It achieves a high sensitivity of 2 pF/g with a full-bridge support, a low noise level, and low temperature sensitivity, which enable it to attain \( \mu \)g and sub-\( \mu \)g performance.

3) Tunneling Devices: Some high-resolution physical sensors, including microaccelerometers [41]–[47], use a constant tunneling current between one tunneling tip (attached to a movable microstructure) and its counterelectrode to sense displacement. Fig. 7 [44] shows the general operating principal of a micromachined tunneling accelerometer. As the tip is brought sufficiently close to its counter-electrode (within a few angstroms) using electrostatic force generated by the bottom deflection electrode, a tunneling current (\( I_{\text{tunnel}} \)) is established and remains constant if the tunneling voltage (\( V_{\text{tunnel}} \)) and distance between the tip and counterelectrode are unchanged. Once the proof mass is displaced due to acceleration, the readout circuit responds to the change of current and adjusts the bottom deflection voltage \( V_d \) to move the proof mass back to its original position, thus maintaining a constant tunneling current. Acceleration can be measured by reading out the bottom deflection voltage in this closed-loop system. Tunneling accelerometers can achieve very high sensitivity with a small size since the tunneling current is highly sensitive to displacement, typically changing by a factor of two for each angstrom of displacement [48]. However, these devices have larger low-frequency noise levels [49].

Micromachined tunneling accelerometers were first introduced by researchers with the Jet Propulsion Laboratory (JPL), Pasadena, CA [41]–[43]. One of their more recent devices provides a bandwidth of several kilohertz, with a noise floor of less than \( 10^{-7} \) g/\( \sqrt{\text{Hz}} \) in the 10–200-Hz frequency range [43]. However, a high supply voltage (tens to hundreds of volts) is required for these devices, thus limiting their application. Using the dissolved wafer process, a low-voltage micromachined tunneling accelerometer has been developed [44], [46], [47]. This accelerometer has a noise spectral density of the sensor-circuit module of \( 4 \) mg/\( \sqrt{\text{Hz}} \) (at 0.5 Hz) and \( 0.1 \) mg/\( \sqrt{\text{Hz}} \) (at 2.5 kHz) and has a minimum detectable acceleration of 8 mg in a 2.5 kHz bandwidth [47]. In continuous operation over 720 h, the accelerometer shows an offset and sensitivity variation of <0.5% [47]. Since these sensors are operated in a closed-loop mode, and since the operating current levels are typically very small, the drift of the tunneling barrier usually does not have a large impact on the performance of the sensor [46].

4) Resonant Devices: The main advantage of resonant sensors is their direct digital output. The first resonant accelerometers were fabricated using quartz micromachining [50], [51]. Silicon resonant accelerometers are generally based on transferring the proof-mass inertial force to axial force on the resonant beams and hence shifting their frequency [52]. To cancel device thermal mismatches and nonlinearities, a differential matched resonator configuration can be used [53]. Recently, two high-sensitivity resonant accelerometers have been reported [54], [55]. The devices use wafer-thick proof mass and achieve high resolution (700 Hz/g with 524 kHz center frequency [55]) and very good stability (2 \( \mu \)g in more than several days [54]). However, these devices typically have small bandwidth (less than a few hertz). Also recently, surface-micromachined resonant accelerometers are developed [56], [57]. The resonator reported in [57] consists of parallel beams, and its operation is based on rigidity change of the resonator due to its cross-sectional shape change, which is induced by the external acceleration. In this manner, the device is expected to achieve high sensitivity of 10%/g [57].

5) Thermal Devices: Another class of accelerometers is based on thermal transduction. One of the first thermal accelerometers used the principle that the temperature flux from a heater to a heat-sink plate is inversely proportional to their separation [58]. Hence, by measuring the temperature using thermopiles, the change in separation between the plates (which is representative of acceleration) can be measured. Devices with a moving thermopile array...
and fixed heater, and vice versa, can be fabricated [58]. Recently, a novel thermal accelerometer was reported that does not have any moving mechanical parts. Its operation is based on free-convection heat transfer of a small hot air bubble in a sealed chamber [59]. The device consists of a thermally isolated heater that forms a hot air bubble. The heat distribution of this bubble changes in the presence of an acceleration and becomes asymmetric with respect to the heater. This heat profile can be sensed by two symmetrically placed temperature sensors and is a measure of the acceleration. The initial prototypes achieved a 0.6 mg sensitivity, and future devices are expected to achieve sub-μg performance [59].

6) Other Devices: In addition to the aforementioned device types, accelerometers also use many other principles, including optical [60], [61], electromagnetic [62], and piezoelectric [63], [64]. The motivation for the development of optical accelerometers has been combining optics and silicon micromachining to exploit advantages of both, as well as achieving miniature devices with very high EMI noise immunity [60] or good linearity [61]. The electromagnetic accelerometer reported in [62] utilizes two coils, one on top of the proof mass and the other separated by an air gap at the bottom, where the proofmass displacement changes the mutual inductance of the two coils. By using a simple readout circuit with a 2.5
mW power dissipation, a very linear response over a 50 g range was achieved [62]. Piezoelectric materials, mainly ZnO, have also been used in accelerometers to directly convert the force affecting the proof mass to an electrical signal [63], [64]. The piezoelectric charge generated by acceleration can be directly coupled to the gate of an MOS transistor and amplified [63]. One of the problems with piezoelectric materials is their leakage that deteriorates the dc response of the device. By isolating the piezoelectric film from the leakage paths, a more flat response at near dc frequencies can be obtained [63].

