

Ray propagation, the thick and thin lens, ABCD matrices (1)

Consider a spherical wave emanating from s, a point source

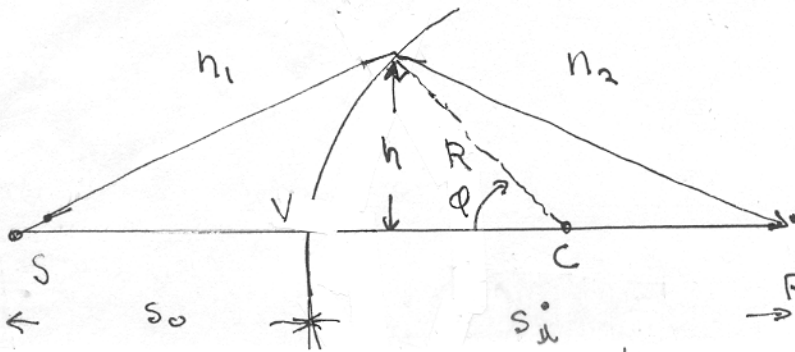


$$\text{Field} = \text{Re} \left(\frac{A}{r} e^{i(\omega t - k \cdot \vec{r})} \right)$$

ray vector $\frac{k \vec{r}}{k} = \hat{r}$

We have shown

$$\frac{n_1}{s_o} + \frac{n_2}{s_i} = \frac{n_2 - n_1}{R}$$



Signs For s_o, s_i (Hecht: 154)
 s_o is positive
 (s_o to the left of V)
 s_i is positive
 (s_i to the right of V)
 $R > 0$ (C to right of V) (see page 154 of Hecht)

Lens formula for rays

$$\frac{n_1}{s_o} + \frac{n_2}{s_i} = \frac{n_2 - n_1}{R}$$

A, B, C, D matrix for rays

$$\begin{pmatrix} h \\ s_o \end{pmatrix} = r_o' = \text{slope of o-ray}$$

$$- \begin{pmatrix} h \\ s_i \end{pmatrix} = r_i' = \text{slope of i-ray}$$

$h = r_o$ = location from axis of o ray

$h = r_i$ = location i ray

$$n_1 r_o' - n_2 r_i' = \frac{n_2 - n_1}{R} r_o$$

$$r_i = r_o$$

$$r_i' = - \frac{n_2 - n_1}{n_2 R} r_o + \frac{n_1}{n_2} r_o'$$

In terms of curvature $R > 0$ if the center is at $z' < z$

$$\frac{n_1}{R_1} - \frac{n_2}{R_2} = \frac{n_2 - n_1}{R n_2} \text{ or}$$

in terms of usual notation

$$\frac{n_1}{q_1} - \frac{n_2}{q_2} = \frac{n_2 - n_1}{R n_2}; \quad q_1 = R_1, \quad q_2 = R_2$$

(q will be the "Gaussian beam parameter but here it is just curvature R)

$$\frac{1}{q_2} = \frac{1}{q_1} \frac{n_1}{n_2} - \frac{n_2 - n_1}{n_2 R}$$

$$q_2 = \frac{q_1}{- \frac{n_2 - n_1}{n_1 R} q_1 + \frac{n_1}{n_2}}$$

Observe $q = \frac{A q_1 + B}{C q_1 + D}$

$$\begin{pmatrix} r_i \\ r_i' \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ - \frac{n_2 - n_1}{n_2 R} & \frac{n_1}{n_2} \end{pmatrix} \begin{pmatrix} r_o \\ r_o' \end{pmatrix} = \begin{pmatrix} A, B \\ C, D \end{pmatrix} \begin{pmatrix} r_o \\ r_o' \end{pmatrix}$$

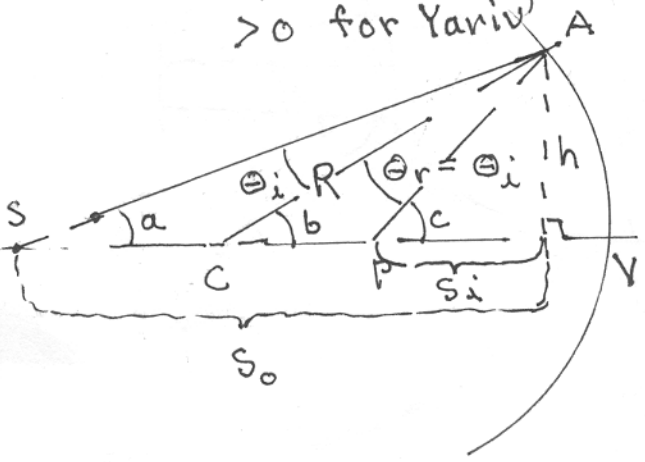
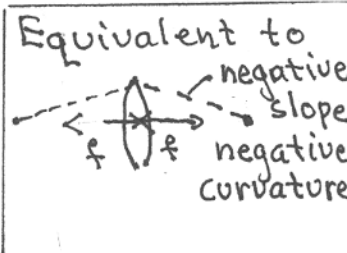
Agrees with Yariv page 41 (since $R > 0$ here)

2: The reflecting lens.

