



The background image shows a detailed microchip layout with various functional blocks highlighted by dashed yellow boxes. The labels include 'RX' (Receiver), 'LO Buffer' (Local Oscillator Buffer), 'Hybrid', 'Wilkinson' (referring to a Wilkinson power divider), 'LO Buffer' (another instance), and 'TX' (Transmitter). The layout features a dense grid of circuit traces and components.

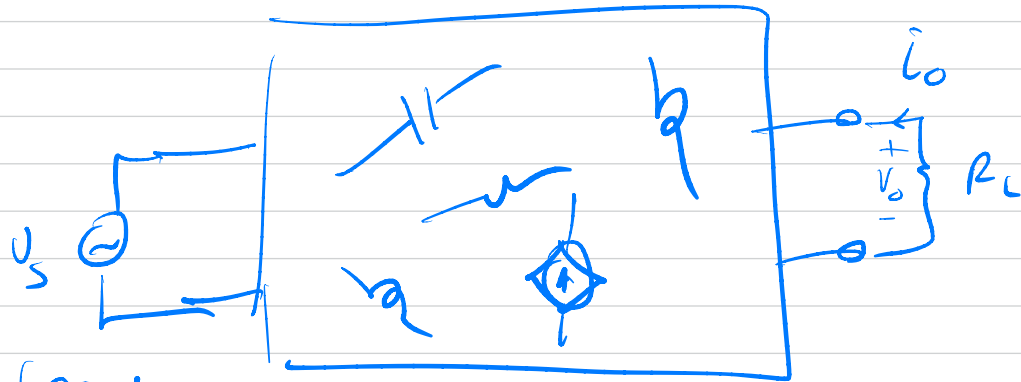
EECS 16B

Designing Information Devices and Systems II

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Module 9: AC Analysis

EECS 16B



v_s
input

Any generic circuit

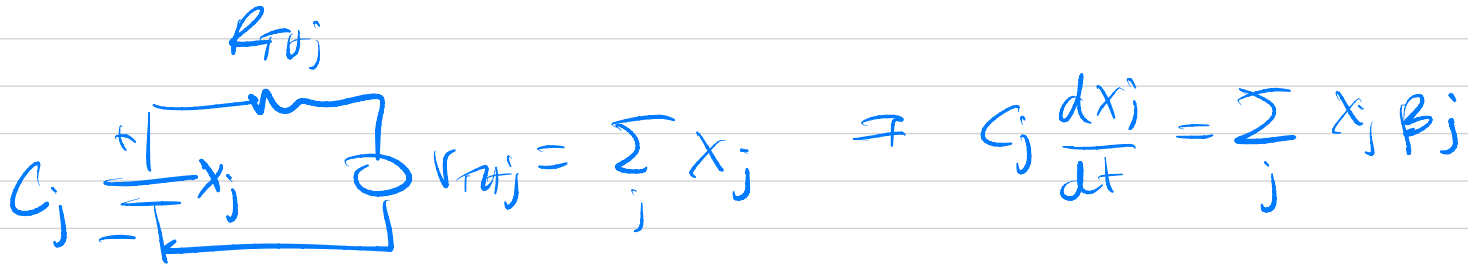
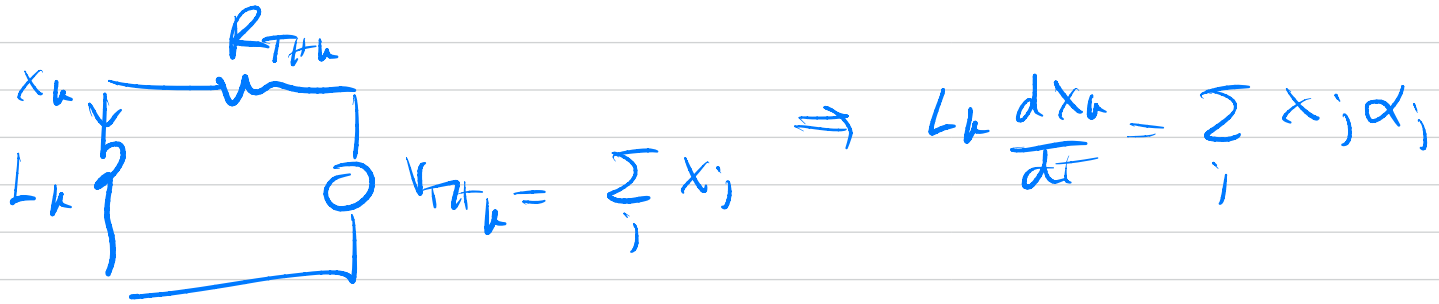
L, R, C

dependent source

$$\Rightarrow \frac{d\vec{x}}{dt} = A\vec{x} + B\vec{k}_s$$

\vec{x} : state

$\left\{ \begin{array}{l} \text{voltages on caps} \\ \text{currents thru ind} \end{array} \right\}$



⋮

$$\frac{d\vec{x}}{dt} = A\vec{x} + B\vec{b}_s$$

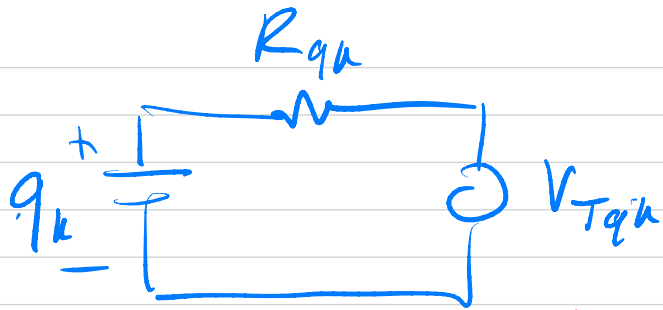
$$\frac{d\vec{x}}{dt} = Q \underbrace{\Lambda Q^{-1}} \vec{x} + B\vec{b}_s$$

$$\frac{d\vec{q}}{dt} = \Lambda \vec{q} + Q^{-1}\vec{b}_s$$

\vec{x} : state space

\vec{q} : state space in the eigenbasis

q_a^i modes of the circuit



eq. circuit
of each row ←

NOT NEEDED ↗

JUST a reminder that we've
problem before.

$$q_u(t) = e^{\lambda_u t} q(0) + e^{\lambda_u t} \int_0^t e^{-\lambda_u s} \tilde{b}_{su}(s) ds$$

General Solution

$$b_s(t) = e^{j\omega t}$$

$\tilde{b}_s(t)$ = source in the eigenbasis

\Rightarrow sum of any number of

$$a_k e^{j\omega t} \Rightarrow \tilde{b} e^{j\omega t}$$

$$= b_1 e^{j\omega t} + b_2 e^{j(\omega t + \phi_2)}$$

$$= b_1 e^{j\omega t} + b_2 e^{j\phi_2} e^{j\omega t}$$

$$= e^{j\omega t} (b_1 + b_2 e^{j\phi_2}) = \hat{b} e^{j\omega t}$$

Summary

- Steady-state solution of VDE
- Concept of Impedance
- AC Circuits
- Examples

Review: Forced Solution

$$q_k^f(t) = \underline{e^{\lambda t}} \hat{b}_k \frac{e^{j(\omega - \lambda)t}}{j\omega - \lambda}$$

$$= e^{j\omega t} \left(\frac{\hat{b}_k}{j\omega - \lambda} \right) = \hat{h}_k(j\omega) e^{j\omega t}$$

