C191 - Homework 2

1. The secular approximation - The *hyperfine contact interaction* in Hydrogen governs the interaction between the electron spin and the nuclear proton spin. This interaction is represented by the Hamiltonian,

$$H_0 = J\vec{\sigma}^{(e)} \cdot \vec{\sigma}^{(p)} \tag{1}$$

$$= J\left(\sigma_x^{(e)} \otimes \sigma_x^{(p)} + \sigma_y^{(e)} \otimes \sigma_y^{(p)} + \sigma_z^{(e)} \otimes \sigma_z^{(p)}\right)$$
(2)

$$= J \left(2\sigma_+^{(e)} \otimes \sigma_-^{(p)} + 2\sigma_-^{(e)} \otimes \sigma_+^{(p)} + \sigma_z^{(e)} \otimes \sigma_z^{(p)} \right)$$
(3)

Where the superscripts (e) and (p) indicate that the operator acts on the electron or proton, respectively, and the spin raising/lowering matrices are:

$$\sigma_{-} = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \qquad \sigma_{+} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

Suppose now that a large magnetic field, B, is applied, interacting with the spins via the following Hamiltonian,

$$H_1 = \frac{1}{2}g\mu_B B\sigma_z^{(e)} - \frac{1}{2}g\mu_n B\sigma_z^{(p)}$$

Notice that the sign and magnitude of the interaction is different for the electron and the proton: the nuclear magneton, $\mu_n = e\hbar/2m_p$, is roughly 2000 times smaller than the Bohr magneton, $\mu_B = e\hbar/2m_e$, because the ratio of the proton mass to the electron mass is $m_p/m_e \simeq 2000$, while the sign difference is due to the opposite signs of the electron and proton charge. The total Hamiltonian is then $H = H_0 + H_1$. Show that, if the magnetic field is sufficiently large, then we can take the *secular approximation*,

$$H \simeq H_{\rm sec} = J\sigma_z^{(e)}\sigma_z^{(p)} + \frac{1}{2}g\mu_B B\sigma_z^e - \frac{1}{2}g\mu_n B\sigma_z^p$$

To do this,

(a) Transform into a rotating frame with respect to H_1 . That is, determine the Hamiltonian governing the evolution of the state

$$|\psi'(t)\rangle = \exp\left(iH_1t/\hbar\right)|\psi(t)\rangle$$

Where $|\psi(t)\rangle$ obeys the Schrödinger equation:

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = (H_0 + H_1) |\psi(t)\rangle$$

It will help to prove the following identities:

$$\exp\left(i\alpha\sigma_z\right)\sigma_+\exp\left(-i\alpha\sigma_z\right) = e^{2i\alpha}\sigma_+$$

$$\exp\left(i\alpha\sigma_z\right)\sigma_{-}\exp\left(-i\alpha\sigma_z\right) = e^{-2i\alpha}\sigma_{-}$$

- (b) Argue that all quickly oscillating terms may be neglected. Why does the magnetic field need to be large?
- (c) Transform back into the lab frame.
- Spin echo Suppose that a qubit, initially in the state |+>, is subjected to a magnetic field of unknown strength, B, experiencing a Hamiltonian,

$$H = KB\sigma_2$$

- (a) What is the state of the qubit after a time, τ ?
- (b) At time τ , a strong field is applied, causing a unitary operator, σ_x to be applied to the qubit. What is the state of the qubit now?

(c) The qubit is allowed to evolve under the Hamiltonian H for an additional time τ . What is the final state of the qubit?

3. Circuit model exercises

(a) Show that a CNOT can be achieved from a CPHASE and two Hadamard operations

(b) Show the reverse,

(c) Show that a SWAP gate may be implemented as three CNOTs

(d) For a matrix $U = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$, what is the matrix representation of the following quantum circuits i.



iii.

iv.

ii.

_____U

 U	

4. Single qubit gates We saw in class that the Hamiltonian,

$$H = -\frac{1}{2}g\mu_B \left(B_0\sigma_z + B_1\cos(\omega t)\sigma_x\right)$$

could generate a unitary operator $U(\tau) \propto \sigma_x$ if applied for a time, τ . Now let's change the phase on the oscillating term,

$$H_1 = -\frac{1}{2}g\mu_B \left(B_0\sigma_z + B_1\sin(\omega t)\sigma_x\right)$$

- (a) If this Hamiltonian is allowed to act for the same time, τ , what unitary operator is generated (in the rotating frame)?
- (b) If we wanted to generate the unitary operator,

$$U \propto \cos(\theta)\sigma_x + \sin(\theta)\sigma_y,$$

what Hamiltonian could we apply to do this? You should check that this operator is, in fact, unitary!