Also, many researchers [65]–[67] have fabricated arrays of accelerometer switches for use as a threshold or digital accelerometer. The main advantage of these devices is a digital output and simple and low-power interface circuitry [67].

In addition to bulk and surface micromachining, electroforming and mixed processes [25], [26], [28], [69] have been used to improve performance further. Electroforming can produce structures with large proof mass. For instance, a lateral accelerometer fabricated by LIGA produces a 1 μg/√Hz resolution in air and a temperature-compensated TCO of 150 ppm/°C [28], [68]. The device presented in [25] uses a combination of thick epigrowth, SIMOX formation, and wafer bonding to attain high precision. The 70 μm proof mass is formed by two bonded epilayers, while the thin uniform air gap is created by etching the sacrificial buried oxide layer. A family of trench refill technologies for MEMS has been developed [69] and applied to fabricate integrated three-axis accelerometers with a medium resolution. The small mass size in surface-micromachined devices [29], [30], [36]–[40] can be somewhat circumvented by using silicon-on-insulator (SOI) wafers, where the structure height is defined by the top silicon layer, or by depositing a thick polysilicon layer using an epitaxial process [70].

**C. Packaging and Interface Circuits**

The accelerometer package has to protect the sensor structure without inducing significant stress or drift. Mounting errors or misalignment directly affect the sense direction of the device and its overall performance. Also, the package should not adversely affect the sensor frequency response or temperature sensitivity. Both metal can [22], [71] and multilayer ceramic hermetic packages [72] have been utilized to house the sensor and its interface circuit. The overall packaging cost can be reduced, and the performance can be improved, by packaging the sensor/circuit at the wafer level using capping glass or silicon wafers that are bonded to the device wafer, then using plastic injection molding for the final external package, as demonstrated by Motorola [73] and Ford [74].

Sensor readout and signal-processing circuit design and development is also an important aspect of micromachined accelerometers. In a piezoresistive accelerometer, the sensor piezoresistors form a half-bridge or full-bridge, and the output voltage of the bridge can be directly amplified by the interface circuit. Tunneling sensors also require simple readout circuitry that convert the tunneling current to voltage, buffer and amplify this voltage, and feed it back to the sensor [47]. Resonant sensors are operated in closed loop, and their output is their frequency, which can be measured by a digital counter. Their interface circuit consists of sense and drive circuitry. A sustaining amplifier in the loop is needed to compensate for the losses and maintain the resonance [76]. Capacitive readout circuits mostly use capacitance-to-frequency converters (oscillators) [77]–[79], capacitive ac-bridges [30], [36], or switched-capacitor circuits [24], [35], [37], [38], [80]–[82].

Accelerometers can be operated open loop or closed loop. In the open-loop mode, the overall linearity, bandwidth, and dynamic range cannot be better than the parameters of the sensor structure by itself. However, open-loop sensors require simpler interface circuitry and are inherently stable in the frequency range below their resonance.

Closed-loop operation improves the overall sensor linearity, dynamic range, and bandwidth. The interface circuitry in a closed-loop device reads the sensor signal and uses feedback to maintain the proof mass in a null position in a stable manner. Loop stability can be ensured if the accelerometer is overdamped with a dominant pole [35] or is provided using a lead compensator in the feedback path [36], or a lead-lag [23] or lead compensator [38] in the forward path. Furthermore, the dynamics of the loop can be controlled by adding an electronic pole in the loop [75]. Closed-loop accelerometers using different schemes, including phase locked loop [79], analog [82], and pulse width modulation [83], have been developed. Recently, electromechanical, oversampled, sigma–delta modulators have become more attractive, since they provide direct digital output and force-feedback control of the proof mass simultaneously over a wide dynamic range [23], [24], [33], [35]–[38].

**D. Commercialization**

The application of micromachined accelerometers in automotive industry has been a major driver for their initial commercialization. EG&G IC Sensors and Lucas Nova Sensors were among the first companies in the United States that commercialized mainly piezoresistive, and later capacitive, microaccelerometers [14]–[16], [31]. However, one of the most successful micromachined accelerometers
in the market has been the ADXL50 by Analog Devices, which is a surface-micromachined, closed-loop, fully monolithic accelerometer integrated with its interface circuitry [71]. This accelerometer has been mainly developed for air-bag deployment and has a range of ±50 g with 6.6 mg/√Hz noise floor and shock survivability of >2000 g. Motorola has also been mass producing a capacitive accelerometer for air-bag applications [84]. This sensor has a range of ±40 g, a sensitivity of 40 mV/g, and a peak noise level of 15 mV in 400 Hz bandwidth, and utilizes a plastic dual in pin integrated circuit (IC) package to reduce cost. Other companies, including Delco and Ford, mass produce accelerometers for automotive applications. Fig. 8 shows Ford’s hybrid packaged accelerometer, which consists of a sensor chip, fabricated using bulk silicon micromachining, and a CMOS sigma–delta interface chip [33]. This accelerometer is designed for automotive passive restraint systems, has a range of 50 to 30 g with 400 Hz bandwidth, provides serial digital output with sensitivity of 3 kHz/g, has a self-test feature, and is available with a low-cost plastic surface mount device (SMD) package [74]. Endevco produces a large variety of piezoelectric, piezoresistive, and capacitive accelerometers with full-scale ranges from 2 to over 4000 g for a wide range of applications, including impact and crash sensing and vibration analysis [85]. In Europe, silicon accelerometers are in volume production by Bosch, and in Japan by Denso.

