#### Lecture 2

#### **OUTLINE**

- Basic Semiconductor Physics (cont'd)
  - Carrier drift and diffusion
- PN Junction Diodes
  - Electrostatics
  - Capacitance

Reading: Chapter 2.1-2.2

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## **Dopant Compensation**

- An N-type semiconductor can be converted into P-type material by counter-doping it with acceptors such that  $N_{\rm A} > N_{\rm D}$ .
- A compensated semiconductor material has both acceptors and donors.

## Types of Charge in a Semiconductor

- Negative charges:
  - Conduction electrons (density = n)
- Ionized acceptor atoms (density = N<sub>△</sub>)
- Positive charges:
  - Holes (density = p)
  - Ionized donor atoms (density =  $N_D$ )
- The net charge density (C/cm³) in a semiconductor is

$$\rho = q(p - n + N_D - N_A)$$

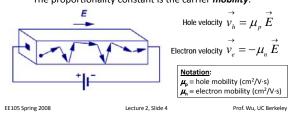
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#### **Carrier Drift**

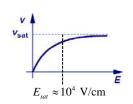
- The process in which charged particles move because of an electric field is called *drift*.
- Charged particles within a semiconductor move with an average velocity proportional to the electric field.
  - The proportionality constant is the carrier *mobility*.



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## **Velocity Saturation**

 In reality, carrier velocities saturate at an upper limit, called the saturation velocity (v<sub>sat</sub>).



$$\mu = \frac{\mu_0}{1 + bE}$$

$$v_{sat} = \frac{\mu_0}{b}$$

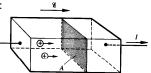
$$v = \frac{\mu_0}{1 + \frac{\mu_0 E}{v}} E$$

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#### **Drift Current**

• Drift current is proportional to the carrier velocity and carrier concentration:



 $v_h t A = volume from which all holes cross plane in time t$ 

 $p v_h t A = #$  of holes crossing plane in time t

 $q p v_h t A =$ charge crossing plane in time t

 $q p v_h A$  = charge crossing plane per unit time = hole current

ightharpoonup Hole current per unit area (i.e. current density)  $J_{\mathrm{p,drift}}$  =  $q~p~v_{\mathrm{h}}$ 

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## **Conductivity and Resistivity**

 In a semiconductor, both electrons and holes conduct current:

The current. 
$$J_{p,drift} = qp\mu_p E \qquad J_{n,drift} = -qn(-\mu_n E)$$

$$J_{tot,drift} = J_{p,drift} + J_{n,drift} = qp\mu_p E + qn\mu_n E$$

$$J_{tot,drift} = q(p\mu_p + n\mu_n) E \equiv \sigma E$$

Conductivity

 $\sigma \equiv qp\mu_p + qn\mu_n \quad \text{[unit: mho/cm = S/cm]}$ 

Resistivity

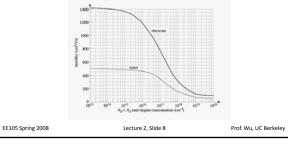
$$\rho = \frac{1}{\sigma} \quad \text{[Unit: } \Omega\text{-cm]}$$

• Typical resistivity range for Si:  $10^{-3} \sim 10^3 \ \Omega\text{-cm}$ 

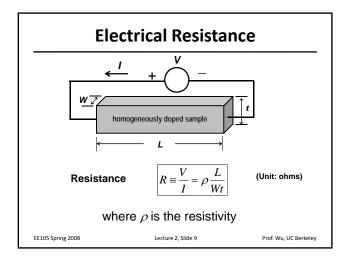
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# Resistivity Example the resistivity of a Si sample doped with

 Estimate the resistivity of a Si sample doped with phosphorus to a concentration of 10<sup>15</sup> cm<sup>-3</sup> and boron to a concentration of 10<sup>17</sup> cm<sup>-3</sup>. The electron mobility and hole mobility are 800 cm<sup>2</sup>/Vs and 300 cm<sup>2</sup>/Vs, respectively.

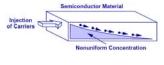


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- Due to thermally induced random motion, mobile particles tend to move from a region of high concentration to a region of low concentration.
  - Analogy: ink droplet in water
- Current flow due to mobile charge diffusion is proportional to the carrier concentration gradient.
  - The proportionality constant is the *diffusion constant*.



