## 1 Unitarity of a Fourier Transform

The Fourier transform mod N is the  $N \times N$  matrix given by

$$FT_N = \frac{1}{\sqrt{N}} \begin{pmatrix} 1 & 1 & 1 & \cdots & 1\\ 1 & \omega & \omega^2 & \cdots & \omega^{N-1}\\ 1 & \omega^2 & \omega^4 & \cdots & \omega^{2(N-1)}\\ \vdots & \vdots & \vdots & \ddots & \vdots\\ 1 & \omega^{N-1} & \omega^{2(N-1)} & \cdots & \omega^{(N-1)^2} \end{pmatrix} , \qquad (1)$$

where  $\omega = e^{2\pi i/N}$  is a primitive Nth root of unity. That is, the i, j'th element of  $FT_N$  is  $\frac{1}{\sqrt{N}}\omega^{ij}$ , for  $i, j = 0, \dots, N-1$ .

Equivalently, in ket notation, for  $j \in \{0, 1, ..., N-1\}$ ,  $FT_N|j\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} \omega^{ij} |i\rangle$ . We need to check the inner product between the *i*th and *j*th columns of  $FT_N$ , that  $\langle i|FT_N^{\dagger}FT_N|j\rangle = \delta_{ij} \equiv \begin{cases} 1 \text{ if } i=j\\ 0 \text{ if } i\neq j \end{cases}$ . Indeed, this inner product is

$$\frac{1}{N} \sum_{k=0}^{N-1} \overline{\omega^{ik}} \omega^{jk} = \frac{1}{N} \sum_{k=0}^{N-1} \omega^{k(j-i)}$$
(2)

This is a geometric series with ratio between terms  $\omega^{j-i}$  and so can easily be evaluated. If  $i = j \mod N$ , then each term is  $\omega^0 = 1$ , so the inner product is N/N = 1. If  $i \neq j$ , then the sum is

$$1 + \omega^{j-i} + \omega^{2(j-i)} + \dots + \omega^{(N-2)(j-i)} + \omega^{(N-1)(j-i)} .$$
(3)

Multiplying the sum by  $\omega^{i-j} \neq 1$  gives

$$\omega^{j-i} + \omega^{2(j-i)} + \omega^{3(j-i)} + \dots + \omega^{(N-1)(j-i)} + \omega^{N(j-i)} .$$
(4)

But  $\omega^{N(j-i)} = (\omega^N)^{j-i} = 1$ , so we have just rearranged the terms of the summation; the sum itself doesn't change when multiplied by  $\omega^{j-i}$ . Therefore the sum is zero, as claimed.

## 2 Fourier Transforms and the uncertainty principle

a) Prove that for any quantum state  $|\psi\rangle$  on *n* qubits,  $S(|\psi\rangle) \leq 2^{n/2}$ .

**Answer**: We need to show that if  $\sum_{x \in \{0,1\}^n} |\alpha_x|^2 = 1$ , then  $\sum_x |\alpha_x| \leq 1$  $2^{n/2}$ . Using the Cauchy-Schwarz inequality  $\langle v|w\rangle \leq ||v|| \cdot ||w||$ , we get

$$\sum_{x} |\alpha_{x}| = \sum_{x} (|\alpha_{x}| \cdot 1)$$
$$\leq \left(\sum_{x} |\alpha_{x}|^{2}\right)^{1/2} \left(\sum_{x} 1^{2}\right)^{1/2}$$
$$= 1 \cdot 2^{n/2} = 2^{n/2} ,$$

with equality iff  $|\alpha_x| = 1/2^{n/2}$  for all x. (b) **Answer**: Using the normalization condition,  $1 = \sum_x |\alpha_x|^2 \le \sum_x a |\alpha_x| = aS(|\psi\rangle)$ . (Notice that this inequality is an equality iff all  $\alpha_x$  are zero or exactly a – that is, to minimize the spread, concentrate the probability mass as much as possible while still satisfying the constraint  $|\alpha_x| \leq a$ .)

$$H^{\otimes n}|x\rangle = \bigotimes_{i=1}^{n} \left( \begin{array}{c} |0\rangle + |1\rangle \text{ if } x_{i} = 0\\ |0\rangle - |1\rangle \text{ if } x_{i} = 1 \end{array} \right)$$
$$= \sum_{z} \left( \bigotimes_{i=1}^{n} \begin{array}{c} |z_{i}\rangle \text{ if } x_{i} = 0 \text{ or } z_{i} = 0\\ -|z_{i}\rangle \text{ if } x_{i} = 1 \text{ and } z_{i} = 1 \end{array} \right)$$
$$= \sum_{z} \left( \prod_{i=1}^{n} (-1)^{x_{i}z_{i}} \right)|z\rangle$$
$$= \sum_{z} (-1)^{x \cdot z}|z\rangle$$

(d) Use (c) to prove that for all y,  $|\beta_y| \leq \frac{1}{2^{n/2}}S(|\psi\rangle)$ . **Answer**: Using the triangle inequality  $|a+b| \leq |a|+|b|$ ,

$$|\beta_y| = \frac{1}{2^{n/2}} |\sum_y (-1)^{x \cdot y} \alpha_y|$$
  

$$\leq \frac{1}{2^{n/2}} \sum_y |(-1)^{x \cdot y} \alpha_y|$$
  

$$= \frac{1}{2^{n/2}} \sum_y |\alpha_y|$$
  

$$= \frac{1}{2^{n/2}} S(|\psi\rangle) .$$

(e) Prove the uncertainty relation  $S(|\psi\rangle)S(H^{\otimes n}|\psi\rangle) \geq 2^{n/2}$ . Justify why it makes sense to call this an uncertainty relation.

**Answer**: By part d,  $|\beta_y| \leq \frac{1}{2^{n/2}} S(|\psi\rangle)$  for all y. So by part b,  $S(H^{\otimes n}|\psi\rangle) \geq 2^{n/2}/S(|\psi\rangle)$ , which is the desired inequality.

This is an uncertainty relation because it gives a tradeoff between the spread in one basis and the spread in another. For example, if a state is well-concentrated in the standard basis, then it has high spread – and therefore high uncertainty – in the Fourier basis.