

EECS 16B

Designing Information Devices and Systems II

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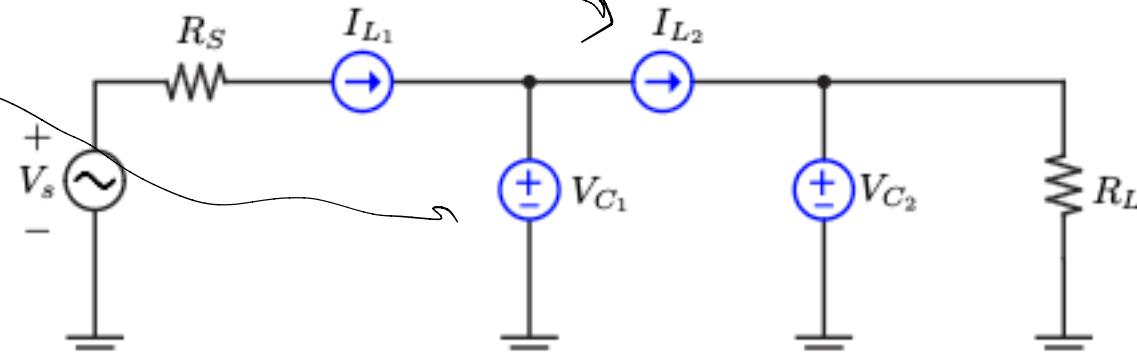
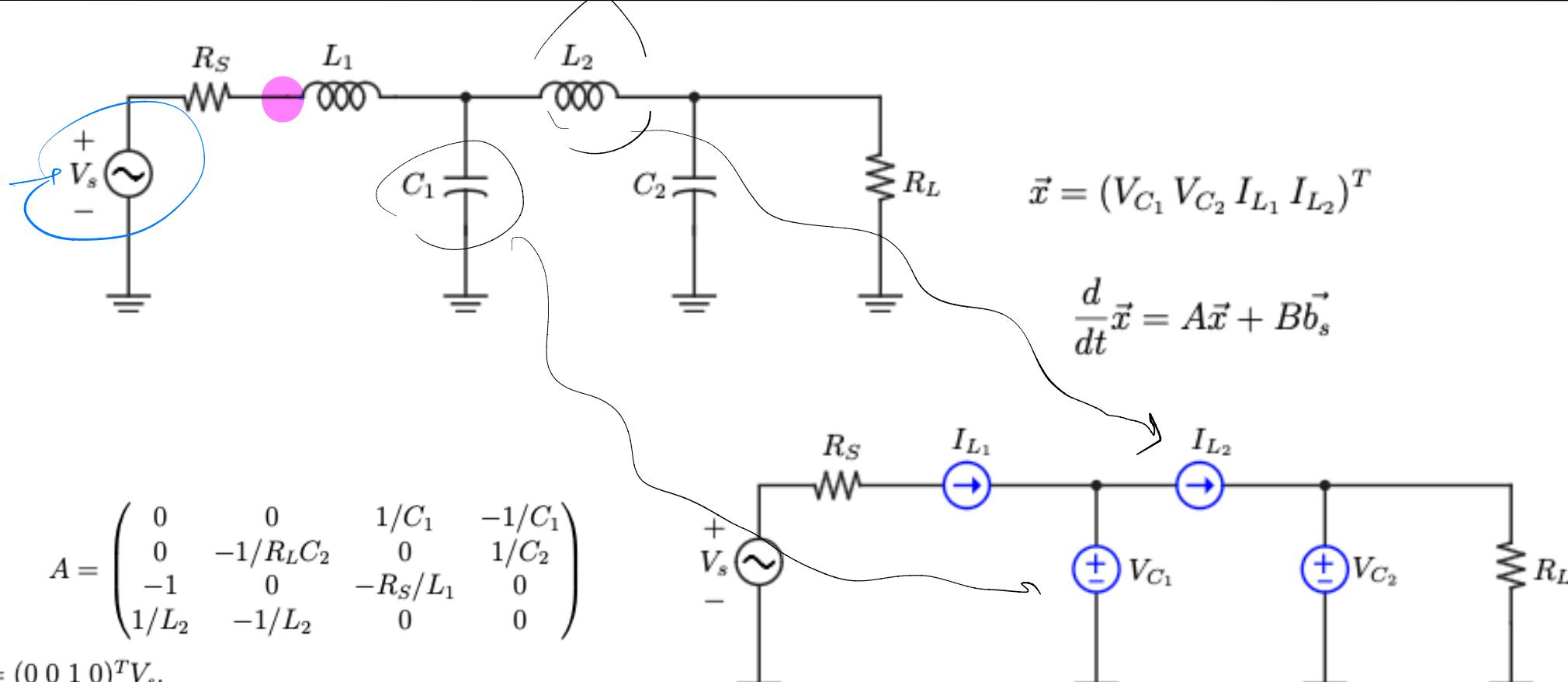
Module 6: Sinusoidal Steady-State Solution to Vector Differential Equations