↖ Constant, not function of time

$$X_k^A(t) = \sum_j \alpha_j^k q_k(t) = h_k(j\omega) e^{j\omega t}$$

$$= \left(\frac{\alpha_1^k}{j\omega - \lambda_1} + \frac{\alpha_2^k}{j\omega - \lambda_2} + \dots \right) \underbrace{e^{j\omega t}}$$

function
of time

$$\underbrace{h_k(j\omega)}$$

$$= \text{Amp \& Phase Shift} \\ H(j\omega)$$

$$H(j\omega) = \left(\frac{\alpha_1^k}{j\omega - \lambda_1} + \frac{\alpha_2^k}{j\omega - \lambda_2} + \dots \right)$$

$$= \frac{N(j\omega)}{D(j\omega)} = \frac{(j\omega - z_1) \dots (j\omega - z_n)}{(j\omega - \lambda_1) (j\omega - \lambda_2) \dots (j\omega - \lambda_m)}$$

) Transfer Function

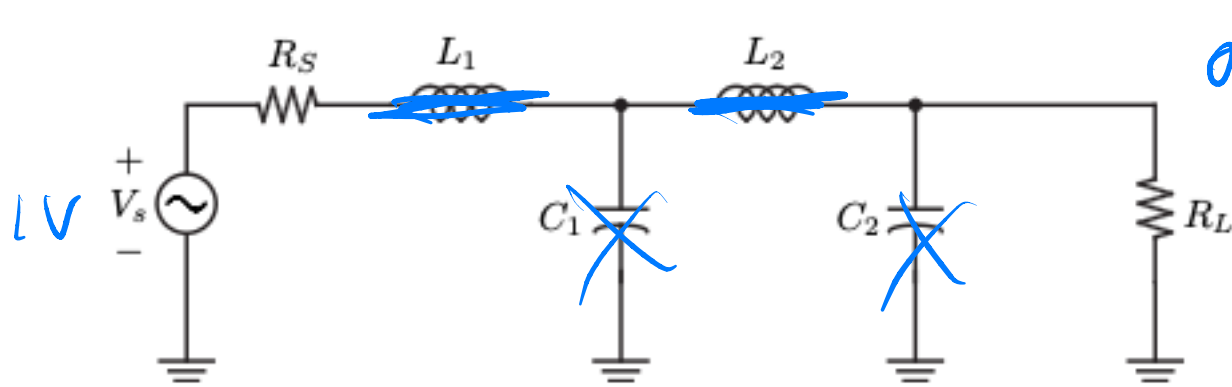
z, λ

zeros

poles

Forced Sinusoidal Solution

Example Circuit

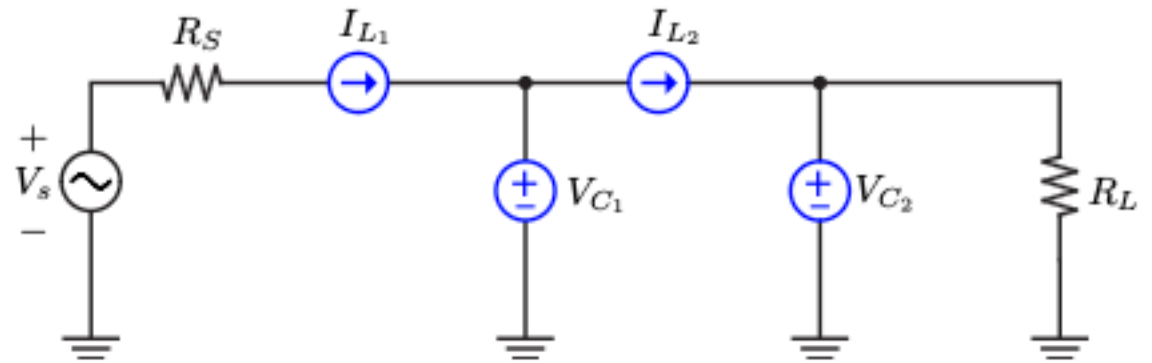


$$\vec{x} = (V_{C_1} \ V_{C_2} \ I_{L_1} \ I_{L_2})^T$$

$$\frac{d}{dt}\vec{x} = A\vec{x} + Bb_s$$

$$A = \begin{pmatrix} 0 & 0 & 1/C_1 & -1/C_1 \\ 0 & -1/R_L C_2 & 0 & 1/C_2 \\ -1/L_1 & 0 & -R_S/L_1 & 0 \\ 1/L_2 & -1/L_2 & 0 & 0 \end{pmatrix}$$

$$B = (0 \ 0 \ 1/L_1 \ 0)^T V_s.$$



Numerical Example

```
In[120]:= EV = Eigenvalues[A] ; EV // MatrixForm
```

```
Out[120]//MatrixForm=
```

$$\begin{pmatrix} -1.9034 \times 10^8 + 7.2085 \times 10^8 i \\ -1.9034 \times 10^8 - 7.2085 \times 10^8 i \\ -4.8545 \times 10^8 + 2.82637 \times 10^8 i \\ -4.8545 \times 10^8 - 2.82637 \times 10^8 i \end{pmatrix}$$

λ_1
 λ_2
 λ_3
 $\lambda_4 = \sigma_4 + j\omega_4$

Note that the eigenvalues have an imaginary part close to the cut-off frequency of the filter (100 MHz):

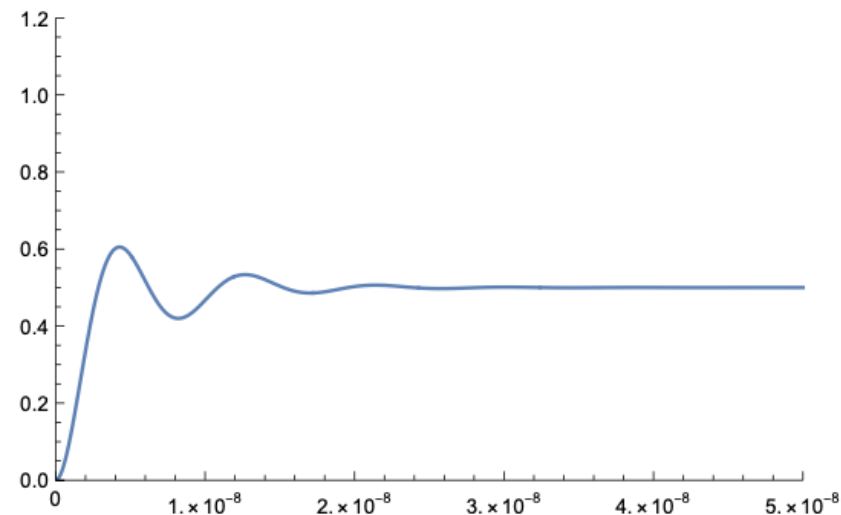
```
In[121]:= Im[EV] / (2 Pi 100 x 10 ^ 6) // N
```

```
Out[121]= {1.14727, -1.14727, 0.449831, -0.449831}
```

