

Lecture 25: High Frequency Circuit Analysis

Announcements:

- HW#8 online and due Friday via Gradescope
- Lab#5 due Tuesday, Nov. 12, 5 p.m.
- This is a recorded video, since the PG&E shutdown forced class cancellation on Monday, 10/28

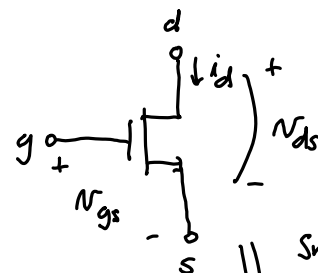
Lecture Topics:

- ↔ MOS High Frequency Model
- ↔ Brute Force CE HF Analysis
- ↔ Open Circuit Time Constant (OCTC) Analysis

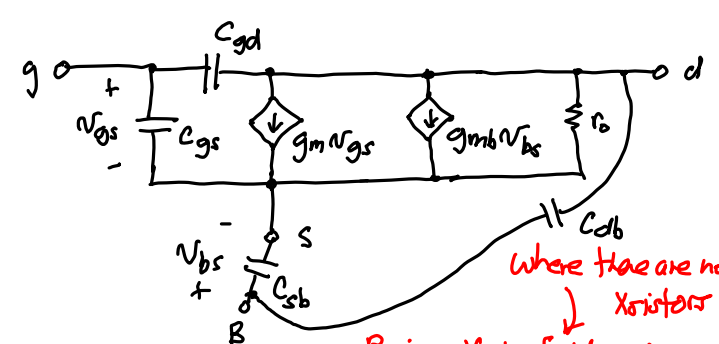
Last Time:

- Covered BJT high frequency model
- Now, continue with MOS high frequency model ...

MOS High Frequency Small-Signal Model

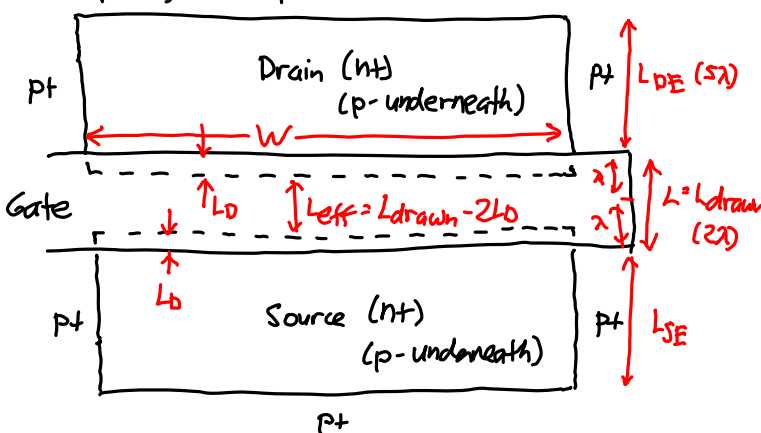


Small-Signal Equivalent Ckt.



Where there are no resistors
 Raise V_t in field regions

MOS Layout View - (top view)



MOS High Frequency S.S. Parameters (saturated NMOS)

$C_g = \text{gate capacitance} = W L_{\text{eff}} C_{ox}$ (linear) *in literature, be careful about the "out the"*
 $= \frac{2}{3} W L_{\text{eff}} C_{ox}$ (saturation)

C_{ox} just understood as a C per unit area

Field Channel Stop Implant

$C_{gsOL} = \text{gate-to-source overlap capacitance} = W L_D C_{ox}$

$C_{gdOL} = \text{gate-to-drain overlap capacitance} = W L_D C_{ox}$

Shielded Out Depletion Capacitance, C_D (when inversion layer present)