 $D_p \equiv$  hole diffusion constant (cm<sup>2</sup>/s)  $D_n \equiv$  electron diffusion constant (cm<sup>2</sup>/s)

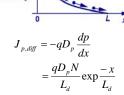
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# **Diffusion Examples**

- · Linear concentration profile → constant diffusion current

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Non-linear concentration profile → varying diffusion current



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## **Diffusion Current**

• Diffusion current within a semiconductor consists of hole and electron components:

$$\begin{split} J_{p,diff} &= -qD_p \frac{dp}{dx} \qquad J_{n,diff} = qD_n \frac{dn}{dx} \\ J_{tot,diff} &= q(D_n \frac{dn}{dx} - D_p \frac{dp}{dx}) \end{split}$$

The total current flowing in a semiconductor is the sum of drift current and diffusion current:

$$J_{tot} = J_{p,drift} + J_{n,drift} + J_{p,diff} + J_{n,diff}$$

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EE105 Fall 2007 3

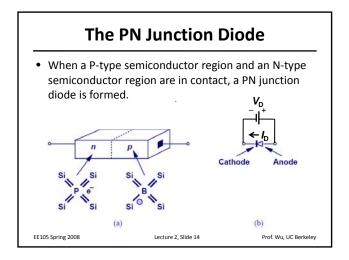
#### The Einstein Relation

 The characteristic constants for drift and diffusion are related:

$$\frac{D}{\mu} = \frac{kT}{q}$$

- Note that  $\frac{kT}{q}\cong 26 \mathrm{mV}$  at room temperature (300K)
  - This is often referred to as the "thermal voltage".

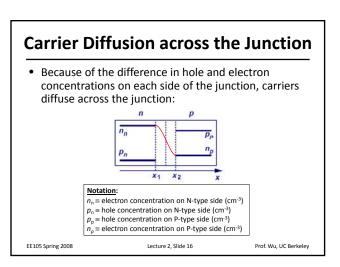
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#### **Diode Operating Regions** • In order to understand the operation of a diode, it is necessary to study its behavior in three operation regions: equilibrium, reverse bias, and forward bias. $V_{\rm D} = 0$ $V_{\rm D} < 0$ $V_{\rm D} > 0$ **PN Junction PN Junction PN Junction** in Equilibrium Under Reverse Bias **Under Forward Bias** Depletion Region Built-in Potential Junction Capacitance I/V Characteristics

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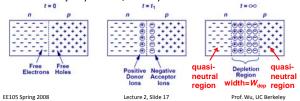


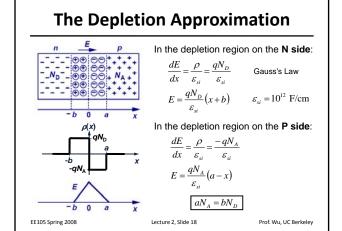
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## **Depletion Region**

- As conduction electrons and holes diffuse across the junction, they leave behind ionized dopants. Thus, a region that is depleted of mobile carriers is formed.
  - The charge density in the depletion region is not zero.
  - The carriers which diffuse across the junction recombine with majority carriers, i.e. they are annihilated.





## **Potential Distribution**

- In the depletion region, the electric potential is quadratic since the electric field is linear
- The potential difference between the N and the P side is called built-in potential, V<sub>0</sub>

$$E = -\frac{dV}{dx}$$

$$V = -\int E \cdot dx$$

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 $J_{tot} = J_{p,drift} + J_{n,drift} + J_{p,diff} + J_{n,diff} = 0$ 

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**PN Junction in Equilibrium** 

• In equilibrium, the drift and diffusion components of

current are balanced; therefore the net current

flowing across the junction is zero.

$$\begin{split} \boldsymbol{J}_{p,d\textit{rift}} &= -\boldsymbol{J}_{p,\textit{diff}} \\ \boldsymbol{J}_{n,\textit{drift}} &= -\boldsymbol{J}_{n,\textit{diff}} \end{split}$$

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## Built-in Potential, $V_0$