EECS 16B

Summary

- Solution of VDE in eigenspace with sources
- Steady-state solution of VDE
- Concept of Impedance
- AC Circuits
- Examples

Example Circuit



$$\frac{d}{dt} \vec{x} = A \vec{x} + B \vec{b}_s \quad \text{Ind. sources}$$

State vector

Numerical Example

$$L1 = 7410^{\wedge} - 9; L2 = 114.210^{\wedge} - 9; C1 = 45.61$$

Let's see the eigenvalues for this matrix:

Eigenvalues[A]/MatrixForm

$$\lambda_n \quad \left(\begin{array}{c} -6.75676 \times 10^8 + 0.i \\ -1.64354 \times 10^8 + 5.88916 \times 10^8 i \\ -1.64354 \times 10^8 - 5.88916 \times 10^8 i \\ -3.47197 \times 10^8 + 0.i \end{array} \right)$$

$\lambda_{n+1} = \lambda_n^*$

Re¹ 4 distinct eigenvalues

complex conjugate

Re²

Stability

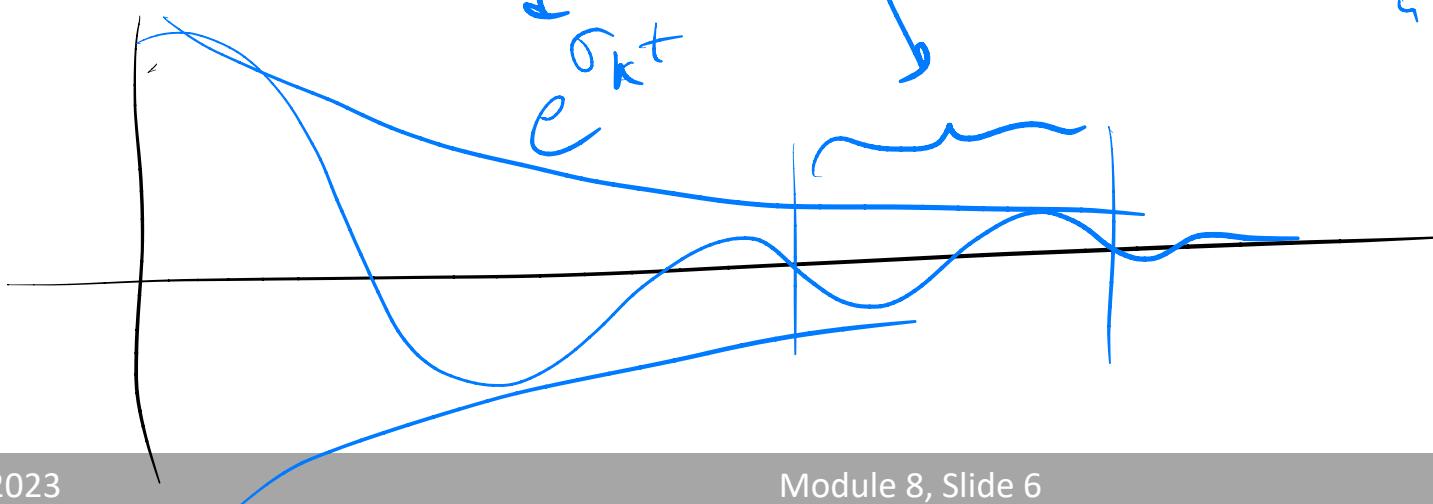
$$\operatorname{Re}(\lambda_n) < 0$$

$$q_k(t) = q_{k_0} e^{\sigma_k t}$$