CSEM in Europe has produced one of the first high-precision silicon microaccelerometers [21], [22], [24]. These devices are capacitive bulk micromachined. One of the latest versions, MS6100, is targeted for low-power applications and achieves sub-mg resolution with a TCO of 50 μg/°C and stability of 1 mg for a 2 g device with less than 0.5 mW power dissipation [86]. In the United States, some companies active in navigation and guidance, such as Litton, have also been pursuing development of μg accelerometers [25].

E. Challenges and Future Trends

Today, microaccelerometers, with low and medium sensitivities, are in large volume production with prices under $5. This is a good indication that their technology is maturing and that numerous challenges involved in their design and mass production are resolved. Although a number of high-precision accelerometers have been already developed [21]–[25], most are not evaluated completely as inertial grade yet, and they are expensive for several applications and still need improvement to overcome long-term stability and temperature-sensitivity problems. Therefore, the development of low-cost, inertial-grade accelerometers with sub-μg noise levels, good long-term stability, and low temperature sensitivity still remains a challenge. The development of low-stress, low-drift packaging technologies for inertial-grade devices without affecting device performance and stability is also being pursued. A major challenge is the interface circuit, where low-drift readout/control circuitry with high sensitivity, low noise level, and large dynamic range is needed. It is believed that highly stable (inertial-grade) multiaxis devices with sub-μg resolution and over 120 dB dynamic range in several hundred hertz bandwidth will be developed in the near future.

III. MICROMACHINED GYROSCOPES

A. Operating Principles and Specifications

Almost all reported micromachined gyroscopes use vibrating mechanical elements to sense rotation. They have no rotating parts that require bearings, and hence they can be easily miniaturized and batch fabricated using micromachining techniques. All vibratory gyroscopes are based on the transfer of energy between two vibration modes of a structure caused by Coriolis acceleration. Coriolis acceleration, named after the French scientist and engineer G. G. de Coriolis (1792–1843), is an apparent acceleration
that arises in a rotating reference frame and is proportional to the rate of rotation. To understand the Coriolis effect, imagine a particle traveling in space with a velocity vector \( \mathbf{v} \). An observer sitting on the \( \mathbf{x} \)-axis of the \( \mathbf{x'y'z'} \) coordinate system, shown in Fig. 9(a), is watching this particle. If the coordinate system rotates around the \( \mathbf{z} \)-axis with an angular velocity \( \mathbf{\Omega} \), the observer thinks that the particle is changing its trajectory toward the \( \mathbf{x} \)-axis with an acceleration equal to \( 2\mathbf{v} \times \mathbf{\Omega} \). Although no real force has been exerted on the particle, to an observer, attached to the rotating reference frame an apparent force has resulted that is directly proportional to the rate of rotation. This effect is the basic operating principle underlying all vibratory structure gyroscopes.

Resolution, drift, zero-rate output (ZRO), and scale factor are important factors that determine the performance of a gyroscope. In the absence of rotation, the output signal of a gyroscope is a random function that is the sum of white noise and a slowly varying function [87]. The white noise defines the resolution of the sensor and is expressed in terms of the standard deviation of equivalent rotation rate per square root of bandwidth of detection \( \sqrt{\text{deg/s} / \sqrt{\text{Hz}}} \) or \( \sqrt{\text{deg/h} / \sqrt{\text{Hz}}} \). The so-called “angle random walk” in \( \sqrt{\text{deg} / \sqrt{\text{h}}} \) may be used instead. The peak-to-peak value of the slowly varying function defines the short- or long-term drift of the gyroscope and is usually expressed in \( \text{deg/s} \) or \( \text{deg/h} \) [87]. Scale factor is defined as the amount of change in the output signal per unit change of rotation rate and is expressed in \( \text{V/deg/s} \). Last, an important factor for any gyroscope that is primarily defined by device imbalances is the ZRO, which represents the output of the device in the absence of a rotation rate.

In general, gyroscopes can be classified into three different categories based on their performance: inertial-grade, tactical-grade, and rate-grade devices. Table 2 summarizes the requirements for each of these categories [87], [88]. Over the past few years, much of the effort in developing micromachined silicon gyroscopes has concentrated on “rate-grade” devices, primarily because of their use in automotive applications. This application requires a full-scale range of at least 50\( \circ/\text{s} \) and a resolution of about 0.1\( \circ/\text{s} \) in a bandwidth of 50 Hz, all at a cost of $10–20 [89]. The operating temperature is in the range from -40 to 85\( ^\circ \text{C} \). There are also several other applications that require improved performance, including inertial navigation, guidance, robotics, and some consumer electronics. Today, optical gyroscopes are the most accurate gyroscopes available in the market. Among these, ring laser gyroscopes have demonstrated inertial-grade performance, while fiber-optic gyroscopes are mainly used in tactical-grade applications. Delco’s hemispherical resonator gyroscope (HRG) is a vibratory gyroscope that has achieved impressive inertial-grade performance [90]. Although highly accurate, these devices are too expensive and bulky for many low-cost applications. Achieving “tactical- and inertial-grade” performance levels has proven to be a tough challenge for micromachined gyroscopes, and new technologies and approaches are being developed. Because of their greater compatibility with batch fabrication technologies, this paper will only review silicon micromachined vibratory gyroscopes.