\nearrow 115 MHz
 \nearrow 45 MHz

```
bs = {{0}, {0}, {1/L1}, {0}};
```

```
Plot[Re[xforce[t][[2]]], {t, 0, 5010^-9}, PlotRange->{{0, 5010^-9}, {0, 1.2}}
```



Forced Response

- Since we've reduced the system to simple constant coefficient linear differential equations, we know the forced response !
- Note that the sources are also transformed into the eigenspace through the matrix Q^{-1}

```
xforce[t_] = Q.MatrixExp[Dlam t].Qi.xi + Q.MatrixExp[Dlam t].Qi .Integrate[Q.MatrixExp[-Dlam sigma].Qi .bs , {sigma, 0, t}];
```

homogeneous solution

Q

$e^{-\lambda t}$

Q^{-1}

$Q^{-1} \int Q e^{-\lambda \sigma} Q^{-1} \underline{b}_s d\sigma$

Steady-State DC Response

```
In[140]:= Limit[Q.MatrixExp[Dlam t].Qi . Integrate[Q.MatrixExp[-Dlam sigma].Qi . bs, {sigma, 0, t}], t -> Infinity]
```

```
Out[140]= {{0.5}, {0.5 + 0. i}, {0.01 + 0. i}, {0.01}}
```

The final values are also easily derived from the basic vector equation by setting all derivatives to zero:

```
In[141]:= -Inverse[A] . bs
```

```
Out[141]= {{0.5}, {0.5}, {0.01}, {0.01}}
```

$$\frac{d\vec{x}}{dt} = A\vec{x} + \vec{b}_s = \vec{0}$$

DC steady state: $\left. \begin{array}{l} L \rightarrow \text{short} \\ C \rightarrow \text{open} \end{array} \right\} \text{obvious}$

$$\vec{x}_s = -A^{-1}\vec{b}_s$$

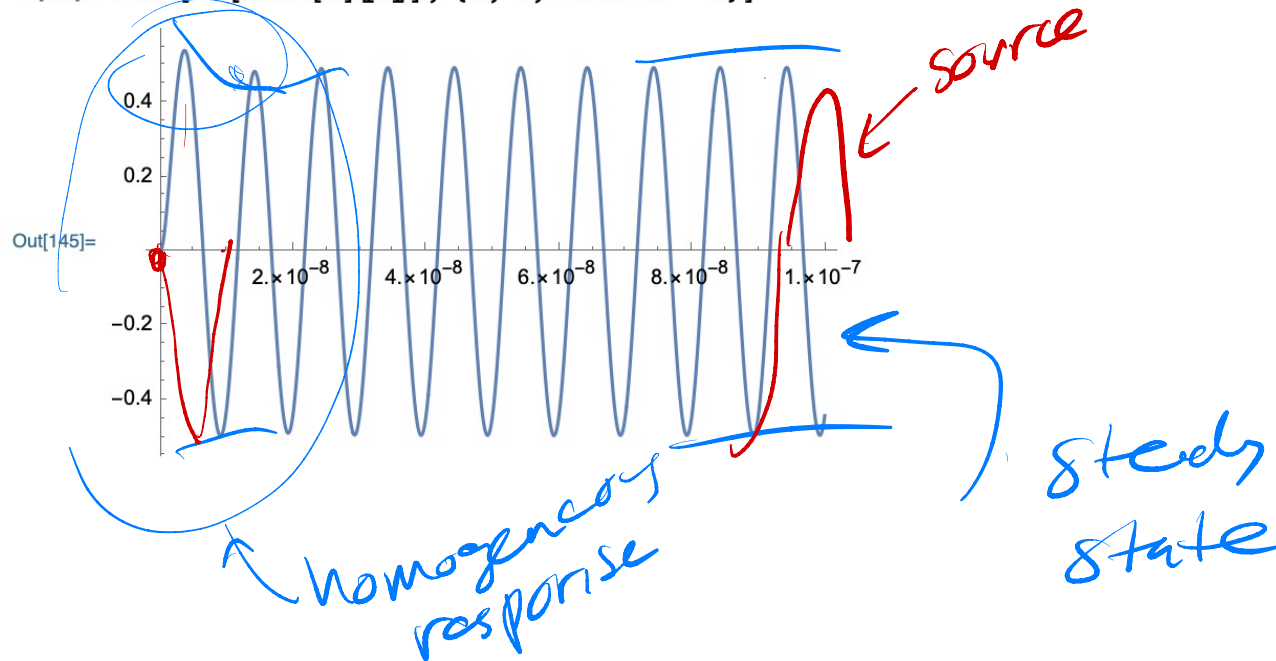
Sinusoidal Steady-State

```
w1 = 2 Pi 100 x 10 ^ 6;
```

100MHz

```
In[143]:= xac1[t_] = Q.MatrixExp[Dlam t].Qi.xi + Q.MatrixExp[Dlam t].Qi.Integrate[Q.MatrixExp[-Dlam sigma].Qi.bs Cos[w1 sigma], {sigma, 0, t}];
```

```
In[145]:= Plot[Re[xac1[t][[2]], {t, 0, 100 x 10 ^ -9}]
```