$C_{sb} = \text{source-to-bulk junction capacitance (underside area)}$

$C_{db} = \text{drain-to-bulk junction capacitance (underside area)}$

$C_{dbsw} = \text{drain-to-bulk junction sidewall capacitance}$

smaller depletion region $\rightarrow \therefore$ larger C

Gate-to-Source Capacitor, C_{gs} : (saturated MOS)

$$C_{gs} = C_{OL} + \frac{2}{3} W L_{\text{eff}} C_{ox}$$

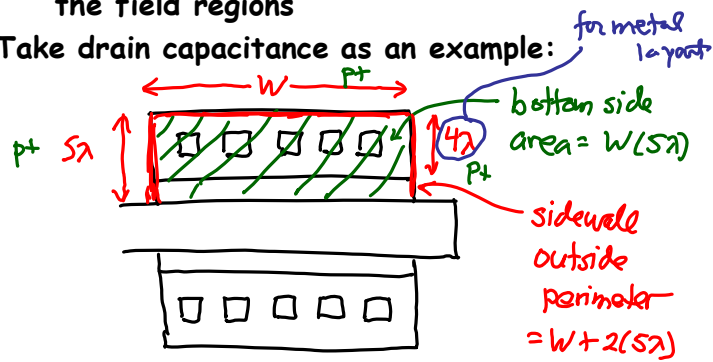
$W L_D C_{ox}$ \leftarrow accounts for the fact that the inversion charge is only under a portion of the gate

Gate-to-Drain Capacitance, C_{gd} :

$$C_{gd} = C_{OL} \quad (\text{no inversion charge near the drain in the saturation region})$$

\uparrow
 $W L_D C_{ox}$

- Source/Drain Junction Capacitance, C_{sb} & C_{db} :
- These are depletion capacitors associated with bulk-to-drain and bulk-to-source pn junctions
- Bottom capacitance per unit area differs from sidewall capacitance due to higher p+ bulk doping at the sidewalls
- The higher doping is near the silicon surface and designed to raise the threshold voltage in field areas, i.e., areas between transistors
 - \rightarrow This way unwanted inversion does not occur in the field regions
- Take drain capacitance as an example: *for metal layout*



$$C_{db} = \frac{C_{dbo}}{\sqrt{1 + \frac{V_{DS}}{\phi_j}}}$$

where:

$$C_{dbo}^{\Delta} = \text{depletion capacitance w/ } V_{SB} = 0V$$

$$= (\text{junction bottom-side area}) \times C_{j0}$$

$$+ (\text{junction outside perimeter}) \times C_{jsw}$$

$$= W(5\lambda) \times C_{j0} + (W + 2(5\lambda)) \times C_{jsw}$$

and where:

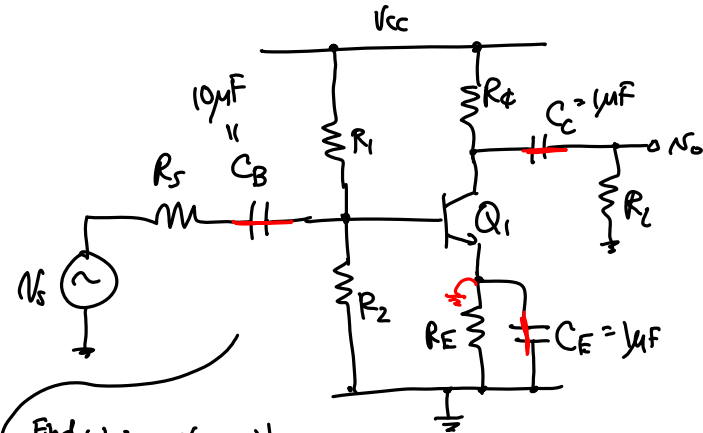
$$C_{j0} = \sqrt{\frac{q\epsilon_s N_B}{2\phi_j}}$$

bulk doping concentration

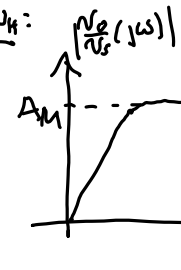
$$C_{jsw} = \sqrt{\frac{q\epsilon_s N_C}{2\phi_j}} \times x_j$$

channel-stop implant doping concentration
 S/D junction depth

C.E. High Freq Analysis (Brute Force)