$$q_k(t) = q_{k_0} e^{\sigma_{n_t} t} e^{j\omega_n t}$$

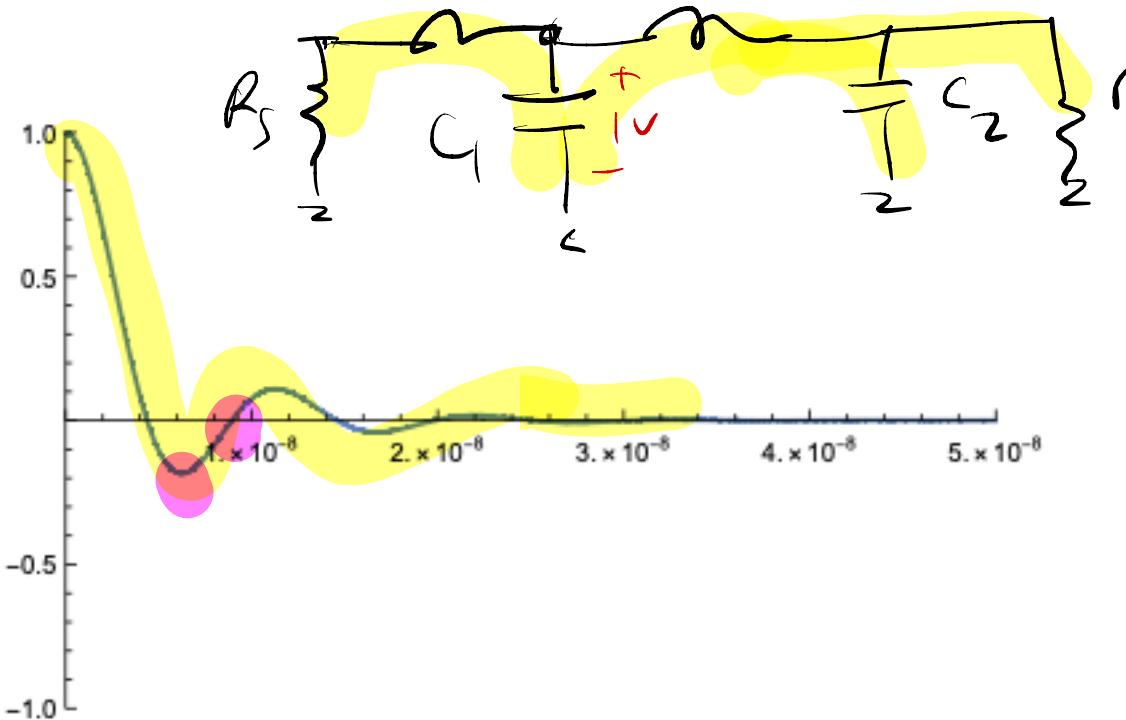
$$\lambda_n = \sigma_n + j\omega_n$$

$$= q_{k_0} e^{\sigma_n t} (\cos \omega_n t + j \sin \omega_n t)$$



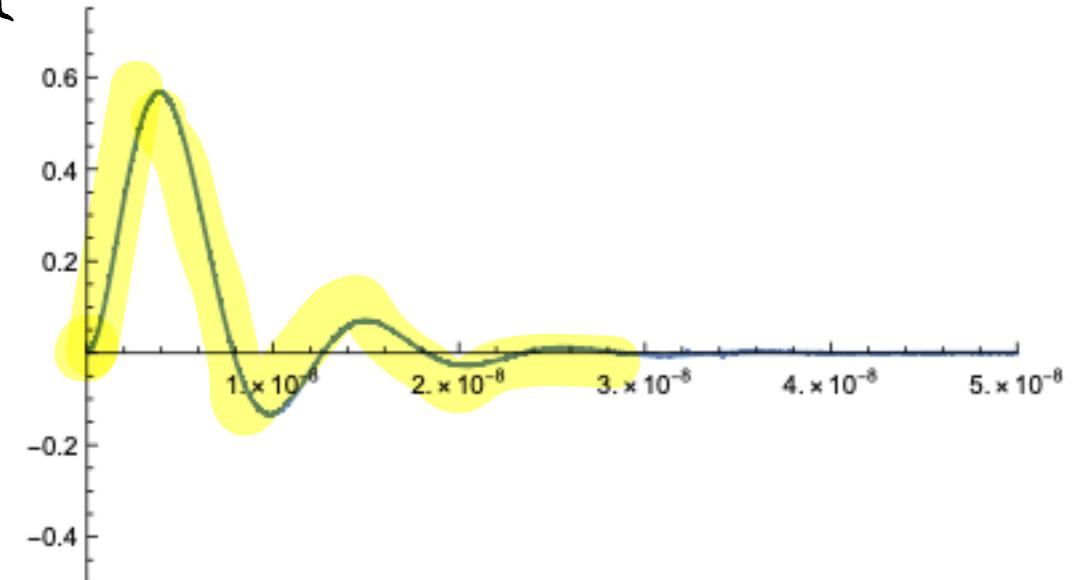
"natural freq"
 ω_n
decay rate: σ_n

Homogeneous Solution



State 1

$$x_1 = v_{C_1}$$
$$\vec{x}_0 = (1, 0, 0, 0)^T$$



State 2

$$x_2 = v_{C_2}$$

Forced Response

- Since we've reduced the system to simple constant coefficient linear differential equations, we know the forced response !
- Note that the sources are also transformed into the eigenspace through the matrix Q^{-1}

$$\vec{q}_k^h(t) = \vec{q}_k^h(0) e^{\lambda_k t}$$
$$\vec{q}(t) = \vec{q}(0) e^{\Lambda t}$$

Euler const Matrix

Notation : $e^{\Lambda} = \begin{bmatrix} e^{\lambda_1 t} & 0 \\ 0 & e^{\lambda_n t} \end{bmatrix}$