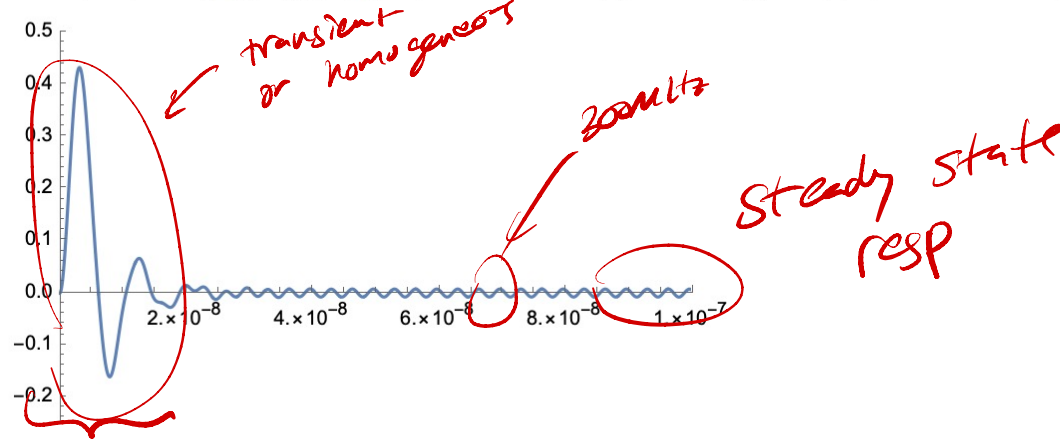


Sinusoidal source

Higher Frequency Response

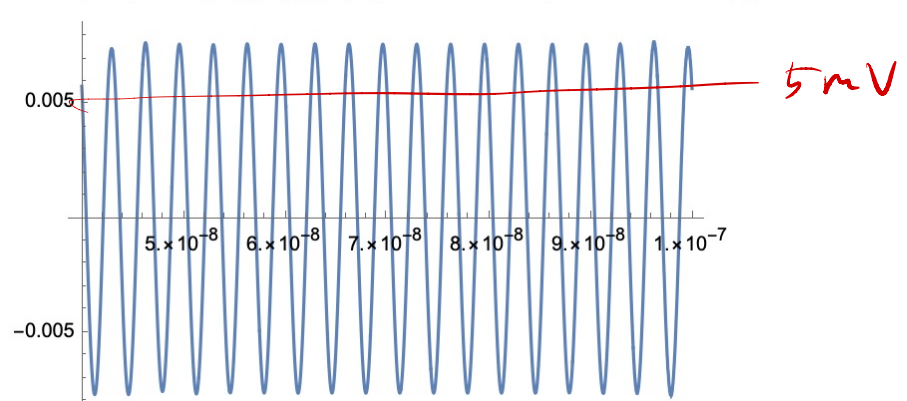
```
xac2[t_] = Q.MatrixExp[Dlam t].Qi.xi + Q.MatrixExp[Dlam t].Qi.Integrate[Q.MatrixExp[-Dlam sigma].Qi.bs Cos[3 w1 sigma], {sigma, 0, t}];
```

```
Plot[Re[xac2[t][[2]], {t, 0, 100*10^-9}, PlotRange -> {{0, 100*10^-9}, {-0.25, 0.5}}]
```



3 × 100 MHz

```
Plot[Re[xac2[t][[2]], {t, 40*10^-9, 100*10^-9}]
```



$$\frac{|H(3\omega)|}{|H(\omega)|} \approx \frac{5\text{mV}}{0.5\text{V}}$$

100 MHz

Steady-State Sinusoidal Response

- Weighted integral perspective

Filter Design

Filter Properties

Response: Lowpass
Type: Chebyshev

Topology: Series First
Order: 4

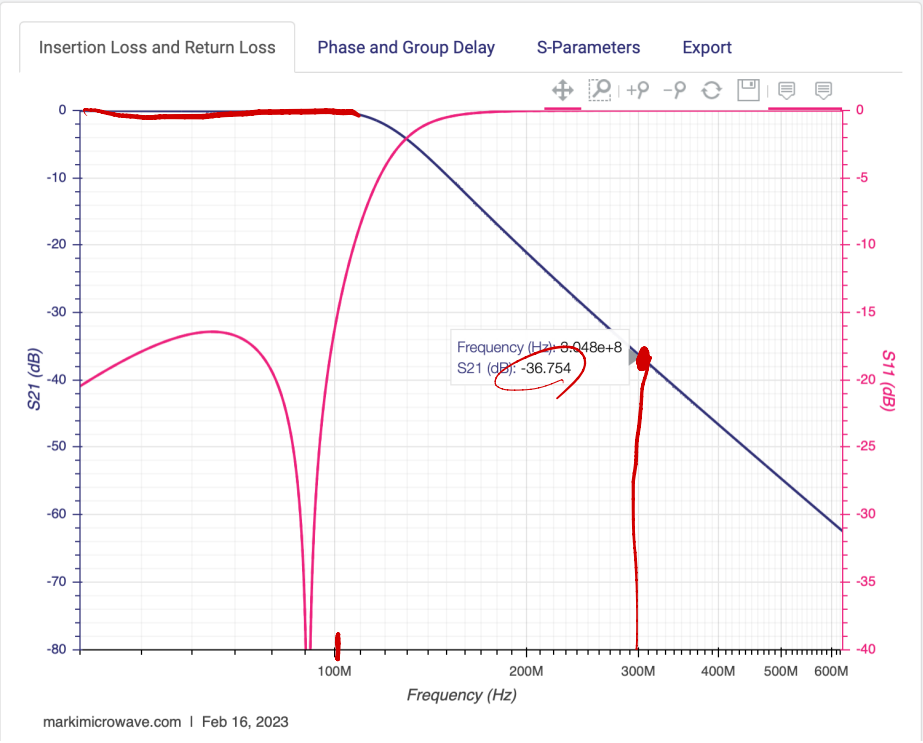
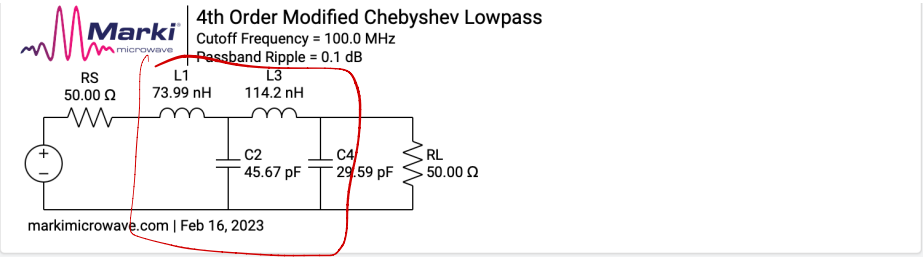
Cutoff Frequency: 100 MHz

Passband Ripple (dB): 0.10

Input Impedance (Ω): 50
Output Impedance (Ω): 50

Additional Settings
Component Values: Exact

Compute Reset



AC Analysis

- Can we get here directly without setting up all the VDE and solving it ?

Goal: Directly Calculate
Sinusoidal Steady-State Resp

“Phasors”

- A complex number by another name

\hat{I} , \hat{V} : a complex number

$$v(t) = V_0 \cos(\omega t + \phi) \Rightarrow \hat{V} \Leftarrow V_0 e^{j(\omega t + \phi)}$$
$$\Leftarrow \underbrace{V_0 e^{j\phi}}_{\hat{V}} e^{j\omega t}$$