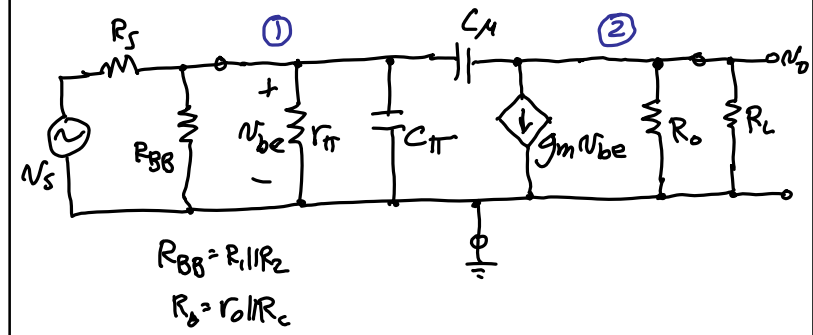
Find ω_H :



We want this frequency, which then gives us

$$\frac{N_o}{N_s}(s) = A_M \frac{1}{1 + \frac{s}{\omega_H}}$$

Small Signal Ckt.



KCL①: $\frac{V_s - V_{be}}{R_s} = \frac{V_{be}}{r_{\pi} \parallel R_{BB}} + V_{be}(sC_{\pi}) + (V_{be} - V_o)(sC_{\mu})$

KCL②: $(V_{be} - V_o)(sC_{\mu}) = g_m V_{be} + \frac{V_o}{R_o \parallel R_L} = g_m V_{be} + R''$
[$R'' = R_o \parallel R_L$]

Rearrange:

$$\frac{V_s}{R_s} = V_{be} \left[\frac{1}{R_s} + \frac{1}{r_{\pi} \parallel R_{BB}} + s(C_{\pi} + C_{\mu}) \right] - V_o(sC_{\mu})$$

$$\frac{1}{r_{\pi} \parallel R_s \parallel R_{BB}} = \frac{1}{R'}$$

...math...

$$\frac{V_o}{V_s}(s) = - \frac{g_m R' R''}{R_s} \left\{ \frac{(1 - s \frac{C_{\mu}}{g_m})}{1 + s[R'(C_{\pi} + C_{\mu}) + R''C_{\mu} + g_m C_{\mu} R' R''] + s^2 R' R'' C_{\pi} C_{\mu}} \right\}$$

A_M
 $= -g_m \frac{(r_{\pi} \parallel R_{BB})(R_o \parallel R_L \parallel R_L)}{R_s + r_{\pi} \parallel R_{BB}}$
 constant $\hat{=}$ midband gain

$F_H(s)$
 \uparrow
 models the freq. response

$F_H(s) = \frac{1 - s \frac{C_{\mu}}{g_m}}{1 + s[R'(C_{\pi} + C_{\mu}) + R''C_{\mu} + g_m C_{\mu} R' R''] + s^2 R' R'' C_{\pi} C_{\mu}}$
 $= \frac{1 - \frac{s}{z_1}}{(1 - \frac{s}{p_1})(1 - \frac{s}{p_2})} = \frac{1 + \frac{s}{\omega_z}}{(1 + \frac{s}{\omega_{p1}})(1 + \frac{s}{\omega_{p2}})} = \frac{N(s)}{D(s)}$

Zeros:

\Rightarrow identify a RHP zero: $z_1 = + \frac{g_m}{C_{\mu}} \rightarrow \omega_z = \frac{g_m}{C_{\mu}}$

Note that $\frac{g_m}{C_{\mu}} \gg \omega_T$ (since $\frac{g_m}{C_{\mu}} \gg \frac{g_m}{C_{\pi} + C_{\mu}}$ and $C_{\pi} \gg C_{\mu}$)

$\therefore \omega_z$ is a very high freq.!
 \hookrightarrow can ignore relative to the lower freq. poles

Poles:
 \Rightarrow often, $\omega_{p1} \ll \omega_{p2}$ (i.e., one pole frequency is much lower than other)
 (ω_{p1} dominates over ω_{p2})

$$D(s) = \left(1 + \frac{s}{\omega_{p1}}\right) \left(1 + \frac{s}{\omega_{p2}}\right)$$

$$= 1 + s \left(\frac{1}{\omega_{p1}} + \frac{1}{\omega_{p2}}\right) + \frac{s^2}{\omega_{p1}\omega_{p2}}$$

$$= 1 + \frac{s}{\omega_{p1}} + \frac{s^2}{\omega_{p1}\omega_{p2}} \quad \left. \vphantom{\frac{s^2}{\omega_{p1}\omega_{p2}}}\right\} \text{Dominant pole approximation}$$

