

The following notes are useful for this discussion: [Note 9](#), [Note 10](#)

1. System Identification by Means of Least Squares

(a) Consider the scalar discrete-time system

$$x[i + 1] = ax[i] + bu[i] + w[i] \tag{1}$$

Where the scalar state at timestep i is $x[i]$, the input applied at timestep i is $u[i]$ and $w[i]$ represents some (small) external disturbance that also participated at timestep i (which we cannot predict or control, it's a purely random disturbance).

Assume that you have measurements for the states $x[i]$ from $i = 0$ to ℓ and also measurements for the controls $u[i]$ from $i = 0$ to $\ell - 1$. Further assume $\ell \geq 2$.

Show that we can set up a linear system as in eq. (2) to find constants a and b . How do we solve this system?

$$\underbrace{\begin{bmatrix} x[1] \\ x[2] \\ \vdots \\ x[\ell] \end{bmatrix}}_{\vec{s}} \approx \underbrace{\begin{bmatrix} x[0] & u[0] \\ x[1] & u[1] \\ \vdots & \vdots \\ x[\ell - 1] & u[\ell - 1] \end{bmatrix}}_D \underbrace{\begin{bmatrix} a \\ b \end{bmatrix}}_{\vec{p}} \tag{2}$$

Solution: Our model is of the form

$$x[i + 1] = ax[i] + bu[i] + w[i] \tag{3}$$

where $w[i]$ is our error term and we are interested in a and b . Since we cannot predict the disturbance $w[i]$ (and therefore cannot have a parameter in our solution associated with the effect of the disturbance on our system), we will solve the adjusted equation in eq. (4).

$$x[i + 1] \approx ax[i] + bu[i] \tag{4}$$

We have measurements from $i = 1$ to $i = m$, and so our least squares formulation is:

$$\underbrace{\begin{bmatrix} x[1] \\ x[2] \\ \vdots \\ x[\ell] \end{bmatrix}}_{\vec{s}} \approx \underbrace{\begin{bmatrix} x[0] & u[0] \\ x[1] & u[1] \\ \vdots & \vdots \\ x[\ell - 1] & u[\ell - 1] \end{bmatrix}}_D \underbrace{\begin{bmatrix} a \\ b \end{bmatrix}}_{\vec{p}} \tag{5}$$

D is not necessarily a square matrix (it is tall), so we cannot invert it and solve for \vec{p} . Hence, we use least squares like previously mentioned. Thus, our best approximation for \vec{p} is

$$\hat{\vec{p}} = (D^T D)^{-1} D^T \vec{s} \tag{6}$$

Since we are using least squares, we can also group our estimation error (remember, $\hat{\vec{p}} \neq \vec{p}$ necessarily) into $w[i]$.

- (b) What if there were now two distinct scalar inputs to a scalar system

$$x[i + 1] = ax[i] + b_1u_1[i] + b_2u_2[i] + w[i] \quad (7)$$

and that we have measurements as before, but now also for both of the control inputs.

Set up a least-squares problem that you can solve to get an estimate of the unknown system parameters a, b_1, b_2 .

Solution: Our new model is of the form

$$x[i + 1] = ax[i] + b_1u_1[i] + b_2u_2[i] + w[i] \quad (8)$$

where $w[i]$ is our error term and we are interested in a, b_1, b_2 . As we did before, we will modify the system and drop the disturbance term, converting the equality to an approximation.

$$x[i + 1] \approx ax[i] + b_1u_1[i] + b_2u_2[i] \quad (9)$$

As before, we have $[1, m]$ measurements, and so our least squares formulation is:

$$\underbrace{\begin{bmatrix} x[1] \\ x[2] \\ \vdots \\ x[\ell] \end{bmatrix}}_{\vec{s}} \approx \underbrace{\begin{bmatrix} x[0] & u_1[0] & u_2[0] \\ x[1] & u_1[1] & u_2[1] \\ \vdots & \vdots & \vdots \\ x[\ell-1] & u_1[\ell-1] & u_2[\ell-1] \end{bmatrix}}_D \underbrace{\begin{bmatrix} a \\ b_1 \\ b_2 \end{bmatrix}}_{\vec{p}} \quad (10)$$