Concept of Impedance

$$Z_C = \frac{1}{j\omega C}$$

$$Z_L = j\omega L$$

Reactance

If Z is purely imaginary

Z is reactive

$$Z = R + jX$$

resistance reactance

$$Y = G + jB = Z^{-1}$$

conductance susceptance

Ohm's Law

$$V = I \cdot R$$

$$I = G V$$



$$\hat{V} = \hat{I} Z$$

$$\hat{I} = Y \hat{V}$$

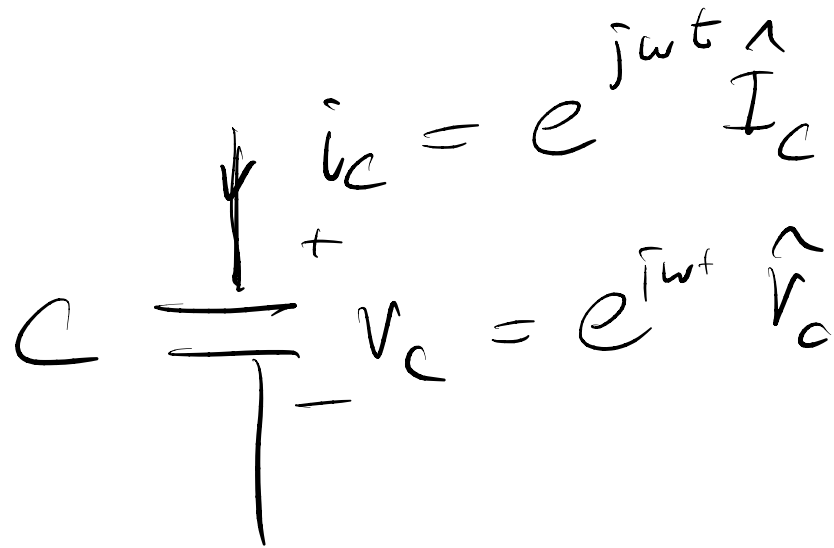
Sinusoidal
equiv.

Steady-state

Impedance

Admittance

Capacitors in AC Circuits



\hat{I}_c
 \hat{V}_c } (complex #)
constants

$$i_c = C \frac{dv_c}{dt}$$

$$e^{j\omega t} \hat{I}_c = C \frac{d}{dt} (e^{j\omega t} \hat{V}_c) = C \hat{V}_c j\omega e^{j\omega t}$$

$$\cancel{e^{j\omega t}} \hat{I}_C = C - j\omega \cdot \hat{V}_C \cancel{e^{j\omega t}}$$

A capacitor in sinusoidal steady

state looks like a "complex

resistor"

$$I = G \cdot V$$

$$\hat{I}_C = j\omega C \hat{V}_C$$

"Complex resistor" \Rightarrow impedance

Inductors in AC Circuits



S.S.S.

$$v_L = L \frac{di_L}{dt}$$

$$\hat{V}_L = j\omega L \hat{I}_L$$

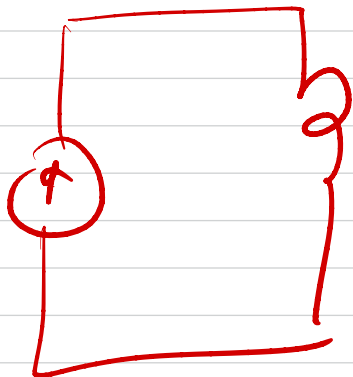
Complex

"resistor"

$$Z_L = j\omega L$$

$$Z_C = \frac{1}{j\omega C}$$

$i(t)$



Back In Time :

$$\hat{V}_c = \frac{1}{j\omega C} \hat{I}_c$$

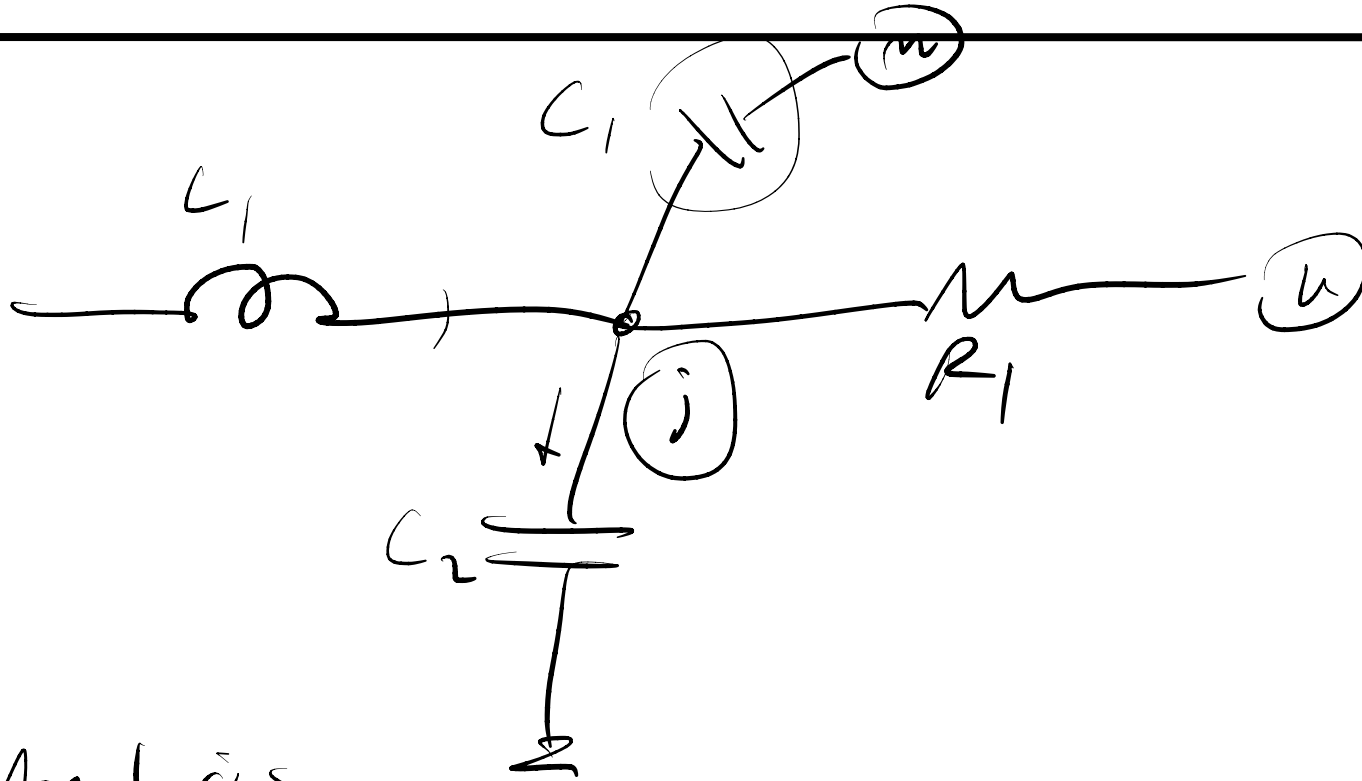
$$v(t) \stackrel{?}{\Leftrightarrow} \frac{1}{j\omega C} e^{j\phi_c} e^{j\omega t}$$

$$= -j \frac{1}{\omega C} e^{j\phi_c} e^{j\omega t}$$

$$= e^{-\frac{\pi}{2}j} \frac{1}{\omega C} e^{j\phi_c} e^{j\omega t} = e^{j\phi_c'} \frac{1}{\omega C} e^{j\omega t}$$

$$v(t) = \frac{1}{\omega C} \cos(\omega t + \phi_c')$$

Direct Solution (no VDE required)