B. Review of Micromachined Vibratory Gyroscopes

A number of vibratory gyroscopes have been demonstrated, including tuning forks [91]–[94], vibrating beams [95], and vibrating shells [89]. Tuning forks are a classical example of vibratory gyroscopes. The tuning fork, as illustrated in Fig. 9(b), consists of two tines that are connected to a junction bar. In operation, the tines are differentially resonated to a fixed amplitude, and when rotated, Coriolis force causes a differential sinusoidal force to develop on the individual tines, orthogonal to the main vibration. This force is detected either as differential bending of the tuning fork tines or as a torsional vibration of the tuning fork stem. The actuation mechanisms used for driving the vibrating structure into resonance are primarily electrostatic, electromagnetic, or piezoelectric. To sense the Coriolis-induced vibrations in the second mode, capacitive, piezoresistive, or piezoelectric detection mechanisms can be used. Optical detection is also feasible, but it is too expensive to implement. In general, silicon micromachining processes for fabrication of vibratory gyroscopes fall into one of four categories: 1) silicon bulk micromachining and wafer bonding; 2) polysilicon surface...
micromachining; 3) metal electroforming and LIGA; and 4) combined bulk-surface micromachining or so-called mixed processes.

Piezoelectric vibratory gyroscopes were demonstrated in the early 1980’s. Examples of these devices are fused-quartz HRG by Delco [90], quartz tuning forks [96] like the Quartz Rate Sensor by Systron Donner [97], [98], and a piezoelectric vibrating disc gyro [99]. Although quartz vibratory gyroscopes can yield very high quality factors at atmospheric pressure with improved level of performance, their batch processing is not compatible with IC fabrication technology. In the late 1980’s, after successful demonstration of batch-fabricated silicon accelerometers, some efforts were initiated to replace quartz with silicon in micromachined vibratory gyroscopes. The Charles Stark Draper Laboratory demonstrated one of the first batch-fabricated silicon micromachined rate gyroscopes in 1991. This bulk silicon device was a double gimbal vibratory gyroscope supported by torsional flexures, with the vibrating mechanical element made from p+ silicon [100].

As illustrated in Fig. 10, the outer gimbal was electrostatically driven at a constant amplitude using the drive electrodes, and this oscillatory motion was transferred to the inner gimbal along the stiff axis of the inner flexures. When exposed to a rotation normal to the plane of the device, Coriolis force causes the inner gimbal to oscillate about its weak axis with a frequency equal to the drive frequency. Therefore, maximum resolution is obtained when the outer gimbal is driven at the resonant frequency of the inner gimbal, causing the sensitivity to be amplified by the mechanical quality factor of the sense resonance mode of the structure. A rotation rate resolution of 4°/s in a 1 Hz bandwidth was realized using this structure.

Later in 1993, Draper reported an improved 1 mm² silicon-on-glass tuning fork gyroscope [91] fabricated through the dissolved wafer process [101]. This gyroscope was electrostatically vibrated in its plane using a set of interdigitated comb drives [102] to achieve a large amplitude of motion (10 μm). Any rotation in the plane of the substrate perpendicular to the drive mode will then excite the out-of-plane rocking mode of the structure, which is capacitively monitored. Fig. 11 shows an SEM view of the device with a perforated mass to minimize damping. The in-plane motion of the structure is lightly damped by air, while out-of-plane motion is strongly damped due to squeeze film effects. Therefore, for out-of-plane modes, $Q$ rises rapidly as pressure is reduced, in contrast to the in-plane $Q$, which shows a small increase as the pressure drops. At pressures of 100 mTorr, a $Q$ of 40,000 was observed for the drive mode and 5000 for the sense mode. The silicon-on-glass technology used in this device has the advantage of low stray capacitance. The noise equivalent rate observed by this structure was 470°/h in a 60 Hz bandwidth, equivalent to 0.02°/s in a 1 Hz bandwidth or angle random walk of 0.72°/√h [103]. The scale-factor accuracy was better than 0.1%, and bias stability was 55°/h overnight. The projected performance was 10–100°/h for bias stability and resolution in a 60 Hz bandwidth.

If the sense and drive resonant modes of a tuning fork have equal frequencies, the output signal will be amplified by the quality factor $Q$ of the sense mode, resulting in much higher sensitivity. However, this involves extreme control of device dimensions and may lead to temperature drift problems if these natural frequencies do not track with temperature. Because of these difficulties, most tuning-fork designs are not based on matched vibration mode frequencies.

Other tuning-fork designs have used electromagnetic excitation to obtain a large amplitude of motion [92], [93], [104]. Bosch’s silicon yaw rate sensor [93] achieves vibration amplitudes as large as 50 μm using a permanent magnet mounted inside a metal package. This device was fabricated through a combination of bulk- and surface-micromachining processes, and it consists of two bulk-micromachined oscillating masses, each of which supports two surface-micromachined accelerometers for detection of Coriolis force. The sensor chip is anodically bonded to a supporting glass wafer and is covered by another silicon cap wafer. Operating at atmospheric pressure, the device has shown a resolution of 0.3°/s in a 100 Hz bandwidth, thanks to its large amplitude of vibration. Although such a large amplitude of oscillation (50 μm) can increase the
output signal level, it increases the total power consumption and may cause fatigue problems over long-term operation. Cross talk between the sense and drive modes was minimized through mechanical decoupling of these modes by separating the oscillator and sense proof masses, resulting in a stable ZRO.