Comparing w/ the denominator of $F_H(s)$:

$$\omega_{p1} = -p_1 \cong \frac{1}{R'(C_{\pi} + C_u) + R'_s C_u + g_m R'_s R''}$$

$$\omega_{p1} \cong \frac{1}{(C_{\pi} + C_u (1 + g_m R'' + \frac{R''}{R'})) R'}$$

$$\omega_{p1}\omega_{p2} \cong \frac{1}{R' R'' C_{\pi} C_u} \Rightarrow \omega_{p2} \cong \frac{[C_{\pi} + C_u (1 + g_m R'' + \frac{R''}{R'})] R'}{R' R'' C_{\pi} C_u}$$

$$\therefore \omega_{p2} = \frac{g_m}{C_{\pi}} + \frac{1}{R'' C_u} + \frac{1}{C_{\pi}} \left(\frac{1}{R'} + \frac{1}{R''}\right)$$

\uparrow

$$\frac{g_m}{C_{\pi}} = \omega_T (= \frac{g_m}{C_{\pi} C_u})$$

$\therefore \omega_{p2}$ is indeed at a very high freq. $\Rightarrow \omega_{p1}$
 \therefore satisfies our dominant pole approx. \checkmark

Take a closer look @ the form of ω_{p1} :

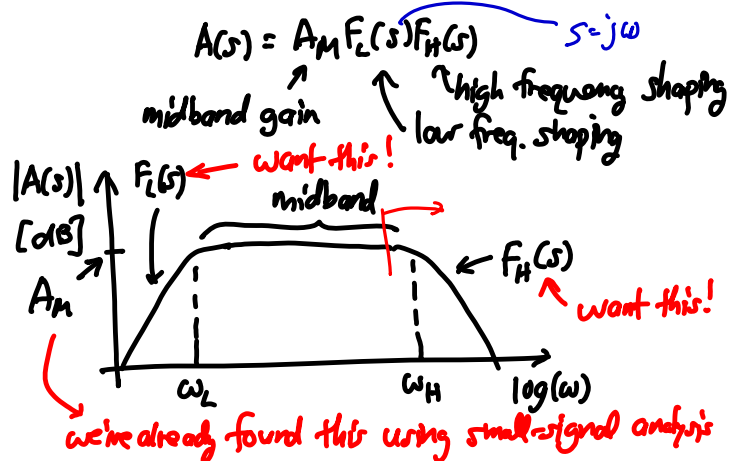
$$\omega_{p1} \cong \frac{1}{\underbrace{R' C_{\pi}}_{\tau_{\pi 0}} + \underbrace{(R' + R'' + g_m R' R'') C_u}_{\tau_{u 0} = R'' C_u}}$$

*

Can think of the denominator as the sum of time constants associated w/ each of the capacitor in the ckt!

Freq. Response

Recall that the transfer function of a general amplifier can be expressed as a function of frequency via:



High Freq. Response Determination Using Open Ckt. Time Constant (OCTC) Analysis

In general:

$$F_H(s) = \frac{1 + a_1 s + a_2 s^2 + \dots + a_{n_z} s^{n_z}}{1 + b_1 s + b_2 s^2 + \dots + b_{n_p} s^{n_p}}, \quad n_p > n_z$$

$$= \frac{\prod_{j=1}^{n_z} (1 - \frac{s}{z_j})}{\prod_{i=1}^{n_p} (1 - \frac{s}{p_i})} = \frac{\prod_{j=1}^{n_z} (1 + \frac{s}{\omega_{zj}})}{\prod_{i=1}^{n_p} (1 + \frac{s}{\omega_{pi}})}$$

from which:

$$b_1 = \frac{1}{\omega_{p1}} + \frac{1}{\omega_{p2}} + \dots + \frac{1}{\omega_{pn_p}} = \sum_{i=1}^{n_p} \frac{1}{\omega_{pi}} = \sum_{k=1}^{n_p} \tau_{pk}$$