Nodal Analysis

$$C_2 \frac{dV_j}{dt} + \frac{V_j - V_u}{R_1} + C_1 \frac{d(V_j - V_m)}{dt} + i_{L_1}$$

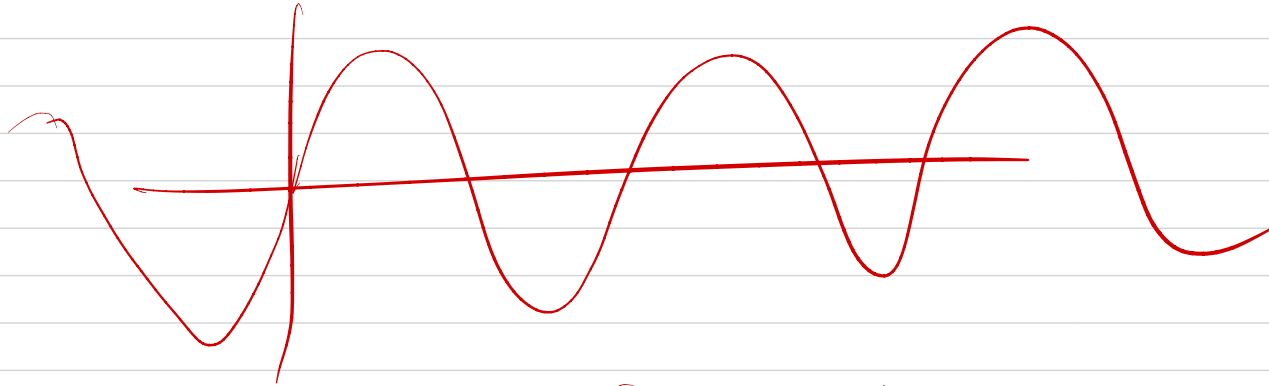
NODE (k)

$$0 = \overset{I_s}{I_s} + \sum_i C_{ij} \frac{d(V_k - V_j)}{dt} + \frac{V_k - V_j}{R_{ij}} + i_{L_j}$$

$$0 = \overset{I_s}{I_s} + \sum_i C_{ij} \frac{d(V_k - V_j)}{dt} + \frac{V_k - V_j}{R_{ij}} + \frac{1}{L_j} \int (V_k - V_j) dt$$

$$V_k = \hat{V}_k e^{j\omega t}$$

$$0 = \cancel{e^{j\omega t}} \sum_i C_{ij} (j\omega) (\hat{V}_k - \hat{V}_j) + \frac{\hat{V}_k - \hat{V}_j}{R_{ij}} + \frac{1}{L_j} \frac{1}{j\omega} (\hat{V}_k - \hat{V}_j) + \overset{I_s}{I_s} \cancel{e^{j\omega t}}$$



$$v(t) = V_0 \sin(\omega t + \phi)$$

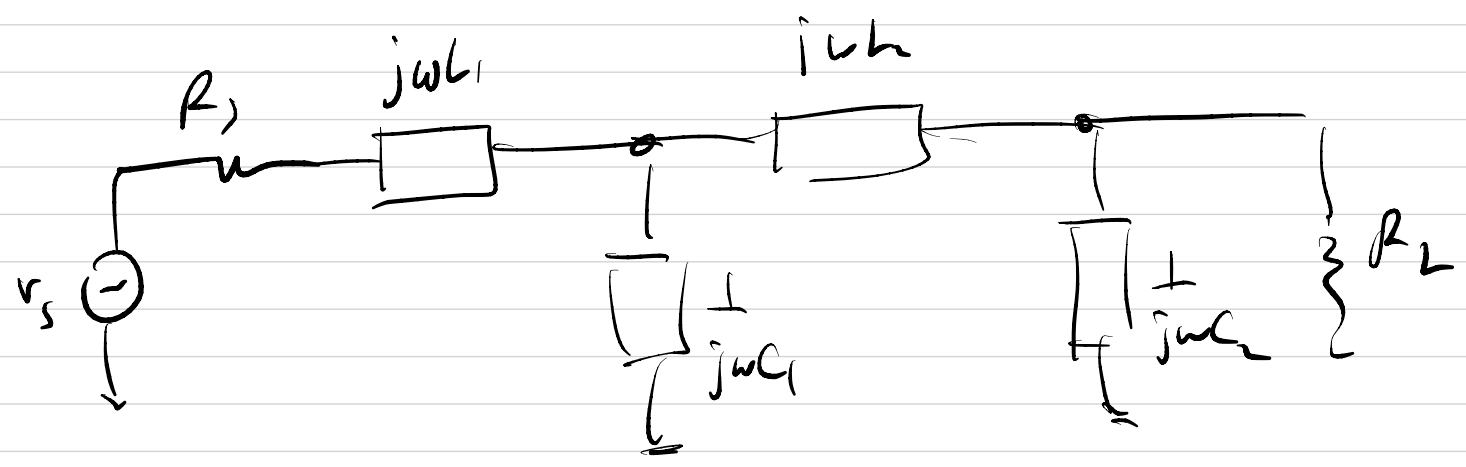
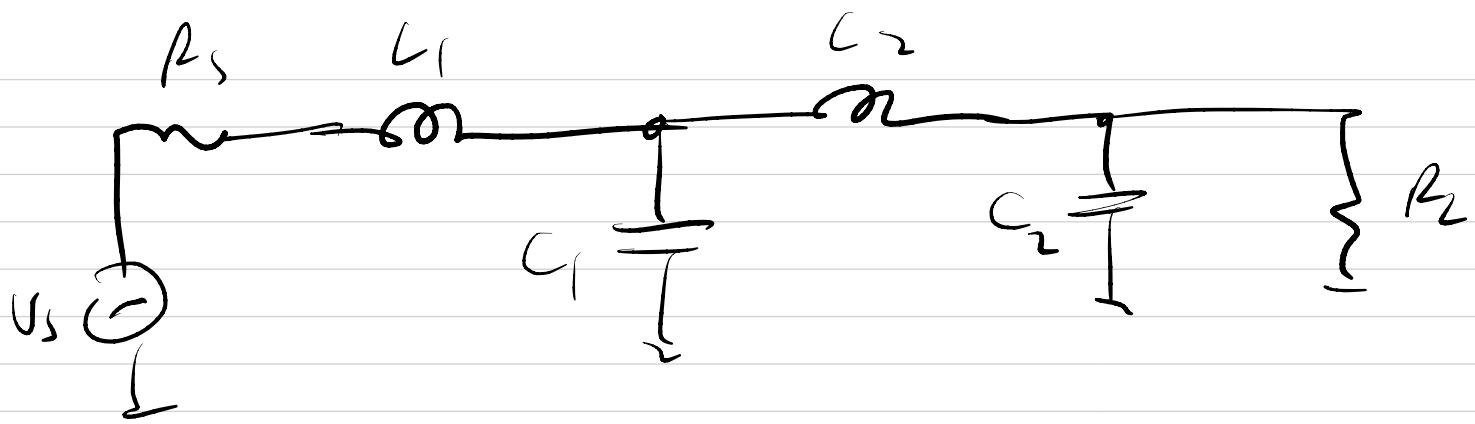
↑ ↑ ↑