Piezoresistive detection has also been used in some gyroscope designs. Daimler Benz has demonstrated a tuning-fork angular rate sensor for automotive applications that piezoresistively measures the rotation-induced shear stress in the stem of the tuning-fork device [94]. In this device, a piezoelectric actuation mechanism was used by depositing a piezoelectric aluminum nitride (AlN) thin-film layer on one of the tines. The use of piezoelectric thin films such as AlN and ZnO on silicon degrades $Q$ and causes large temperature variation of offset and sensitivity [105]. This device was fabricated through a combination of bulk micromachining and bonding of SOI wafers. Researchers at the University of Neuchatel, Switzerland, have demonstrated a tuning-fork design based on two isolated vibrating proof masses, each supported by a four-beam bridge-type suspension [104]. These proof masses are electromagnetically vibrated in plane and antiphase, and the rotation-induced out-of-plane motion is then detected by means of four piezoresistors connected in a Wheatstone bridge configuration, showing a sensitivity of 4 nV/$\text{deg/s}$ with excellent linearity up to 750°/s. This device was fabricated through silicon bulk micromachining and was wafer-level vacuum packaged by anodic bonding of the silicon wafer to encapsulating glass wafers. In general, package-induced stress on the sensor structure can be lowered by low-temperature anodic bonding of glass wafers with silicon [106]. Although piezoresistive devices are easier to fabricate and require a simpler electronic interface due to their lower output impedance compared to capacitative devices, they have large temperature sensitivity and poor resolution.

Also reported in the literature are capacitive bulk-micromachined silicon-on-glass vibrating beams [107], vibrating membranes [108], and double-gimbaled structures [109]. Since the Young’s modulus of single-crystal silicon changes with crystallographic orientation, symmetric vibrating structures made of single-crystal silicon may show excessive mechanical coupling between drive and sense modes (due to this anisotropy), resulting in a large ZRO with unacceptable drift characteristics [107].

Surface-micromachined vibratory gyroscopes have also been demonstrated. Some have been integrated with the readout electronic circuitry on a single silicon chip, reducing parasitic capacitances and hence increasing the signal-to-noise ratio. In addition, the vibrating structure is made of polysilicon, which has a high quality factor and an orientation-independent Young’s modulus. Single- and dual-axis polysilicon surface-micromachined gyroscopes have been realized by researchers at Berkeley [110], [111] and Samsung [112]–[114]. Berkeley’s $z$-axis vibratory rate gyroscope [110] resembles a vibrating beam design and consists of an oscillating mass that is electrostatically driven into resonance using comb drives. Any deflections that result from Coriolis acceleration are detected differentially in the sense mode using interdigitated comb fingers. This device, 1 mm across, was integrated with a transresistance amplifier on a single die using the Analog Devices BiMEMS process. The remaining control and signal-processing electronics were implemented off chip. Quadrature error nulling and sense-mode resonant frequency tuning can be accomplished in this design by applying a control dc bias voltage to the position sense fingers. The dc bias voltage generates an electrostatic negative stiffness, which reduces the resonant frequency of the sense mode. By slightly changing this dc bias voltage on the differential comb fingers ($\pm \Delta V$), a lateral electrostatic field arises that can be used to align the drive mode oscillations and reduce the quadrature error. This device demonstrated a resolution of $1^\circ/s/\sqrt{Hz}$ with performance projected to improve to $0.1^\circ/s/\sqrt{Hz}$ in a second-generation device.

Samsung has also reported a very similar surface-micromachined $z$-axis device, shown in Fig. 12, with a 7-$\mu$m-thick polysilicon resonating mass supported by four fishhook-shaped springs [113]. This device, though not integrated with electronics, has demonstrated a resolution of $0.1^\circ/s$ at 2 Hz, an operating bandwidth of 100 Hz, and a linearity of 1% full scale in a range of 90°/s. Hybrid attachment of the sensor chip to a CMOS application-specific integrated circuit (ASIC) chip used for readout and closed-loop operation of the gyro was done in a vacuum-packaged ceramic case.

Murata has presented a surface-micromachined polysilicon gyroscope that is sensitive to lateral ($z$- or $y$-axis) angular rate [115]. The sense electrode was made underneath the perforated polysilicon resonator by diffusing phosphorous into the silicon substrate (junction isolation). This device showed an open-loop noise-equivalent rate of $2^\circ/s/\sqrt{Hz}$. The junction-isolation scheme used in this device, although simple, has the disadvantage of relatively large parasitic capacitance and large amount of shot noise associated with the existing pn junction, which in turn

![Fig. 12. SEM view of a comb-driven polysilicon surface micromachined $z$-axis vibratory gyroscope [113].](image-url)
Berkeley has reported a surface-micromachined dual-axis gyroscope based on rotational resonance of a 2-μm-thick polysilicon rotor disk, as shown in [116] (Fig. 13). Since the disk is symmetric in two orthogonal axes, the sensor can sense rotation equally about these two axes. This device, integrated with electronics, yielded a random walk as low as 10°/√h with cross-axis sensitivity ranging 3–16%. Resolution can be further improved to 2°/√h by frequency matching at the cost of excessive cross-axis sensitivity. Also reported in the literature is a cross-shaped nickel-on-glass two-axis micromachined gyroscope [117], which has shown a rate sensitivity of 0.1 mV/°/s.

