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EE C245 - ME C218 Introduction to MEMS Design Fall 2012

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Lecture Module 7: Mechanics of Materials

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Outline

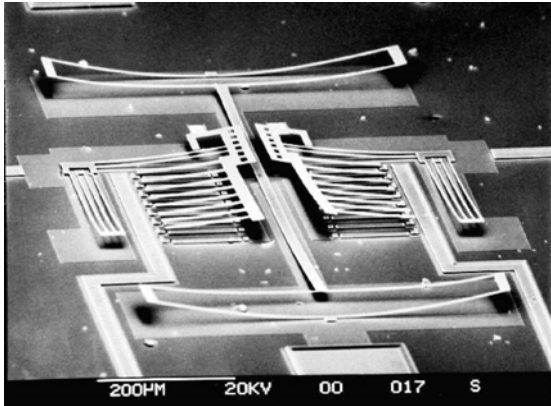
- Reading: Senturia, Chpt. 8
- Lecture Topics:
 - ↗ Stress, strain, etc., for isotropic materials
 - ↗ Thin films: thermal stress, residual stress, and stress gradients
 - ↗ Internal dissipation
 - ↗ MEMS material properties and performance metrics

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Vertical Stress Gradients

- Variation of residual stress in the direction of film growth
- Can warp released structures in z-direction



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Elasticity

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Normal Stress (1D)

If the force acts normal to a surface, then the stress is called a **normal stress**

Force assumed uniform over the whole area A

Stress = $\left\{ \begin{array}{l} \text{Force per} \\ \text{unit area} \end{array} \right\} = \sigma = \frac{F}{A} \quad [N/m^2 = Pa]$
 \leftarrow standard mks unit

\Rightarrow Microscopic Definition: force per unit area acting on the surface of a differential volume element of a solid body

\Rightarrow Note: assume stress acts uniformly across the entire surface of the element, not at just a point

Differential volume element

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Strain (1D)

Sometimes a unit called the "microstrain" is used, where $10^{-6} = 1 \mu\epsilon = \frac{\Delta L}{L}$ of 1 part in 10^6

Strain = $\left\{ \begin{array}{l} \text{Fractional Change} \\ \text{in length} \end{array} \right\} = \epsilon = \frac{L' - L}{L} = \frac{\Delta L}{L}$
 [unitless]

In the elastic regime (i.e., for "small" stresses at "low" temperatures), strain is found to be proportional to stress

For solids: MPa \rightarrow GPa

slope = E = Young's modulus of elasticity

$\sigma = \epsilon E \rightarrow \epsilon = \frac{\sigma}{E}$ [unitless]

Thus, the units of E are the same as $\sigma \rightarrow Pa$

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The Poisson Ratio

Apply normal stress to a free-standing object $\left\{ \begin{array}{l} \rightarrow \text{uniaxial strain} \\ \rightarrow \text{but also get contraction in directions transverse to the uniaxial strain} \end{array} \right.$

\Rightarrow contraction creates a (-) strain:

$$\epsilon_y = \frac{W' - W}{W} = \frac{\Delta W}{W} = -\nu \epsilon_x$$

$\leftarrow \nu$ = Poisson ratio [unitless]

\leftarrow typical values: 0 \rightarrow 0.5

\Rightarrow inorganic solids: 0.2 \rightarrow 0.3

\Rightarrow elastomers (e.g., rubber): ~ 0.5

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Shear Stress & Strain (1D)

Note: Assume compensating forces are applied to the vertical faces to avoid a net torque. (This by convention.)