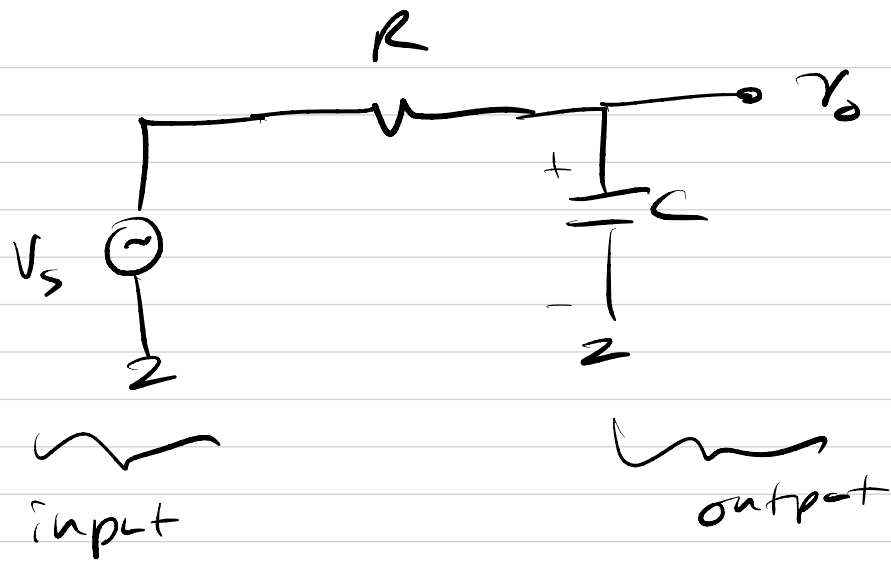
↓ frequency

$$\left(\hat{V}_0 e^{j\omega t} \right)$$

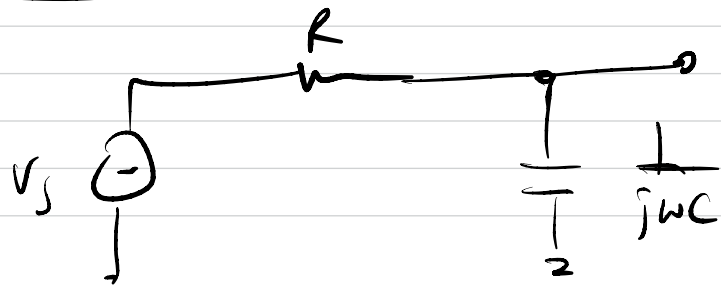
Phasor Algebra

- Same as any complex number



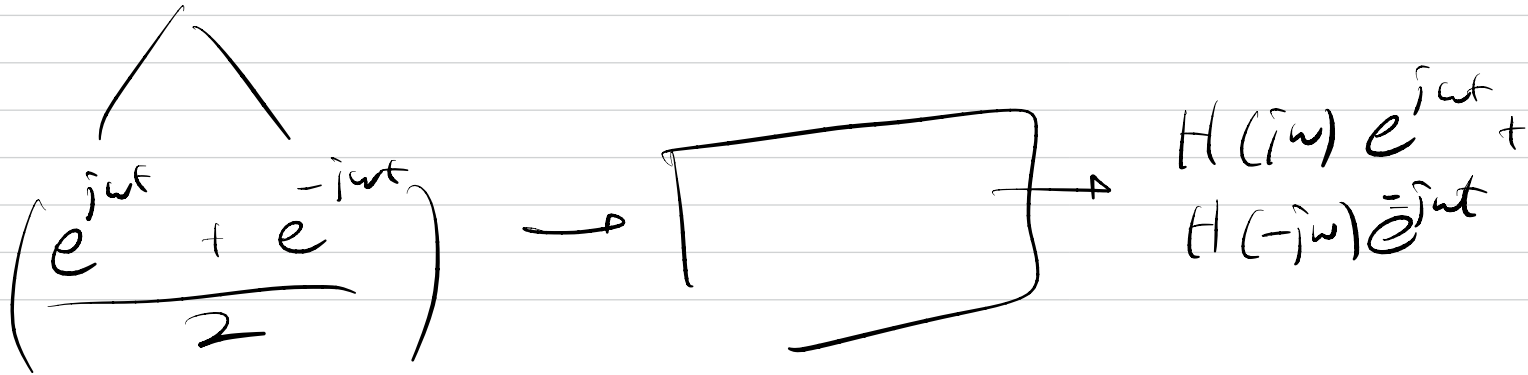
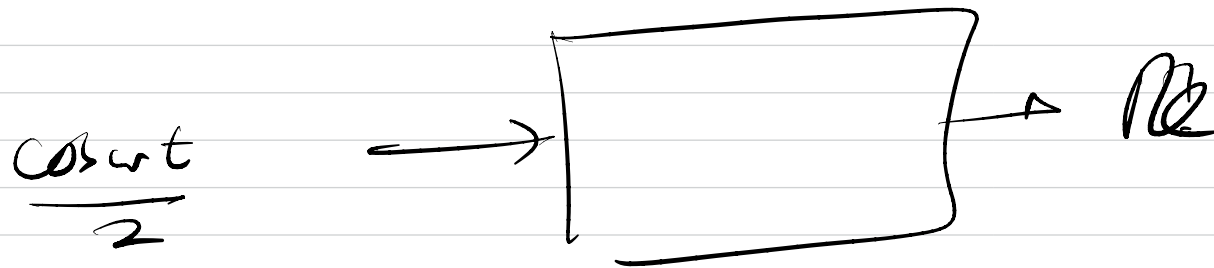
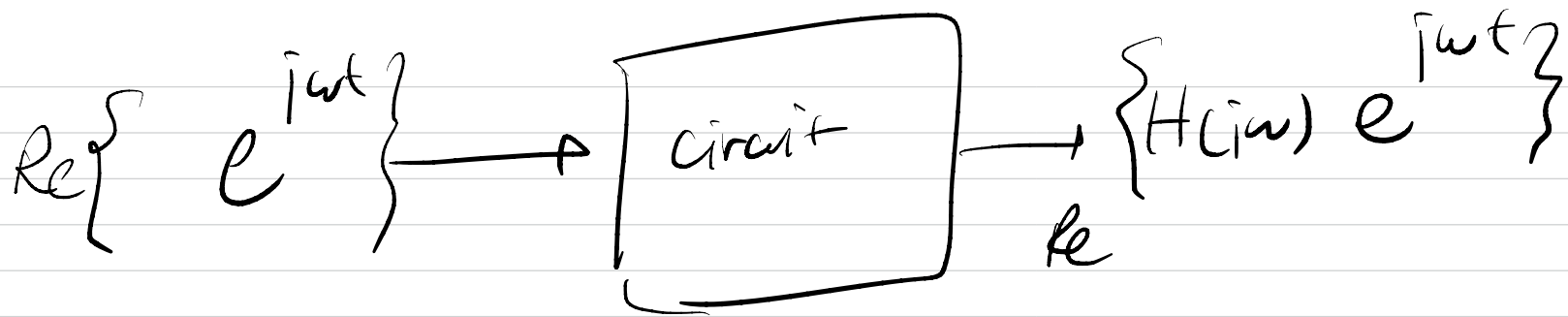