Λ : diagonal matrix

$$\vec{x}(t) = Q \vec{q}(t)$$

Eigenvector matrix

$$\vec{q}(t) = \vec{q}(0) e^{\lambda t}$$

$$\vec{x}(t) = Q \vec{q}(t) = \underbrace{Q \vec{q}(0)}_? e^{\lambda t}$$

$$\vec{q}(0) = Q^{-1} \vec{x}(0)$$

initial state

transformation
matrix

Steady-State DC Response

$$\cancel{\frac{d\vec{x}}{dt}} = A\vec{x} + \vec{b}_s$$

$$\lim_{t \rightarrow \infty} e^{\lambda t} = 0$$

All eigen values have negative real part

$$0 = A\vec{x} + \vec{b}_s$$

$$\vec{x} = -A^{-1}\vec{b}_s$$
$$= -A^{-1}B\vec{b}_s$$

Review: General Solution to 1st-order ODE

- With constant coefficients, we derived

$$\frac{d\vec{x}}{dt} = \underbrace{A\vec{x}}_{\text{Vector of states}} + \underbrace{B\vec{b}_s}_{\text{Vector of sources}}$$

$$\frac{d\vec{x}}{dt} = \underbrace{A\vec{x}}_{\text{Vector of states}} + \underbrace{B\vec{b}_s}_{\text{Vector of sources}}$$

$$\frac{d\vec{x}}{dt} = Q^{-1}Q^T\vec{x} + B\vec{b}_s$$

$$\frac{d(Q^T\vec{x})}{dt} = -\omega Q^T\vec{x} + Q^T B \vec{b}_s$$

$$\boxed{\frac{d\vec{q}}{dt} = -\omega \vec{q} + \vec{b}_s}$$

Apply Solution to VDE

$$\frac{dq_k}{dt} = \lambda_k q_k + \tilde{b}_{sk}$$

$$e^{-\lambda_k t} \left(\frac{dq_k}{dt} - \lambda_k q_k \right) = \tilde{b}_{sk} e^{-\lambda_k t}$$

$$\left(e^{-\lambda_k t} \frac{dq_k}{dt} - \lambda_k e^{-\lambda_k t} q_k \right)$$

$$\int (q_k(t) e^{-\lambda_k t})' ds = \int_{-\infty}^t \tilde{b}_{sk} e^{-\lambda_k s} ds + k_k$$

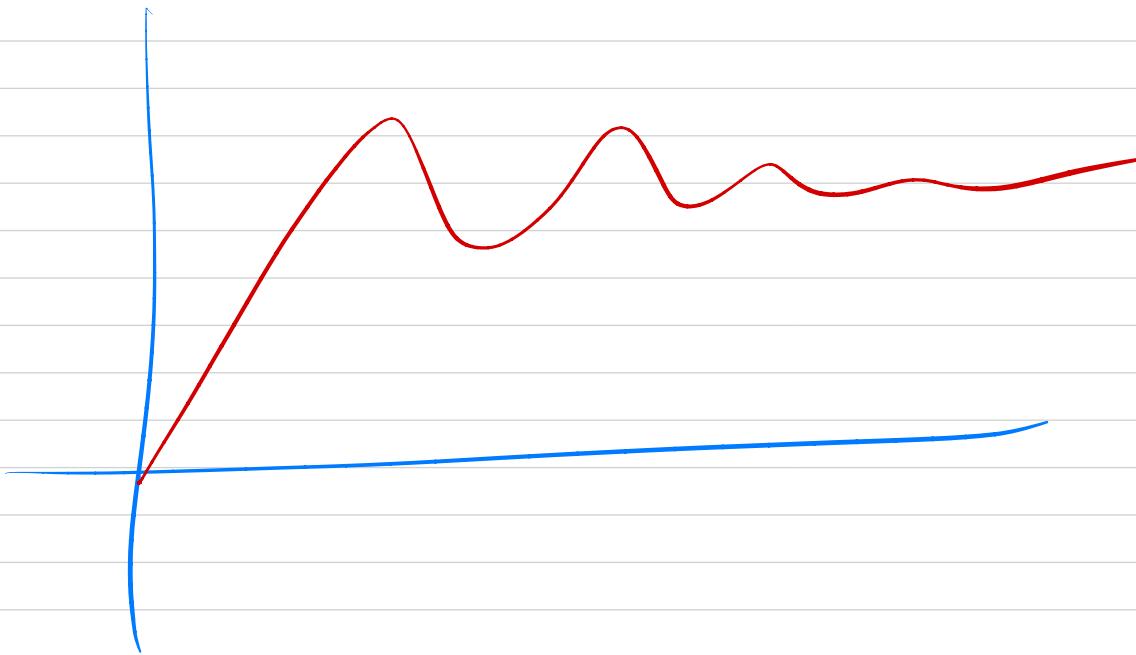
$$\int \left(q_k(t) e^{-\lambda_k t} \right)' ds = \int_{-\infty}^t \tilde{b}_{sk} e^{-\lambda_k s} ds + k_k$$