Shear Stress = $\left\{ \begin{array}{l} \text{Force per Unit Area} \\ \text{Parallel to the Surfaces} \end{array} \right\} = \tau = \frac{F}{A} \quad [Pa]$
 \leftarrow not a rotator

Generates a shear strain:

Shear Strain = $\theta = \frac{\tau}{G}$

$G \triangleq$ shear modulus

$$G = \frac{E}{2(1 + \nu)}$$

\leftarrow differential volume element

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2D and 3D Considerations

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- Important assumption: the differential volume element is in static equilibrium \rightarrow no net forces or torques (i.e., rotational movements)
 - Every σ must have an equal σ in the opposite direction on the other side of the element
 - For no net torque, the shear forces on different faces must also be matched as follows:

Stresses acting on a differential volume element

$$\tau_{xy} = \tau_{yx} \quad \tau_{xz} = \tau_{zx} \quad \tau_{yz} = \tau_{zy}$$

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2D Strain

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- In general, motion consists of
 - rigid-body displacement (motion of the center of mass)
 - rigid-body rotation (rotation about the center of mass)
 - Deformation relative to displacement and rotation

Area element experiences both displacement and deformation

- Must work with displacement vectors
- Differential definition of axial strain: $\epsilon_x = \frac{u_x(x + \Delta x) - u_x(x)}{\Delta x} = \frac{\partial u_x}{\partial x}$

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2D Shear Strain

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Rotate clockwise by θ_1

\Rightarrow For shear strains, must remove any rigid body rotation that accompanies the deformation

\hookrightarrow use a symmetric definition of shear strain:

$$\tau_{xy} = \theta_2 + \theta_1 \approx \left(\frac{\Delta u_x}{\Delta y} + \frac{\Delta u_y}{\Delta x} \right) = \left(\frac{\partial u_x}{\partial y} + \frac{\partial u_y}{\partial x} \right)$$

For small amplitude deformations.

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Volume Change for a Uniaxial Stress

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Stresses acting on a differential volume element

Given an x -directed uniaxial stress, σ_x :

$$\begin{aligned} \Delta x &\rightarrow \Delta x (1 + \epsilon_x) \\ \Delta y &\rightarrow \Delta y (1 - \nu \epsilon_x) \\ \Delta z &\rightarrow \Delta z (1 - \nu \epsilon_x) \end{aligned}$$

The resulting change in volume ΔV

$$\begin{aligned} \Delta V &= \Delta x \Delta y \Delta z [(1 + \epsilon_x)(1 - \nu \epsilon_x)^2 - 1] \\ &= \Delta x \Delta y \Delta z [(1 + \epsilon_x)(1 - 2\nu \epsilon_x + \nu^2 \epsilon_x^2) - 1] \end{aligned}$$


{Assume small strains} $\Rightarrow \Delta V = \Delta x \Delta y \Delta z [(1 + \epsilon_x)(1 - 2\nu \epsilon_x) - 1]$

$[(1 + m)x]^n \approx 1 + nm x \Rightarrow \Delta V \approx \Delta x \Delta y \Delta z [1 + \epsilon_x - 2\nu \epsilon_x - 1]$

$\Delta V \approx \Delta x \Delta y \Delta z (1 - 2\nu) \epsilon_x$

For $\nu = 0.5$ (rubber) \rightarrow no ΔV !
 $\nu < 0.5 \rightarrow$ finite ΔV

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Isotropic Elasticity in 3D

- Isotropic = same in all directions
- The complete stress-strain relations for an isotropic elastic solid in 3D: (i.e., a generalized Hooke's Law)

$$\begin{aligned}\varepsilon_x &= \frac{1}{E} \left[\sigma_x - \nu (\sigma_y + \sigma_z) \right] & \gamma_{xy} &= \frac{1}{G} \tau_{xy} \\ \varepsilon_y &= \frac{1}{E} \left[\sigma_y - \nu (\sigma_z + \sigma_x) \right] & \gamma_{yz} &= \frac{1}{G} \tau_{yz} \\ \varepsilon_z &= \frac{1}{E} \left[\sigma_z - \nu (\sigma_x + \sigma_y) \right] & \gamma_{zx} &= \frac{1}{G} \tau_{zx}\end{aligned}$$

Basically, add in off-axis strains from normal stresses in other directions

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