The JPL, in collaboration with the University of California, Los Angeles, has demonstrated a bulk-micromachined, precision silicon MEMS vibratory gyroscope for space applications [118], [119]. This clover-leaf-shaped gyroscope consists of three major components: a silicon clover-leaf vibrating structure; a silicon baseplate, which is bonded to the clover-leaf structure; and a metal post, which is epoxied inside a hole on the silicon resonator. A hermetically sealed package, 1 × 1 × 0.7 in, houses the microgyroscope and most of its control electronics. This packaged gyroscope has a 7 Hz split between its drive and sense mode (f_{res} = 1.44 kHz), a scale factor of 24 mV/°/s, a bias stability of 70°/h, and an angle random walk of 6.3°/√h.

Recently, researchers at HSG-IMIT, Germany, have demonstrated and reported a surface-micromachined precision z-axis vibratory gyroscope (MARS-RR) with a very small ZRO achieved by mechanical decoupling of the drive and sense vibration modes [120]. This device (6 mm²), shown in Fig. 14, was fabricated through the standard Bosch foundry process featuring a 10-μm-thick structural polysilicon layer in addition to the buried polysilicon layer, which defines the sense electrodes. The reported performance of this device is quite impressive: a random angle walk of 0.27°/√h, a bias stability of 65°/h, and a scale-factor nonlinearity of <0.2%.

Researchers at General Motors and the University of Michigan have developed a vibrating ring gyroscope [89], schematically shown in Fig. 15. This device consists of a ring, semicircular support springs, and drive, sense, and balance electrodes, which are located around the structure. Symmetry considerations require at least eight springs to result in a balanced device with two identical flexural modes that have equal natural frequencies [121]. The ring is electrostatically vibrated into an in-plane elliptically shaped primary flexural mode with a fixed amplitude. When it is subjected to rotation around its normal axis, Coriolis force causes energy to be transferred from the primary mode to the secondary flexural mode, which is located 45° apart from the primary mode, causing amplitude to build up proportionally in the latter mode; this buildup is capacitively monitored. The vibrating ring structure has some important features compared to other types of vibratory gyroscopes. First, the inherent symmetry of the structure makes it less sensitive to spurious vibrations. Only when the ring has mass or stiffness asymmetries can environmental vibrations induce a spurious response. Second, since two identical flexural modes of the structure “with nominally equal resonant frequencies” are used to sense rotation, the sensitivity of the sensor is amplified by the quality factor of the structure, resulting in higher sensitivity. Third, the vibrating ring is less temperature sensitive since the vibration modes are affected equally by temperature. Last, electronic balancing of the structure is possible. Any frequency mismatch due to mass or stiffness asymmetries that occurs during the fabrica-
tion process can be electronically compensated by use of the balancing electrodes that are located around the structure.

The first micromachined version of the vibrating ring gyroscope was fabricated by electroforming nickel into a thick polyimide (or photoresist) mold on a silicon substrate in a post circuit process [89], [121], [122]. The gyroscope demonstrated a resolution of $\sim 0.5^\circ$/s in a 25 Hz bandwidth limited by the readout electronic noise. The sensor was integrated with a low-input capacitance source-follower buffer and the amplifier on a silicon chip. The zero bias drift was $<10^\circ$/s over the temperature range $-40$ to $85^\circ$C, and the sensitivity of the device varied by less than 3% over the same temperature range. Scale-factor nonlinearity in a $\pm 100^\circ$/s rate range was $<0.2^\circ$ [122].

To improve performance further, a new polysilicon ring gyroscope (PRG) [123] was recently fabricated through a single-wafer, all-silicon, high-aspect-ratio polysilicon trench-refill technology [124] at the University of Michigan. In this new process, the vibrating ring and support springs are created by refilling deep dry-etched trenches with polysilicon deposited over a sacrificial LPCVD oxide layer. Each sense electrode is made from a polysilicon island (12 $\mu$m deep) hanging over an ethylenediamine-pyrocatechol (EDP)-etched pit. Fig. 16 shows a SEM picture of a 1.7 $\times$ 1.7 mm$^2$ PRG. This device provides several important features required for high-performance gyroscopes, including small ring-to-electrode gap spacing ($<1 \mu$m) for increasing the sense capacitance; large structural height for increasing the radius and sense capacitance and reducing the resonant frequency; and a better structural material (polysilicon) for increasing $Q$ with an orientation-independent Young’s modulus. By taking advantage of these features, a tactical-grade ring microgyroscope with a reported root-mean-square noise floor of 0.15$^\circ$/s in a 30 Hz bandwidth and an in-run drift of approximately 0.05$^\circ$/s [125]. This device was fabricated through deep dry etching of a 100-$\mu$m-thick silicon wafer, which was then anodically bonded to a glass support wafer.

Levitated micromachined spinning-disc gyroscopes have also been investigated [126], [127]. The concept was based on a rotor disc, levitated using electromagnetic or electrostatic means and spun at a very high rate by means of a motor to produce angular momentum. With additional electrostatic fields, the rotor can be held in equilibrium even if the sensor is tilted or inverted. It is predicted that spinning microgyroscopes can yield a lower drift than a vibrating structure gyroscope [127]. The performance of these devices is yet to be demonstrated.

