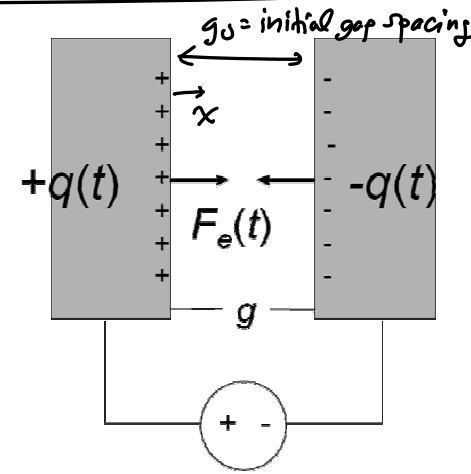


Lecture 18: ke and Comb Drive

- Announcements:
- HW#5 online; due Wednesday, April 15
- Project slide #1 due Friday, April 10
- I will be on travel next week; will either record lectures or do make-up lectures next Friday
- Will announce via Piazza
- -----
- Reading: Senturia, Chpt. 5, Chpt. 6
- Lecture Topics:
  - ↳ Energy Conserving Transducers
    - Charge Control
    - Voltage Control
  - ↳ Parallel-Plate Capacitive Transducers
    - Linearizing Capacitive Actuators
    - Electrical Stiffness
  - ↳ Electrostatic Comb-Drive
    - 1<sup>st</sup> Order Analysis
    - 2<sup>nd</sup> Order Analysis
- -----
- Last Time:
- Looking at methods for linearizing the voltage-to-force transfer function of a parallel-plate capacitive-gap transducer

↻ over

Linearizing the Voltage-to-force Transfer Function



$$v(t) = \underbrace{V_P}_{\text{Bias (DC)}} + \underbrace{v_i(t)}_{\text{Signal (AC)}}$$

$$F_e(t) = \frac{\partial W'}{\partial x} = \frac{\partial}{\partial x} \left[ \frac{1}{2} C [v(t)]^2 \right]$$

$$= \frac{1}{2} \frac{\partial C}{\partial x} [v(t)]^2 = \frac{1}{2} \frac{\partial C}{\partial x} (V_P + v_i(t))^2$$

$$= \frac{1}{2} [V_P^2 + 2V_P v_i(t) + \cancel{v_i(t)^2}] \frac{\partial C}{\partial x}$$

$$[V_P \gg v_i(t)] \Rightarrow F_e(t) = \underbrace{\frac{1}{2} V_P^2 \frac{\partial C}{\partial x}}_{\text{DC offset}} + \underbrace{V_P \frac{\partial C}{\partial x} v_i(t)}_{\text{AC drive signal}}$$

What's  $\frac{\partial C}{\partial x}$ ?

$$C_0 = \frac{\epsilon A}{g_0} \rightarrow C(x) = \frac{\epsilon A}{g_0 - x} = C_0 \left(1 - \frac{x}{g_0}\right)^{-1}$$

$$[x \ll g_0] \Rightarrow \approx C_0 \left(1 + \frac{x}{g_0}\right) = C(x)$$

$\therefore \frac{\partial C}{\partial x} = \frac{C_0}{g_0} = \frac{\epsilon A}{g_0^2}$   
 $\Rightarrow F_e(t) = \frac{1}{2} \frac{C_0}{g_0} V_p^2 + V_p \frac{C_0}{g_0} V_i(t)$

$V_i(t) = V_i \cos \omega t$   
 this is a linear dependence on  $V_i(t)$   
 holds for small amplitudes

$\frac{|X|}{|F_e|}$  vs  $\omega$   
 DC offset  
 Signal Component  
 $\omega_0$   
 very small response  $\rightarrow$  but must still worry about pull in  $V_{PI}$

Cancel the DC offset using Differential Symmetry

$F_{net} = F_{er}(t) - F_{el}(t)$   
 $= \frac{1}{2} \frac{\partial C}{\partial x} \{ [V_R(t)]^2 - [V_L(t)]^2 \}$   
 $= \frac{1}{2} \frac{\partial C}{\partial x} \{ V_p^2 + 2V_p v(t) + [v(t)]^2 - (V_p^2 - 2V_p v(t) + [v(t)]^2) \}$   
 $\therefore F_{net}(t) = 2V_p \frac{\partial C}{\partial x} v(t) = 2V_p \frac{C_0}{g_0} v(t)$

No DC component.  $\rightarrow$  less pull-in problems, but linear w/  $v(t)$ ! only to the extent of matching