AC response



$$\frac{v_o}{v_s} = \frac{1/j\omega C}{R + 1/j\omega C}$$

$$= \frac{1}{1 + j\omega RC}$$



$$H(j\omega) = \frac{1}{1 + j\omega RC}$$

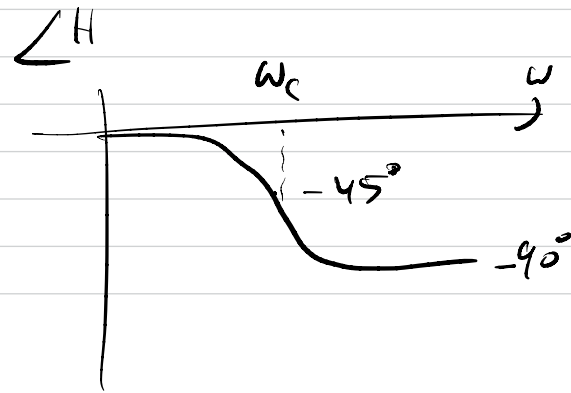
$$|H(j\omega)| = \left| \frac{1}{j\omega RC} \right| = \frac{1}{\omega RC} = \frac{1}{\omega}$$

$$|H(j\omega)| = \frac{1}{\sqrt{1 + \omega^2 (RC)^2}}$$

$$\omega_c = \frac{1}{RC}$$

Magnitude Resp

$$\angle H(j\omega) = -\tan^{-1} \left(\frac{\omega RC}{1} \right)$$



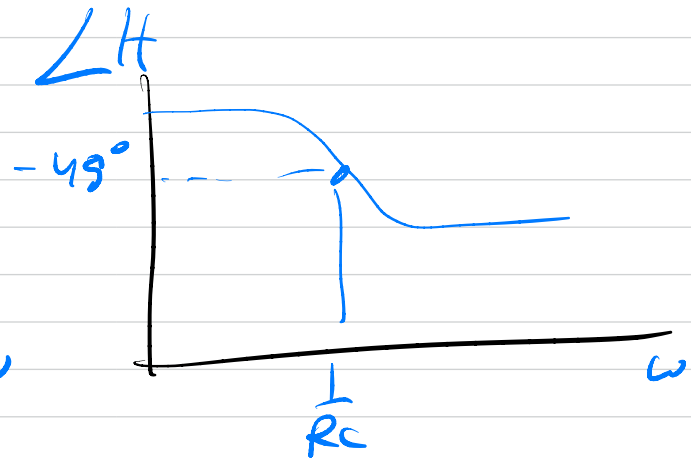
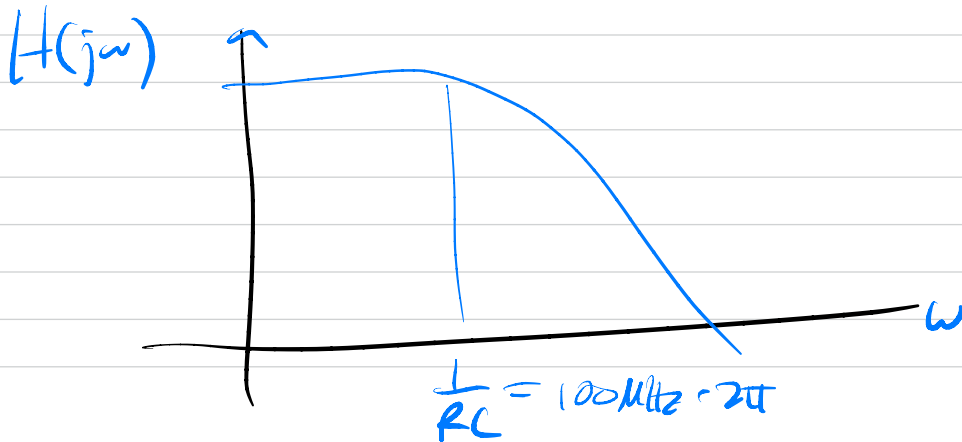
$$\vec{V}_o = H \vec{V}_s$$

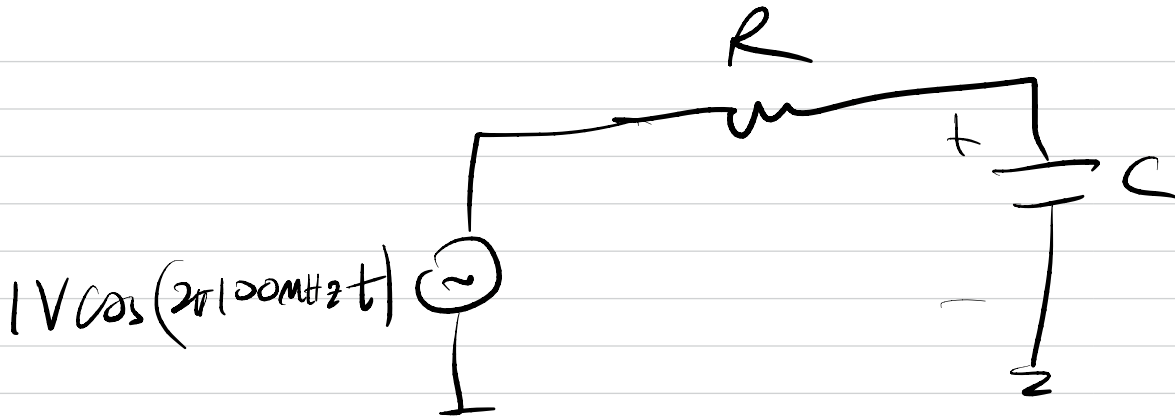
phasor phasor

$$v_s(t) = V_s \cos \omega t$$



$$v_o(t) = |H(\omega)| V_s \cos(\omega t + \angle H(\omega))$$





$$v_o(t) = 1V \cdot \left| H\left(j\frac{1}{RC}\right) \right| \cos\left(2\pi \cdot 100 \text{ MHz } t + \angle H\right)$$

$$= 1V \cdot \frac{1}{\sqrt{2}} \cos\left(2\pi \cdot 100 \text{ MHz } t - 45^\circ\right)$$

$\frac{1}{\sqrt{2}}$
 Atten.

-45°
 Phase
 Shift