- (c) **What could go wrong in the previous case? For what kind of inputs would make least-squares fail to give you the parameters you want?**

Solution: We can take a look at the least squares formula, and think about what the possible failure points are.

$$\hat{\vec{p}} = (D^T D)^{-1} D^T \vec{s}. \quad (11)$$

In this equation, the likely point of failure is the inversion of $D^T D$; the other operations (matrix-matrix multiplications, matrix-vector multiplications) do not have the same issue.

$D^T D$ might not be invertible when D has columns that are not linearly independent. For example, it could be because the inputs \vec{u}_1 and \vec{u}_2 are too similar, as if $\vec{u}_1 = \alpha \vec{u}_2$. We need these two inputs to be different and sufficiently varied so that least-squares does not fail.

- (d) Now consider the two dimensional state case with a single input.

$$\vec{x}[i + 1] = \begin{bmatrix} x_1[i + 1] \\ x_2[i + 1] \end{bmatrix} = \begin{bmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{bmatrix} \vec{x}[i] + \begin{bmatrix} b_1 \\ b_2 \end{bmatrix} u[i] + \vec{w}[i] \quad (12)$$

How can we treat this like two parallel problems to set this up using least-squares to get estimates for the unknown parameters $a_{11}, a_{12}, a_{21}, a_{22}, b_1, b_2$? Write the least squares solution in terms of your known matrices and vectors (including based on the labels you gave to various matrices/vectors in previous parts). *Hint: What work/computation can we reuse across the two problems?*

Solution: We can rewrite eq. (12) as

$$\begin{bmatrix} x_1[i+1] \\ x_2[i+1] \end{bmatrix} = \begin{bmatrix} a_{11}x_1[i] + a_{12}x_2[i] + b_1u[i] \\ a_{21}x_1[i] + a_{22}x_2[i] + b_2u[i] \end{bmatrix} \quad (13)$$

We can set up a problem to solve for a_{11}, a_{12}, b_1 (call this subsystem 1) and another problem to solve for a_{21}, a_{22}, b_2 (call this subsystem 2). We can rewrite the first row of eq. (13) as

$$x_1[i+1] = \begin{bmatrix} x_1[i] & x_2[i] & u[i] \end{bmatrix} \begin{bmatrix} a_{11} \\ a_{12} \\ b_1 \end{bmatrix} \quad (14)$$

and likewise for the second row

$$x_2[i+1] = \begin{bmatrix} x_1[i] & x_2[i] & u[i] \end{bmatrix} \begin{bmatrix} a_{21} \\ a_{22} \\ b_2 \end{bmatrix} \quad (15)$$

To find the unknowns in subsystem 1, we can set up the following least squares problem:

$$\underbrace{\begin{bmatrix} x_1[1] \\ x_1[2] \\ \vdots \\ x_1[\ell] \end{bmatrix}}_{\tilde{s}_1} \approx \underbrace{\begin{bmatrix} x_1[0] & x_2[0] & u[0] \\ x_1[1] & x_2[1] & u[1] \\ \vdots & \vdots & \vdots \\ x_1[\ell-1] & x_2[\ell-1] & u[\ell-1] \end{bmatrix}}_{D_1} \underbrace{\begin{bmatrix} a_{11} \\ a_{12} \\ b_1 \end{bmatrix}}_{\tilde{p}_1} \quad (16)$$

Now, to find the unknowns in subsystem 2, we can set up the following least squares problem:

$$\underbrace{\begin{bmatrix} x_2[1] \\ x_2[2] \\ \vdots \\ x_2[\ell] \end{bmatrix}}_{\tilde{s}_2} \approx \underbrace{\begin{bmatrix} x_1[0] & x_2[0] & u[0] \\ x_1[1] & x_2[1] & u[1] \\ \vdots & \vdots & \vdots \\ x_1[\ell-1] & x_2[\ell-1] & u[\ell-1] \end{bmatrix}}_{D_2} \underbrace{\begin{bmatrix} a_{21} \\ a_{22} \\ b_2 \end{bmatrix}}_{\tilde{p}_2} \quad (17)$$