$$q_k(t) = e^{\lambda_k t} \underbrace{\int_{-\infty}^t \tilde{b}_{sk}(s) e^{-\lambda_k s} ds}_{\text{forced solution}} + e^{\lambda_k t} \underbrace{q_k(0)}_{\text{solution to zero input}}$$

for zero initial

t

$$\vec{q}(t) = e^{\lambda_k t} \xrightarrow{\text{condition}} \vec{q}(0) + e^{\lambda_k t} \int_{-\infty}^t Q^{-1} \tilde{b}_s e^{-\lambda_k s} ds$$



$$\vec{q}(t) = e^{-\lambda t} \overset{\text{def}}{\vec{q}(0)} + e^{-\lambda t} \int_{-\infty}^t Q^{-1} \vec{b}_s e^{-\lambda s} ds$$

$$\vec{x}(t) = Q \vec{q}(t) = \underbrace{Q e^{-\lambda t} Q^{-1}}_{\text{A curve}} \vec{x}(0) +$$

$$Q e^{-\lambda t} \int_{-\infty}^t Q^{-1} \vec{b}_s e^{-\lambda s} ds$$

$$Q e^{-\lambda t} \int_{-\infty}^t \vec{b}_s e^{-\lambda s} ds$$

$$\vec{x}(t) = \underbrace{Q e^{-\lambda t} \vec{x}(0)}_{C^{At}} + \underbrace{Q e^{-\lambda t} \int_{-\infty}^t \tilde{b}_s e^{-\lambda s} ds}_{\text{integral term}}$$

$$\frac{dx}{dt} = ax + b$$

$$\vec{x}(t) = e^{At} \vec{x}(0) + e^{At} \int_{-\infty}^t \tilde{b}_s e^{-\lambda s} ds$$

Elegant Notation

- Let's define the matrix exponential
- While for now you can accept this as merely a convenient notation, it turns out to be a more general idea !
[\(https://youtu.be/O85OWBJ2ayo\)](https://youtu.be/O85OWBJ2ayo)

$$e^{At} \triangleq \underbrace{Q e^{\Lambda t} Q^{-1}}$$

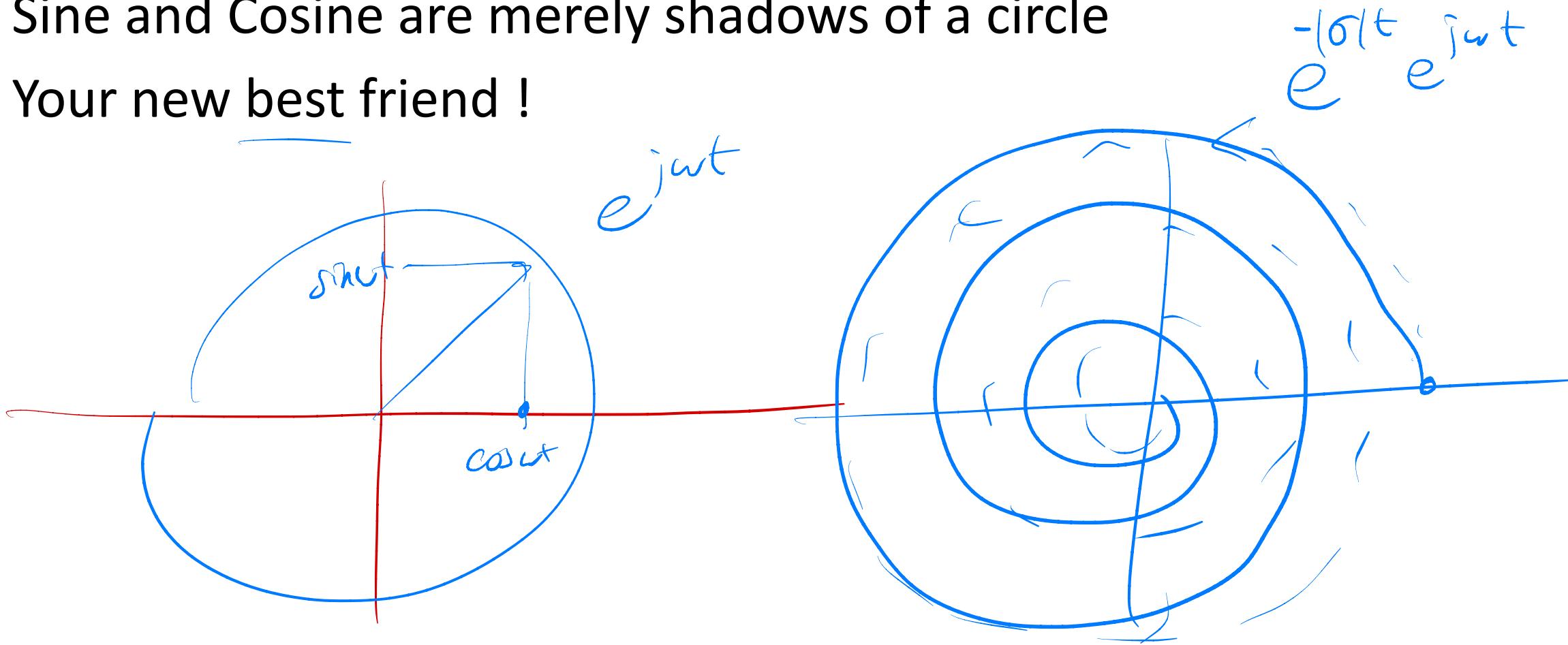
$$e^x = 1 + \frac{x}{1!} + \frac{x^2}{2!} + \dots$$