C. Design Issues and Considerations

Vibratory gyroscopes can be operated open or closed loop to measure rate of rotation (angular velocity). In the open-loop mode, the response to a change in rotation rate is not instantaneous, as time is required for the amplitude of the sense mode to reach its steady-state value. With matched sense and drive resonant modes, the time constant associated with this amplitude buildup is approximately equal to $2Q/\omega$ [121]. This response time limits the bandwidth of the sensor to a few hertz. To obtain larger open-loop bandwidth, gyroscopes are sometimes operated with a slight mismatch in the sense and drive modes; however, as theoretically shown in [128], this is done at the cost of reduced sensitivity. In the closed-loop mode of operation, the sense mode amplitude is continuously monitored and driven to zero, and hence the bandwidth and dynamic range of the sensor can be increased beyond the open-loop values even with matched resonant modes. The bandwidth is then limited by the readout and control electronics and can be increased to values approaching the resonant frequency of the structure.

With matched sense and drive resonant modes, assuming that the readout electronic noise $V_n$ has a white spectrum
around the resonant frequency of the structure and that the
detection circuit has a bandwidth of BW, the minimum
detectable electronic signal for a capacitive device can be expressed by

$$\Omega_{2\text{ (min)}} = \frac{\omega}{Q} \cdot \left( \frac{C_{\text{rest}} + C_{\text{paras}}}{\delta_{\text{rest}}} \right) \cdot V_{n} \cdot \sqrt{BW}, \quad (4)$$

This equation shows that reducing the noise of the readout
circuit, increasing the Coriolis-induced capacitance change
of the device, lowering the resonant frequency, increasing
the mechanical quality factor, and minimizing the parasitic
capacitances are determining factors in improving the reso-
lution. If the closed-loop bandwidth is not extended beyond
the mechanical quality factor, and minimizing the parasitic
circuit, increasing the Coriolis-induced capacitance change
(4 kHz). Stronger Coriolis forces will be obtained by
increasing the amplitude of vibration in the drive mode
($q_{\text{drive}}$), thus improving the resolution.

It should be noted that the fundamental limiting noise
component of the mechanical structure is due to the Brownian
motion of the sense vibration mode, and its total noise
equivalent acceleration is given by (3). This acceleration
noise corresponds to an input rotation rate whose spectral
density is determined using the transfer function
(from input rotation rate to the sense mode acceleration) of
each particular device. At high quality factors (operation in
vacuum), resolution is usually not limited by Brownian motion
of the structure; it is limited by the readout electronic
noise.

An important performance parameter for a vibratory gyroscope is its zero rate output (or zero bias). Geometrical
imperfections in the vibrating mechanical structure and/or
the sense and drive electrodes as well as electrical coupling
between these electrodes can cause an output signal in the
absence of rotation. For instance, in tuning-fork designs,
if due to fabrication flaws the tines’ centers of mass
are not precisely aligned in the plane of vibration, their
inertial forces produce a vibrating torque about the stem
just as the Coriolis torque does. This error, often called
the quadrature error, can be distinguished from the rate
signal simply because it is in phase quadrature with the
Coriolis-induced signal. If too large, it may cause errors in
sensing the rotation rate and can even saturate the amplifier
[110]. In the case of shell gyroscope designs, asymmetric
damping of the structure, which manifests itself as different
sense and drive mode quality factors, can also cause a
finite output when there is no rotation [88]. By electrically
and mechanically decoupling the sense and drive modes,
and by minimizing the fabrication process errors, ZRO
can be significantly reduced [93], [120]. A large quality
factor and lower natural frequency can further reduce the
ZRO, and its drift simply due to the fact that these two
factors increase the open-loop sensitivity of the sensor.
Therefore, high-quality materials with low internal damping
will improve the accuracy of the sensor. Any remaining
zero bias error should then be further reduced electronically.
For instance, the nickel ring gyroscope can be electronically
trimmed to reduce the zero bias from 600/s to less than
0.5/s.

A high-performance gyroscope should have an accurate
scale factor over a wide dynamic range as large as 140
dB. The scale factor should have a small temperature
sensitivity. Special attention should be paid to the structural
material. Metal structures used on silicon substrates or
piezoelectric layers cemented on silicon structures can
cause large variation of scale factor with temperature.
All-silicon devices are suitable for high-performance appli-
cations. Variation of resonant frequency over temperature
(due to temperature dependence of Young’s modulus) can
also affect the scale factor of the sensor, and therefore, some
temperature compensation of the scale factor is required
[121].

To obtain a high mechanical quality factor, both external
and internal energy losses should be minimized. Balanced,
symmetrical, isolated structures that do not radiate acoustic
energy out of their anchors are therefore required. In balanced symmetrical structures, the dependence of the
output signal on rigid body motion of the structure caused
by linear acceleration can be electronically rejected. An
element of a balanced system is a tuning-fork design,
made from two symmetrical tines oscillating antiphase,
which has no net motion at the junction and therefore
can be mounted inertially stable at this junction. Energy
losses due to squeeze film damping by air molecules will
be significantly reduced if the resonant device is operated
in vacuum. This calls for hermetically sealed, low-stress,
robust vacuum-packaging techniques that are capable of
holding vacuum levels (10⁻²–10⁻³ Torr) for extended pe-
riods of time. Silicon or glass wafers [93], [104] bonded to
the sensor substrate can provide hermetically sealed, chip-
level encapsulation [129]. Internal energy loss is reduced if
the structure is fabricated from low-loss materials such as
single- and polycrystalline silicon. Another concern about
a vibrating mechanical element is the long-term drift and
fatigue problem. Small vibration amplitude can help reduce
these problems.