Notice that $D_1 = D_2$. Hence, we can write $D = D_1 = D_2$, and we only need to compute $(D^\top D)^{-1} D^\top$ once. Hence, the solution for the i th subsystem (for $i \in \{1, 2\}$) is

$$\hat{\tilde{p}}_i = (D^\top D)^{-1} D^\top \tilde{s}_i \quad (18)$$

Furthermore, we can horizontally stack the two separate problems for each subsystem as follows:

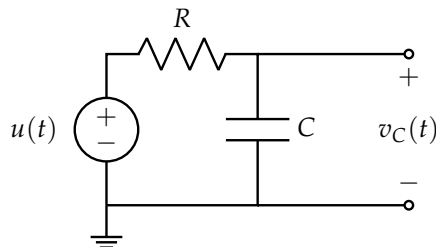
$$\underbrace{\begin{bmatrix} x_1[1] & x_2[1] \\ x_1[2] & x_2[2] \\ \vdots & \vdots \\ x_1[\ell] & x_2[\ell] \end{bmatrix}}_S \approx \underbrace{\begin{bmatrix} x_1[0] & x_2[0] & u[0] \\ x_1[1] & x_2[1] & u[1] \\ \vdots & \vdots & \vdots \\ x_1[\ell-1] & x_2[\ell-1] & u[\ell-1] \end{bmatrix}}_D \underbrace{\begin{bmatrix} a_{11} & a_{21} \\ a_{12} & a_{22} \\ b_1 & b_2 \end{bmatrix}}_P \quad (19)$$

Finally, solving this as a single least squares problem gives us

$$\hat{P} = (D^\top D)^{-1} D^\top S \quad (20)$$

2. Stability Examples and Counterexamples

- (a) Consider the circuit below with $R = 1 \Omega$, $C = 0.5 \text{ F}$, and $u(t)$ is some function bounded between $-K$ and K for some constant $K \in \mathbb{R}$ (for example $K \cos(t)$). Furthermore assume that $v_C(0) = 0 \text{ V}$ (that the capacitor is initially discharged).



This circuit can be modeled by the differential equation

$$\frac{dv_C(t)}{dt} = -2v_C(t) + 2u(t) \quad (21)$$

Show that the differential equation is always stable (that is, as long as the input $u(t)$ is bounded, $v_C(t)$ also stays bounded). Consider what this means in the physical circuit. *HINT: You may want to use the triangle inequality, i.e. $|a + b| \leq |a| + |b|$, and the triangle inequality for integrals, i.e. $\left| \int_a^b f(x) dx \right| \leq \int_a^b |f(x)| dx$. When we use $|\cdot|$ notation here, we will take this to mean the magnitude, rather than the absolute value (since we can be dealing with complex numbers).*

Solution: We can apply the integral solution for a nonhomogeneous differential equation to demonstrate boundedness of the solution. The general solution to $\frac{dx(t)}{dt} = \lambda x(t) + bu(t)$ is $x(t) = x_0 e^{\lambda t} + \int_0^t e^{\lambda(t-\theta)} bu(\theta) d\theta$. Here, we can say that:

$$v_C(t) = v_C(0)e^{-2t} + \int_0^t e^{-2(t-\theta)} 2u(\theta) d\theta \quad (22)$$

$$= v_C(0)e^{-2t} + 2 \int_0^t e^{-2(t-\theta)} u(\theta) d\theta \quad (23)$$

We wish to show $|v_C(t)| \leq M$ for all $t \geq 0$, where $M \in \mathbb{R}$ is some constant (this is another way to say that something is “bounded”). We can take the absolute value around eq. (23) as follows:

$$|v_C(t)| = \left| v_C(0)e^{-2t} + 2 \int_0^t e^{-2(t-\theta)} u(\theta) d\theta \right| \quad (24)$$

$$\leq |v_C(0)e^{-2t}| + \left| 2 \int_0^t e^{-2(t-\theta)} u(\theta) d\theta \right| \quad (25)$$

$$\leq |v_C(0)e^{-2t}| + 2 \int_0^t |e^{-2(t-\theta)} u(\theta)| d\theta \quad (26)$$

$$= |v_C(0)|e^{-2t} + 2 \int_0^t e^{-2(t-\theta)} |u(\theta)| d\theta \quad (27)$$

where we use the traditional triangle inequality to obtain eq. (25) and the integral triangle inequality to obtain eq. (26). We know $v_C(0) = 0$, so the first term is 0. Even if it is nonzero, we may assume that it is some finite constant. Furthermore, $0 \leq e^{-2t} \leq 1$ for $t \geq 0$ (it is a decaying exponential). Hence, the $|v_C(0)|e^{-2t}$ term is bounded. Next, we are allowed to assume that

$|u(t)| \leq K$ from the statement of the problem. This will let us obtain

$$|v_C(t)| \leq 2 \int_0^t e^{-2(t-\theta)} \underbrace{|u(\theta)|}_{\leq K} d\theta \quad (28)$$

$$\leq 2K \int_0^t e^{-2(t-\theta)} d\theta \quad (29)$$

$$= K(1 - e^{-2t}) \quad (30)$$

Because $e^{-2t} \geq 0$, $1 - e^{-2t} \leq 1$. Hence, $|v_C(t)| \leq K$ so $v_C(t)$ is bounded.

- (b) **(PRACTICE)** Now, suppose that in the circuit of part 2.a we replaced the resistor with an inductor as in fig. 1.

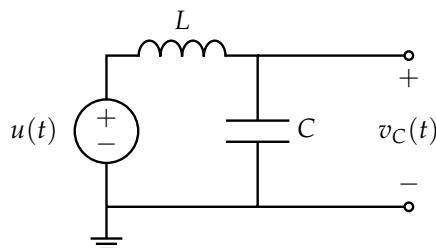


Figure 1: The original circuit with an inductor in place of the resistor.

Let $L = 1$ mH. Repeat part 2.a for the new circuit (with an inductor). Consider the following process to arrive at the result:

- i. Derive the system of differential equations using KCL, KVL, and NVA. Show that the system is $\frac{d}{dt} \begin{bmatrix} v_C(t) \\ i_L(t) \end{bmatrix} = \begin{bmatrix} 0 & \frac{1}{C} \\ -\frac{1}{L} & 0 \end{bmatrix} \begin{bmatrix} v_C(t) \\ i_L(t) \end{bmatrix} + \begin{bmatrix} 0 \\ \frac{1}{L} \end{bmatrix} u(t)$ with the initial condition being $\begin{bmatrix} v_C(0) \\ i_L(0) \end{bmatrix} = \vec{0}$.
- ii. Solve the matrix differential equation, using diagonalization if needed. Show that the diagonalized system has a solution

$$\vec{y}(t) = \begin{bmatrix} \frac{1}{2LC} e^{j\frac{1}{\sqrt{LC}}t} \int_0^t e^{-j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta \\ \frac{1}{2LC} e^{-j\frac{1}{\sqrt{LC}}t} \int_0^t e^{j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta \end{bmatrix} \quad (31)$$

where $\vec{y}(t) = V^{-1} \begin{bmatrix} v_C(t) \\ i_L(t) \end{bmatrix}$ for change of basis matrix V . You may use the fact that the

eigenvalue, eigenvector pairs of $\begin{bmatrix} 0 & \frac{1}{C} \\ -\frac{1}{L} & 0 \end{bmatrix}$ are $\left(j\frac{1}{\sqrt{LC}}, \begin{bmatrix} -j\sqrt{\frac{L}{C}} \\ 1 \end{bmatrix} \right)$ and $\left(-j\frac{1}{\sqrt{LC}}, \begin{bmatrix} j\sqrt{\frac{L}{C}} \\ 1 \end{bmatrix} \right)$.