Forced Response in Steady-State

- Assume source is sinusoidal (very common in practice)
 - RF signals essentially sinusoidal on short time scales
 - Musical notes are superpositions of tones (frequencies)
 - Any general waveform can be decomposed into sinusoids ! (Fourier Series and Transform ... EE120)

Complex Exponential (again)

- Sine and Cosine are merely shadows of a circle
- Your new best friend !



Steady-State Continued

$$q_a(t) = e^{\lambda_a t} q_a(0) + e^{\lambda_a t} \int_{-\infty}^t e^{-\lambda_a s} b_{sa}(\omega) ds$$

steady state $\Omega \rightarrow t=\infty$

$$q_a^f(t) = e^{\lambda_a t} \int_{-\infty}^t e^{-\lambda_a s} Q^{-1} \hat{b}_s e^{j\omega s} ds$$

$$\hat{b}_{sa} = Q^{-1} b_s$$

$$= Q^{-1} \hat{b}_s e^{j\omega t}$$

$$= e^{\lambda_a t} Q^{-1} \hat{b}_s \int_{-\infty}^t e^{(j\omega - \lambda_a)s} ds$$

$$= e^{\lambda_a t} Q^{-1} \hat{b}_s \left\{ \frac{e^{(j\omega - \lambda_a)t}}{j\omega - \lambda_a} \right\}_{-\infty}^t = \frac{e^{(j\omega - \lambda_a)t}}{j\omega - \lambda_a} - \cancel{\left[\frac{e^{(j\omega - \lambda_a)t}}{j\omega - \lambda_a} \right]_{t \rightarrow \infty}}$$

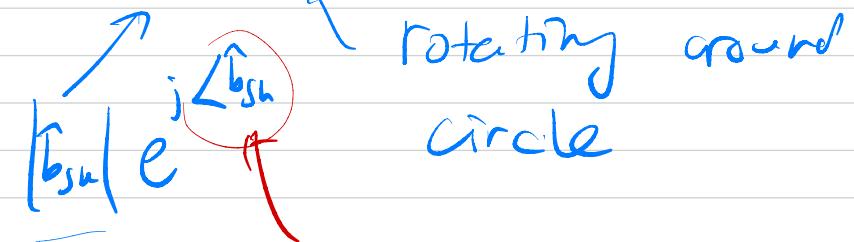
by $\Re(j\omega - \lambda_a) < 0$ stable

~~$\frac{e^{(j\omega + \lambda_a)t}}{j\omega + \lambda_a}$~~

$$\begin{aligned}
 \hat{b}_{sh} &= b_1 e^{j(\omega t + \phi_1)} + b_2 e^{j(\omega t + \phi_2)} + b_3 e^{j(\omega t + \phi_3)} \\
 &= b_1 e^{j\omega t} e^{j\phi_1} + b_2 e^{j\omega t} e^{j\phi_2} + b_3 e^{j\omega t} e^{j\phi_3} \\
 &= e^{j\omega t} (b_1 e^{j\phi_1} + b_2 e^{j\phi_2} + \dots) \\
 &\quad (b_1 \cos \phi_1 + b_1 j \sin \phi_1 + b_2 \cos \phi_2 + j b_2 \sin \phi_2 + \dots)
 \end{aligned}$$