\uparrow coeff. of the 1st order term

Through network theory, one can prove that: (see Gray & Meyer, Chpt. 7)

$$\sum_{i=1}^{n_p} \tau_{pi} = \sum_j C_j R_{j0} = \sum_j \tau_{j0} = b_1$$

where C_j are capacitors in the H.F. ckt., i.e., small ones
 $R_{j0} \hat{=}$ driving pt. resistance seen between the terminals of C_j determined with

- ① all small (<1nF) capacitors open-circuited
- ② all independent sources eliminated (i.e., short voltage sources, open current sources)
- ③ short all large (coupling/bypass) capacitors (i.e., >1μF or >1nF)

In calculating the H.F. response, we use the dominant pole approximation:

- (i) $\omega_{p1} \ll \omega_{p2}, \dots, \omega_{pn_p}$
- (ii) $F_H(s) \cong \frac{1}{1 + b_1 s} = \frac{1}{1 + \frac{s}{\omega_H}}$

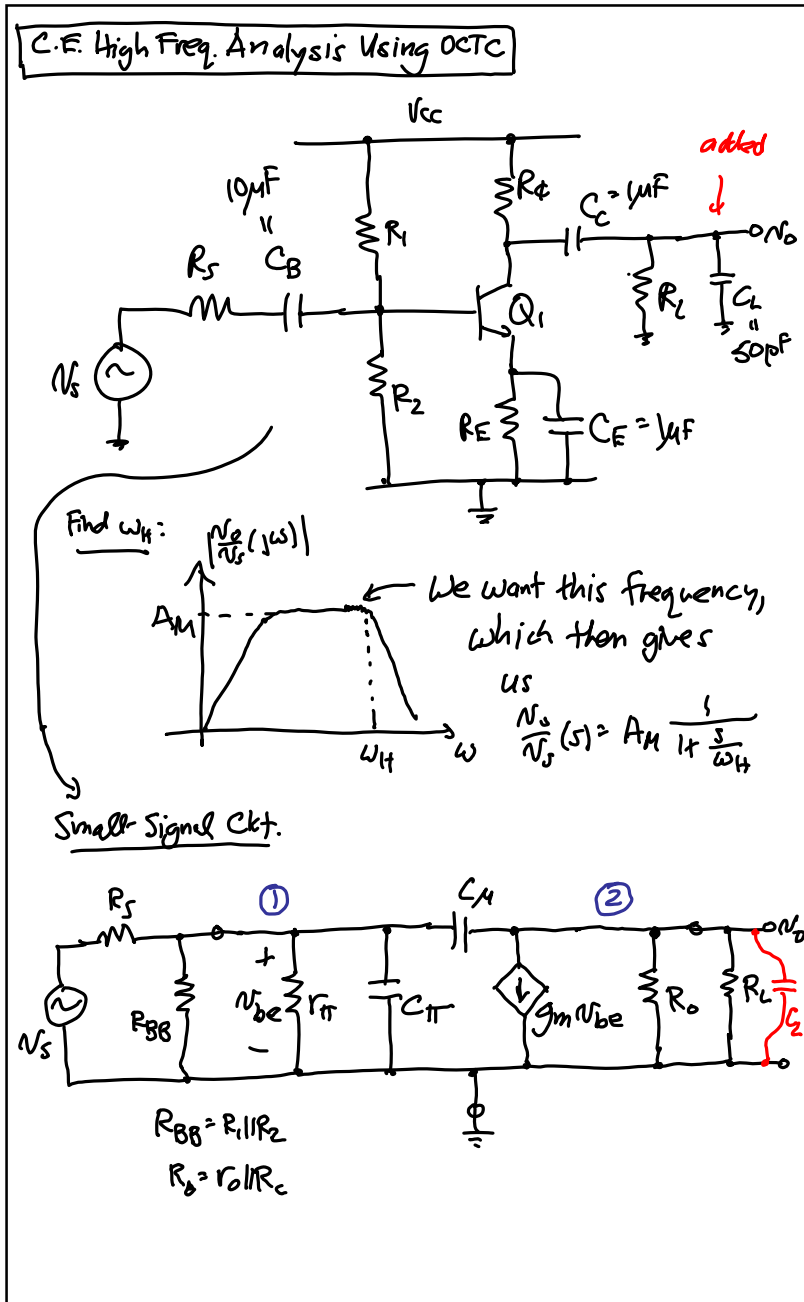
$$\omega_H \cong \frac{1}{b_1} = \frac{1}{\sum_j C_j R_{j0}} = \frac{1}{\sum_j \tau_{j0}}$$