(2)

Note $R < 0$ for Hecht
(C to left of V) but
 > 0 for Yariv

page 182



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$$b = a + \theta_i$$

$$c = b + \theta_i$$

$$a + c = 2b$$

$$a = \frac{h}{S_0} ; b = \frac{h}{R} ; c = \frac{h}{S_i}$$

Thus $\frac{1}{S_0} + \frac{1}{S_i} = \frac{2}{R} = \frac{1}{f}$ + for both Hecht & Yariv

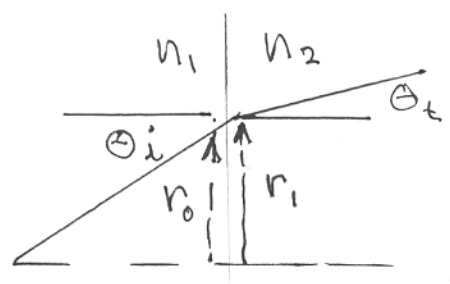
$S_0 = R_0$ $S_i = -R_i$ $R = \frac{R_0}{-\frac{1}{f}R_0 + 1}$

$(A \ B) = (1 \ 0)$
 $(C \ D) = (-\frac{1}{f} \ 1)$
(page 41, Yariv)

Also for ray and ray slopes

$$\begin{pmatrix} r_i \\ r_i' \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} r_0 \\ r_0' \end{pmatrix}$$

Planar boundary



$$r_1 = r_0$$

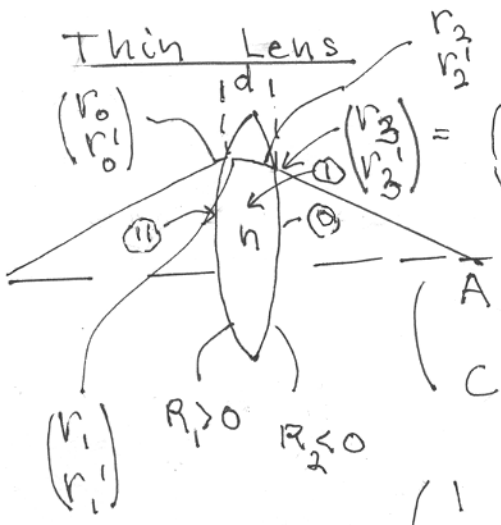
$$n_2 \theta_t = n_1 \theta_i$$

$$r_1' = \frac{n_1}{n_2} \theta_i = \frac{n_1}{n_2} r_0'$$

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{n_1}{n_2} \end{pmatrix}$$

(Yariv page 41)

Thin Lens



$$\begin{pmatrix} r_0 \\ r_0' \end{pmatrix} \begin{pmatrix} r_3 \\ r_3' \end{pmatrix} = \begin{pmatrix} A^0 & B^0 \\ C^0 & D^0 \end{pmatrix} \begin{pmatrix} r_2 \\ r_2' \end{pmatrix} = \begin{pmatrix} A^0 & B^0 \\ C^0 & D^0 \end{pmatrix} \begin{pmatrix} A' & B' \\ C' & D' \end{pmatrix} \begin{pmatrix} A'' & B'' \\ C'' & D'' \end{pmatrix} \begin{pmatrix} r_0 \\ r_0' \end{pmatrix}$$

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} \stackrel{n_1=n, n_2=1}{=} \begin{pmatrix} 1 & 0 \\ -\frac{(1-n)}{R_2}, n \end{pmatrix} \begin{pmatrix} 1 & d=0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{(n-1)}{nR_1}, \frac{1}{n} \end{pmatrix}$$

Assuming $d=0$ $n_2=n, n_1=1$

$$= \begin{pmatrix} 1 & 0 \\ -\frac{(1-n)}{R_2}, n \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{n-1}{nR_1}, \frac{1}{n} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -\frac{(n-1)}{R_1} - \frac{1}{R_2}, 1 \end{pmatrix}$$

$$\text{But } \begin{pmatrix} 1 & 0 \\ -(n-1)\left(\frac{1}{R_1} - \frac{1}{R_2}\right) & 1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix}$$

From the A, B, C, D Law.