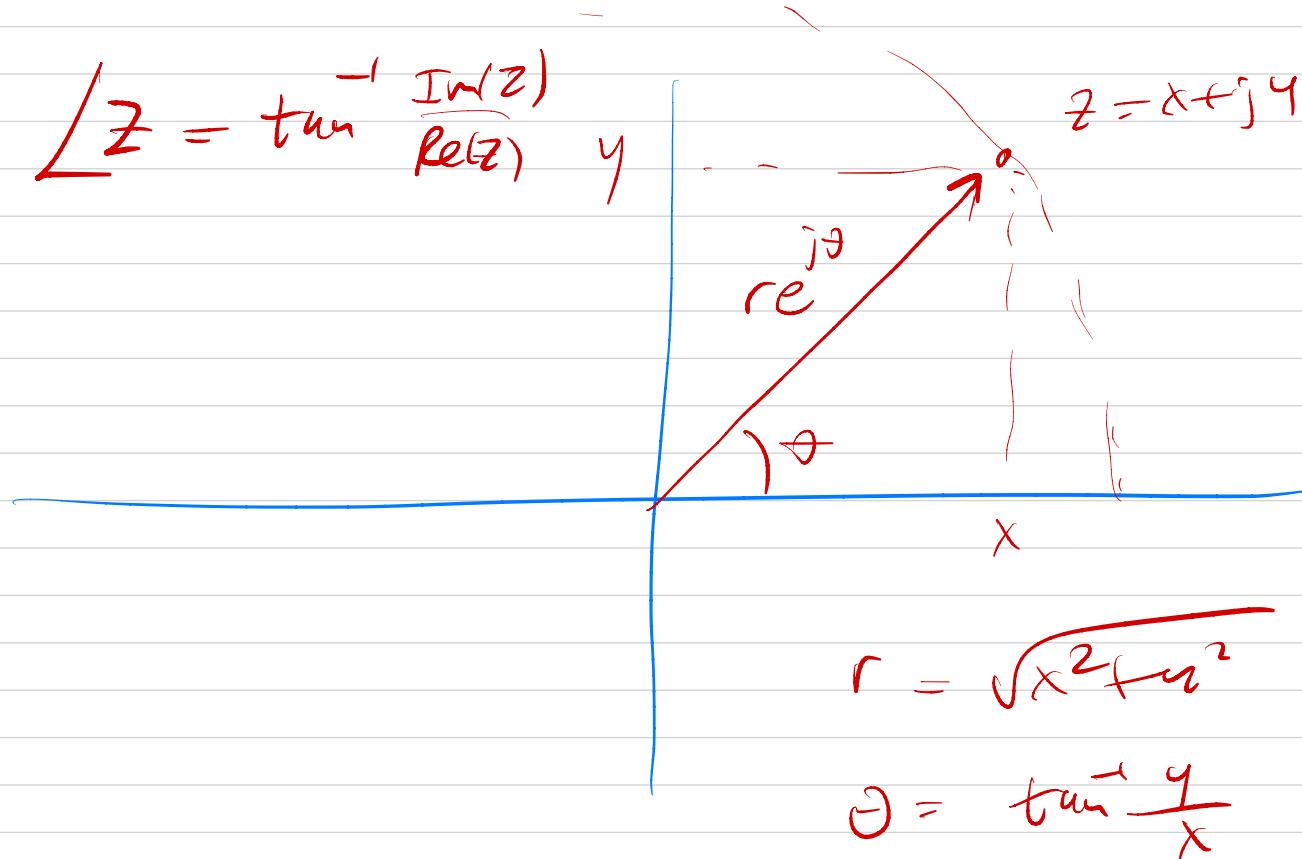
$$= x + jy = r e^{j\theta}$$

$$= \hat{b}_{sh} e^{j\omega t}$$

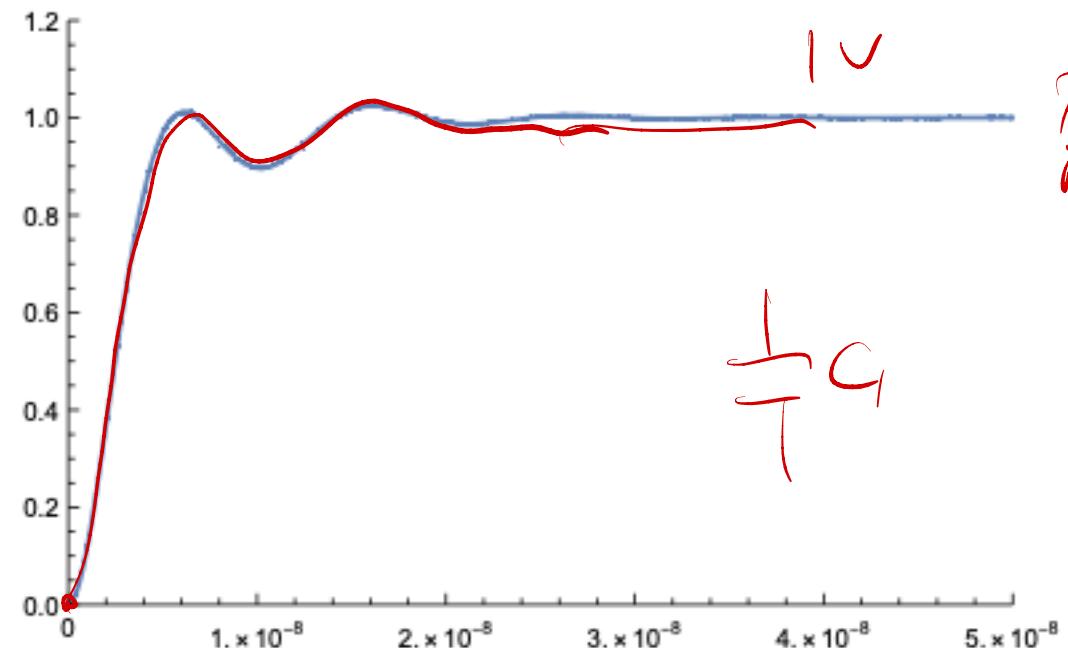
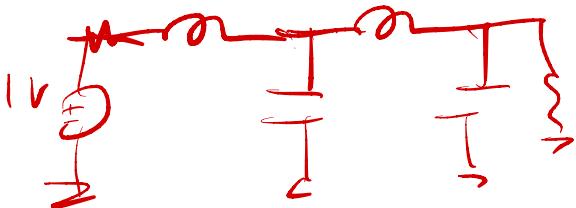
$\hat{b}_{sh} = |\hat{b}_{sh}| e^{j\angle \hat{b}_{sh}}$ 