Last, readout, signal-processing, and control electron-
ics are needed for open- and closed-loop operation of a
vibratory gyroscope. Part of the electronics drives the
first mode of the structure into resonance and keeps the
amplitude of vibration constant. Extremely small Coriolis-
induced motions of the structure are then detected by the
readout electronics. Low-noise readout techniques require
a careful biasing scheme using diodes [89], resistors [130],
and subthreshold MOS devices [110]. Monolithic inte-
gration of the sensor and the readout electronics reduces
the interconnect parasitics and improves the resolution of
the sensor. Parasitic capacitances can be further reduced
by special bootstrapping techniques [121], [130]. Delco
Electronics has demonstrated a CMOS ASIC chip for
closed-loop operation of an electroformed ring gyroscope
[131]. It consists of four feedback loops and provides
the necessary electronics to compensate for the sensor
nonuniformities and its offset and sensitivity variations with temperature (−40 to +85°C). High-performance gyroscopes will continue to need precision, low-noise, parasitic-insensitive interface circuits capable of resolving attofarad changes in capacitance with as small an input capacitance as possible.

D. Commercialization and Future Trends

Mainly driven by the automotive industry, micromachined silicon gyroscopes have been the subject of extensive research and development over the past few years. The performance of micromachined gyroscopes has drastically improved over a rather short period, as illustrated in Fig. 17, which is based on a sample of devices reported in the literature. Since 1991, performance, indicated by the random angle walk, has improved by a factor of 10× every two years for both bulk- and surface-micromachined devices. The plot also shows that surface-micromachined devices have been closing the gap with their bulk counterparts. Although this general trend is not expected to continue indefinitely, it is anticipated that for at least the next 5–10 years, we will see a continuing improvement in the performance of both bulk- and surface-micromachined gyroscopes.

Much effort is also under way for large-volume production of micromachined gyroscopes. Production cost, performance, and reliability are the key factors in commercializing micromachined gyroscopes. Precision micromachining, robust vacuum packaging, and high-performance interface circuit and electronic tuning techniques are required to reduce the production cost to a level that is acceptable for the large-volume automotive market [132]. Since 1993, Draper and Rockwell International have been collaborating to commercialize Draper’s silicon tuning-fork gyroscope targeted for automotive applications [133]. Other companies like General Motors, Analog Devices,
and Samsung have also performed extensive research and development on various gyroscope structures, and prototypes have been demonstrated [132]. Bosch [93] has started volume production of its yaw sensor, and it is clear that several other companies will soon be doing the same.

In general, all-silicon, mixed-mode (bulk/surface) fabrication technologies, combined with high-aspect-ratio deep dry-etching techniques, can provide features that are required for future high-performance microgyroscopes. Formation of submicrometer capacitive gaps through sacrificial layer etching, high-aspect-ratio and thick structures with high quality factor and uniform material properties, along with chip-level vacuum packaging, will help improve the performance by orders of magnitude. Thick, high-aspect-ratio structures formed by deep dry etching can yield a large sense capacitance with a larger oscillating mass, which in turn will help improve performance and simplify packaging. By shrinking the capacitive gaps to submicrometer levels, bias and control voltages will also shift down to CMOS acceptable levels. Future high-performance tactical- and inertial-grade gyroscopes will make use of dynamic electronic tuning of the structure to compensate for temperature and long-term drift effects of the sensor.

Further progress is also anticipated in the development of integrated multiaxis devices on a single chip. Shown in Fig. 18 is a six-degree-of-freedom, fingernail-sized inertial measurement unit (IMU) with integrated signal-processing and control circuitry, which is designed at the University of California, Berkeley, and fabricated by Sandia National Laboratories [37], [116]. This unit combines a dual-axis rate gyroscope, a z-axis gyroscope, and a three-axis accelerometer chip to measure angular rate and acceleration without the need to align individual sensors.

IV. CONCLUSIONS

Micromachined inertial sensors have progressed rapidly since the first device was demonstrated in 1979 by researchers at Stanford University. Microaccelerometers are now in large-volume production, cost a few dollars, incorporate many functions, including self-testing, and have been shown to be extremely reliable for many years. This impressive level of performance has been achieved by resolving many obstacles in areas of device fabrication, interface and readout circuitry, assembly and packaging, and testing. One can routinely measure sub-mg accelerations, and prototypes capable of providing μg performance have been demonstrated.

By taking advantage of many of the technologies developed for microaccelerometers, microgyroscopes have been able to enhance their performance by a factor of 10× every two years since 1991, and there is every indication that this trend will continue for at least the next few years. Based on past performance, it is clear that micromachined gyroscopes will be capable of providing inertial-grade performance at orders of magnitude less cost and size.

In the next decade, much effort will be expended to advance precision micromachining further, develop low-noise and low-drift interface circuitry, and provide reliable low-cost packaging for development of even higher performance, lower cost, and lower power inertial sensors for many emerging and as yet unknown applications.

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REFERENCES


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