When there is no dominant pole, an approximate expression for ω_H is:

$$\omega_H \approx \sqrt{\frac{1}{\omega_{p1}^2 + \omega_{p2}^2 + \dots} - \frac{1}{\omega_{z1}^2} - \frac{1}{\omega_{z2}^2} - \dots}$$

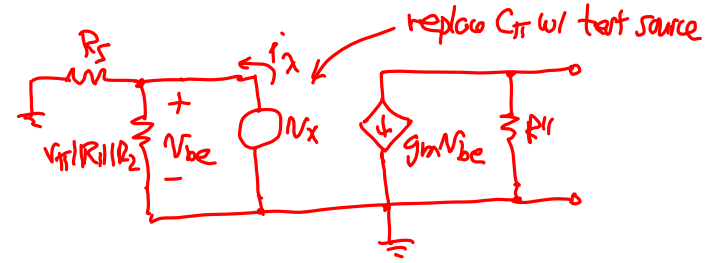
(just FYI)

• Now, use OCTC analysis on the CE amplifier to find the upper cut-off frequency, ω_H



(a) Determine $T_{\pi\pi} = C_{\pi} R_{\pi\pi}$

\Rightarrow open ckt. all C's, zero out all sources, and determine the driving point impedance, $R_{\pi\pi}$:

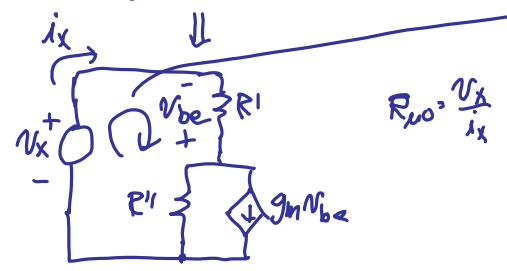
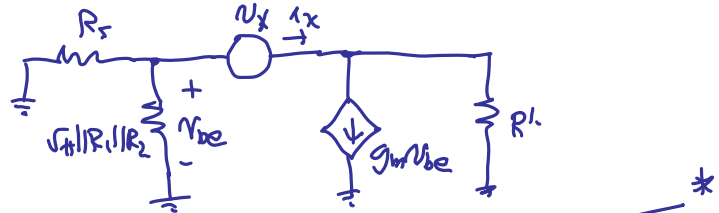


$R_{\pi\pi} = \frac{V_x}{i_x} = r_{\pi} || R_1 || R_2 || R_5 = R'$ (by inspection)

$\therefore T_{\pi\pi} = C_{\pi} R'$

(b) Determine $T_{\mu\mu} = C_{\mu} R_{\mu\mu}$:

\Rightarrow need $R_{\mu\mu}$: replace C_{μ} w/ test source

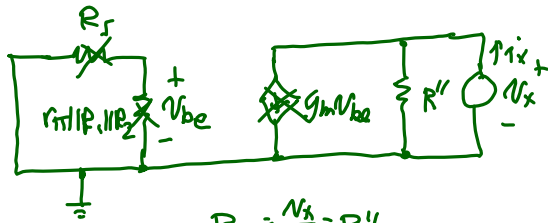


* KVL: $V_x = i_x R' + R''(i_x - g_m V_{be})$
 $[V_{be} = -i_x R'] \rightarrow V_x = i_x R' + R''(i_x + i_x R' g_m)$

$\therefore R_{in} = \frac{V_x}{i_x} = R' + R'' + g_m R' R''$

$\therefore \tau_{in} = C_{in} (R' + R'' + g_m R' R'')$

(c) Determine $\tau_{L0} = C_L R_{L0}$:



$R_{L0} = \frac{V_x}{i_x} = R''$

$\therefore \tau_{L0} = C_L R''$

Thus:

$$\omega_H = \frac{1}{\sum_j \tau_{j0}} = \frac{1}{C_{\pi} R' + C_{in} (R' + R'' + g_m R' R'') + C_L R''}$$

$= \dots *$