• Because of the electric field in the depletion region, there exists a potential drop across the junction:

$$\begin{split} qp\mu_p E &= qD_p \frac{dp}{dx} \quad \Rightarrow \quad p\mu_p \bigg( -\frac{dV}{dx} \bigg) = D_p \frac{dp}{dx} & \stackrel{-b}{\underset{p}{\longrightarrow}} \frac{a}{\underset{p}{\longrightarrow}} \\ \Rightarrow \quad -\mu_p \int\limits_{-b}^a dV = D_p \int\limits_{p_p}^{p_p} \frac{dp}{p} \\ \Rightarrow \quad V(-b) - V(a) &= \frac{D_p}{\mu_p} \ln \frac{p_p}{p_n} = \frac{kT}{q} \ln \frac{N_A}{\left(n_i^2 / N_D\right)} & \stackrel{\text{Depletion}}{\underset{\text{Region}}{\longleftarrow}} \\ \hline V_0 &= \frac{kT}{q} \ln \frac{N_A N_D}{n_i^2} & \text{(Unit: Volts)} \end{split}$$

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## **Built-In Potential Example**

• Estimate the built-in potential for PN junction below.

N P
$$N_D = 10^{18} \text{ cm}^{-3}$$
  $N_A = 10^{15} \text{ cm}^{-3}$ 

$$V_0 = \frac{kT}{q} \ln \left( \frac{N_D N_A}{n_i^2} \right) = (26 \text{mV}) \ln \left( \frac{10^{18} 10^{15}}{10^{20}} \right) = (26 \text{mV}) \ln \left( 10^{13} \right)$$

Note: 
$$\frac{kT}{q}\ln(10) \cong 26\text{mV} \times 2.3 \cong 60\text{mV}$$

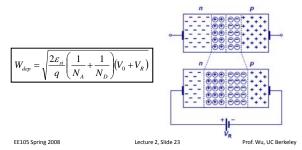
$$V_0 = 60 \,\text{mV} \times 13 = 780 \,\text{mV}$$

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#### **PN Junction under Reverse Bias**

 A reverse bias increases the potential drop across the junction. As a result, the magnitude of the electric field increases and the width of the depletion region widens.



#### **Diode Current under Reverse Bias**

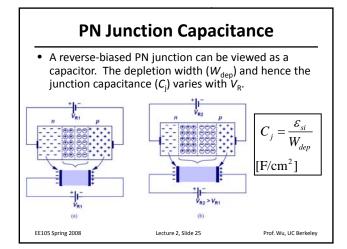
- In equilibrium, the built-in potential effectively prevents carriers from diffusing across the junction.
- Under reverse bias, the potential drop across the junction increases; therefore, negligible diffusion current flows. A very small drift current flows, limited by the rate at which minority carriers diffuse from the quasi-neutral regions into the depletion region.

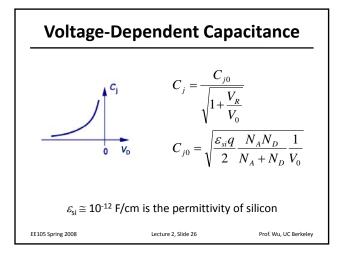
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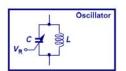
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## **Reverse-Biased Diode Application**

 A very important application of a reverse-biased PN junction is in a voltage controlled oscillator (VCO), which uses an LC tank. By changing V<sub>R</sub>, we can change C, which changes the oscillation frequency.



$$f_{res} = \frac{1}{2\pi} \frac{1}{\sqrt{LC}}$$

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## **Summary**

- Current flowing in a semiconductor is comprised of drift and diffusion components:  $J_{wa} = qp\mu_pE + qn\mu_nE + qD_n\frac{dn}{dx} qD_p\frac{dp}{dx}$
- A region depleted of mobile charge exists at the junction between P-type and N-type materials.
  - A built-in potential drop ( $V_{\rm 0}$ ) across this region is established by the charge density profile; it opposes diffusion of carriers across the junction. A reverse bias voltage serves to enhance the potential drop across the depletion region, resulting in very little (drift) current flowing across the junction.
  - The width of the depletion region ( $W_{\rm dep}$ ) is a function of the bias voltage ( $V_{\rm D}$ ).

 $W_{dep} = \sqrt{\frac{2\varepsilon_{si}}{a}} \left( \frac{1}{N_{\odot}} + \frac{1}{N_{\odot}} \right) (V_0 - V_D)$ 

 $V_0 = \frac{kT}{a} \ln \frac{N_A N_D}{n^2}$ 

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EE105 Fall 2007 7