 rotating around circle

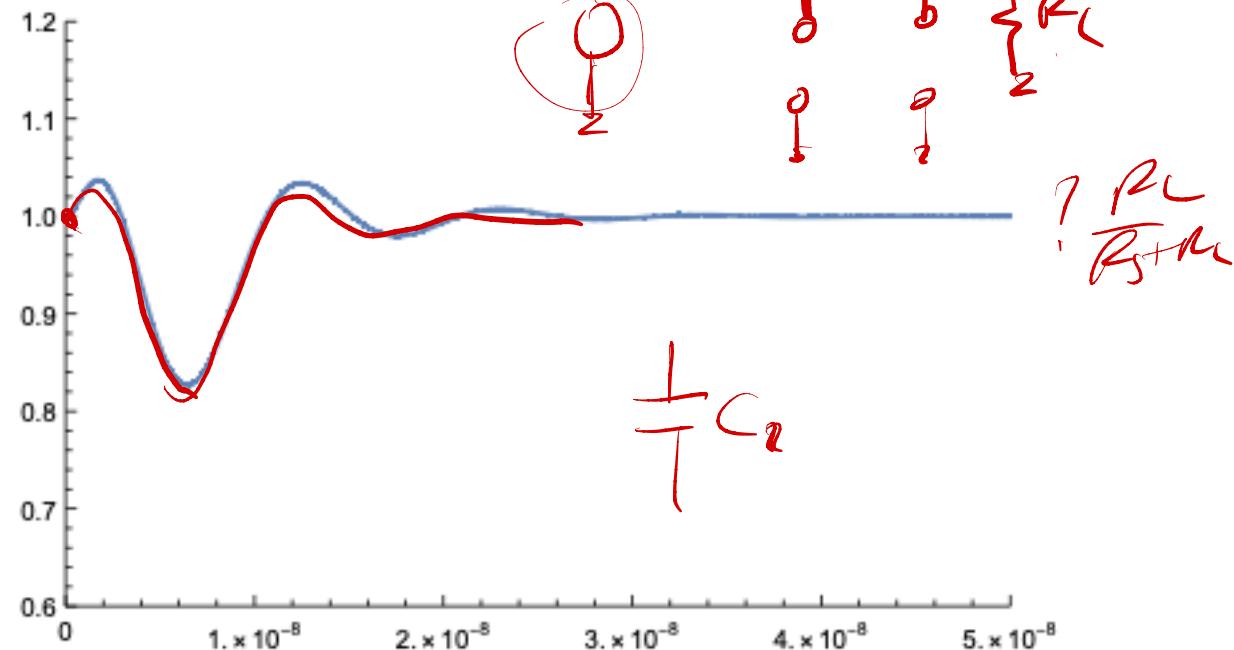
$$\angle z = \tan^{-1} \frac{\operatorname{Im}(z)}{\operatorname{Re}(z)}$$



Example: Forced Solution



State 1



State 2

$$\vec{x}(0) = (0, 1, 0, 0)^T$$

Steady-State

$$q_k(t) = \frac{\hat{b}_k}{\underbrace{jw - \lambda_k}_{+}} e^{jw t} = q_k e^{jw t}$$

+ some complex #

$$\vec{x} = Q \vec{q} \neq$$

$$x_k(t) = \left(\sum_u \hat{h}_k \right) e^{jw t}$$

$$x_k(t) = h_k e^{jw t}$$