$$q_1 = \frac{q_0}{-\frac{1}{f} q_0 + 1}$$

$$\frac{1}{q_1} = \frac{1}{q_0} \left(-\frac{1}{f} q_0 + 1 \right)$$

$$-\frac{1}{q_1} + \frac{1}{q_0} = \frac{1}{f}$$

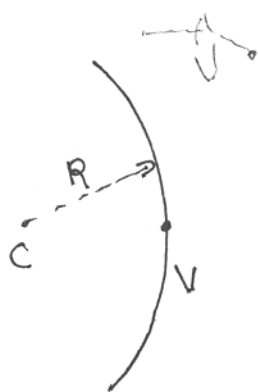
$$\frac{1}{f} = (n-1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

$$\uparrow \quad \uparrow \\ \frac{1}{(s_1)} + \frac{1}{(s_0)} = \frac{1}{f}$$

$$\lim_{s_0 \rightarrow \infty} s_i = f_i$$

$$\lim_{s_i \rightarrow \infty} s_0 = f_0$$

Sign Conventions for curved surfaces in Hecht (154)



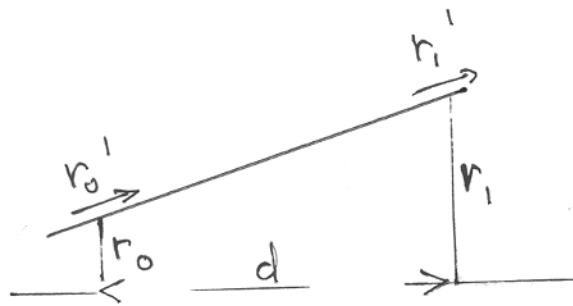
s_0, f_0	+ left of V	- <u>diverging</u> beam
x_0	+ left of V	
s_i, f_i	+ right of V	- <u>converging</u> beam
x_i	+ right of V	
R	+ if C is right of V	- <u>conver</u>

For beam propagation + curvature is diverging - curvature is converging

Free space

(4)

Ray



$$r_1 = r_0 + r_0' d$$

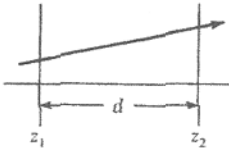
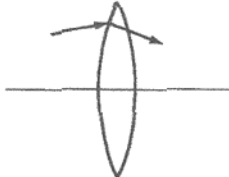
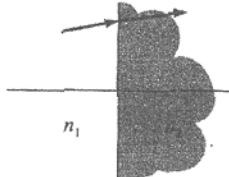
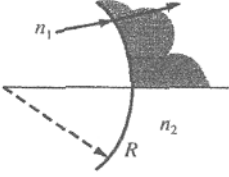
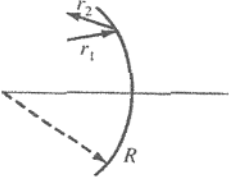
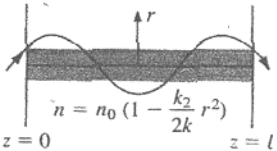
$$r_1' = r_0'$$

$$\begin{aligned} \begin{pmatrix} r_1 \\ r_1' \end{pmatrix} &= \begin{pmatrix} 1 & d \\ 0 & 1 \end{pmatrix} \begin{pmatrix} r_0 \\ r_0' \end{pmatrix} \\ &= \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} r_0 \\ r_0' \end{pmatrix} \end{aligned}$$

Gaussian parameter

$$q_1 = \frac{A q_0 + B}{C q_0 + D} = q_0 + d$$

Table 2-1 Ray Matrices for Some Common Optical Elements and Media

<p>(1) Straight Section: Length d</p>		$\begin{bmatrix} 1 & d \\ 0 & 1 \end{bmatrix}$
<p>(2) Thin Lens: Focal length f ($f > 0$, converging; $f < 0$, diverging)</p>		$\begin{bmatrix} 1 & 0 \\ -1/f & 1 \end{bmatrix}$
<p>(3) Dielectric Interface: Refractive indices n_1, n_2</p>		$\begin{bmatrix} 1 & 0 \\ 0 & n_1/n_2 \end{bmatrix}$
<p>(4) Spherical Dielectric Interface: Radius R</p>		$\begin{bmatrix} 1 & 0 \\ \frac{n_2 - n_1}{n_2 R} & \frac{n_1}{n_2} \end{bmatrix}$
<p>(5) Spherical Mirror: Radius of curvature R</p>		$\begin{bmatrix} 1 & 0 \\ -2/R & 1 \end{bmatrix}$
<p>(6) A medium with a quadratic index profile</p>		$\begin{bmatrix} \cos\left(\sqrt{\frac{k_2}{k}} l\right) & \sqrt{\frac{k}{k_2}} \sin\left(\sqrt{\frac{k_2}{k}} l\right) \\ \sqrt{\frac{k_2}{k}} \sin\left(\sqrt{\frac{k_2}{k}} l\right) & \cos\left(\sqrt{\frac{k_2}{k}} l\right) \end{bmatrix}$

See Chapt 6 of Hecht