Nonlinearity Still Effects Us

More Complete Expressions

$$C_1(x) = \frac{\epsilon A}{d_1 + x} = C_{01} \left(1 + \frac{x}{d_1}\right)^{-1} \rightarrow \frac{\partial C_1}{\partial x} = -\frac{C_{01}}{d_1} \left(1 + \frac{x}{d_1}\right)^{-2}$$

[Expand into Taylor series]

$$\frac{\partial C_1}{\partial x} = -\frac{C_{01}}{d_1} \left(1 + A_1 x + A_2 x^2 + A_3 x^3 + \dots\right)$$

where  $A_1 = -\frac{2}{d_1}, A_2 = \frac{3}{d_1^2}, A_3 = -\frac{4}{d_1^3}, \dots$

$$F_{d1} = \frac{1}{2} \frac{\partial C_1}{\partial x} (V_p - V_1 - v_1)^2 = \frac{1}{2} \frac{\partial C_1}{\partial x} (V_{p1} - v_1)^2$$

[small displacements:  $x \ll d_1$ ]

$$F_{d1} = \frac{1}{2} \left(-\frac{C_{01}}{d_1}\right) (1 + A_1 x) (V_{p1}^2 - 2V_{p1}v_1 + v_1^2)$$

$$= \frac{1}{2} \left(-\frac{C_{01}}{d_1}\right) \left\{ V_{p1}^2 - 2V_{p1}v_1 + v_1^2 + A_1 V_{p1}^2 x - 2A_1 V_{p1}xv_1 + A_1 x v_1^2 \right\}$$

Resonance:

@ resonance:

$$x = \frac{Q F_{d1}}{jk} \approx \frac{Q}{jk} \frac{\partial C_1}{\partial x} V_{p1} v_1$$

approx.  $F_{d1} \leftarrow$  linear term

90° phase shift

$$v_1 = |v_1| \cos \omega t \rightarrow x = |x| \sin \omega t$$

90° phase-shifted

Pick out force terms @  $\omega_0$

$$F_{d1}|_{\omega_0} = V_{p1} \frac{C_{01}}{d_1} \cos \omega_0 t + V_{p1}^2 \frac{C_{01}}{d_1^2} x \sin \omega_0 t$$

drive force term

$k_e \rightarrow$  electrical stiffness

proportional to  $x$

90° phase-shifted from  $\therefore$  in phase w/ displacement!

force term in phase as displacement  $\rightarrow$  it's a stiffener!

Electrical Stiffness:

- ① A negative spring constant!
- ② Derive from  $V_p$ :

$$k_e = V_{p1}^2 \frac{C_{01}}{d_0^2} = V_{p1}^2 \frac{\epsilon A^2}{d_1^3}$$

overlap area of C

DC Bias

gap  $\rightarrow$  3rd order dependence on gap!

$k_e \rightarrow$  can affect the resonance freq,  $f_0$ !

$\omega_0 \triangleq$  radian resonance freq. w/ no  $V_p$  applied (i.e.,  $V_{p1} = 0V$ )

$$\omega_0' = \sqrt{\frac{k}{m}} = \sqrt{\frac{k_m - k_e}{m}}$$

$$= \sqrt{\frac{k_m}{m}} \left(1 - \frac{k_e}{k_m}\right)^{1/2}$$

$$\omega_0' = \omega_0 \left[1 - \frac{V_{p1}^2 \epsilon A^2}{k_m d_1^3}\right]^{1/2}$$

now a fun of dc bias voltage!  
(voltage-controllable!)

- Go through Module 12 slides 26-35

**Electrostatic Comb-Drive**

Top View

Side View

$V_P$

$V_i$

Shuttle Finger

Drive Finger

$d$

$x$

$L_f$

$h \leftarrow \text{thickness}$

$y$

$z$

$x$

$F_d = \frac{\partial W'}{\partial x} = \frac{1}{2} \frac{\partial C}{\partial x} (V_P - V_i)^2$  Need C(x).

$C(x) = \frac{2\epsilon_0 x h}{d} \rightarrow \frac{\partial C}{\partial x} = \frac{2\epsilon_0 h}{d}$  Not a fun of x!

$F_d = \frac{1}{2} \frac{2\epsilon_0 h}{d} (V_P^2 - 2V_P V_i + V_i^2)$   $k_e = 0!$

can balance out by symmetrically placed electrodes  $V_i \ll V_P$

$F_d = -2V_P \frac{\epsilon_0 h}{d} V_i$

- Go through remaining comb-drive slides in Module 12