- iii. Apply a similar process from part 2.a to show that, if we have a bounded input $u(t)$, then the system can grow unboundedly. When showing that a system is unstable, it suffices to choose a bounded $u(t)$ that makes the system unbounded. We can choose $u(t) = 2 \cos\left(\frac{1}{\sqrt{LC}}t\right) = e^{j\frac{1}{\sqrt{LC}}t} + e^{-j\frac{1}{\sqrt{LC}}t}$ ¹. HINT: You may use the fact that $i_L(t) = y_1(t) + y_2(t)$.

¹The natural frequency of this system is $\omega_n = \frac{1}{\sqrt{LC}}$. If we excite this system at a period equal to the natural frequency, we can make it grow unboundedly. This is similar to pushing a swing at the same rate it swings, which makes it swing farther.

Hint: You might find it useful to revisit the process of generating the state-space equations for $v_C(t)$ and $i_L(t)$ as done in Note 4 for the LC Tank. The difference is that here, we have an input voltage.

Solution: 2.(b)i:

First, we begin forming the vector state-space equation, which involves relating $v_C(t)$ and $i_L(t)$ to their derivatives and the input voltage.

$$C \frac{dv_C(t)}{dt} = i_C(t) = i_L(t) \quad (32)$$

$$\implies \frac{dv_C(t)}{dt} = \frac{1}{C} i_L(t) \quad (33)$$

$$L \frac{di_L(t)}{dt} = v_L(t) = u(t) - v_C(t) \quad (34)$$

$$\implies \frac{di_L(t)}{dt} = \frac{1}{L} v_L(t) = -\frac{1}{L} v_C(t) + \frac{1}{L} u(t) \quad (35)$$

Combining this info, we find:

$$\frac{d}{dt} \begin{bmatrix} v_C(t) \\ i_L(t) \end{bmatrix} = \underbrace{\begin{bmatrix} 0 & \frac{1}{C} \\ -\frac{1}{L} & 0 \end{bmatrix}}_A \underbrace{\begin{bmatrix} v_C(t) \\ i_L(t) \end{bmatrix}}_{\vec{x}(t)} + \underbrace{\begin{bmatrix} 0 \\ \frac{1}{L} \end{bmatrix}}_{\vec{b}} u(t) \quad (36)$$

2.(b)ii:

This is not a diagonal system, so we have to diagonalize it first. We start by solving for the eigenvalues and eigenvectors of A :

$$\lambda_1 = j \frac{1}{\sqrt{LC}} \quad \vec{v}_1 = \begin{bmatrix} -j\sqrt{\frac{L}{C}} \\ 1 \end{bmatrix} \quad (37)$$

$$\lambda_2 = -j \frac{1}{\sqrt{LC}} \quad \vec{v}_2 = \begin{bmatrix} j\sqrt{\frac{L}{C}} \\ 1 \end{bmatrix} \quad (38)$$

Note that these eigenvalues are purely imaginary. This will be helpful later. Our change of basis matrix is $V = \begin{bmatrix} -j\sqrt{\frac{L}{C}} & j\sqrt{\frac{L}{C}} \\ 1 & 1 \end{bmatrix}$, so we can define our change of basis as $\vec{y}(t) = V^{-1}\vec{x}(t)$. Note that the new diagonal system will be

$$\frac{d}{dt} \vec{y}(t) = \begin{bmatrix} j\frac{1}{\sqrt{LC}} & 0 \\ 0 & -j\frac{1}{\sqrt{LC}} \end{bmatrix} \vec{y}(t) + V^{-1} \vec{b} u(t) \quad (39)$$

$$= \begin{bmatrix} j\frac{1}{\sqrt{LC}} & 0 \\ 0 & -j\frac{1}{\sqrt{LC}} \end{bmatrix} \vec{y}(t) + \left(\begin{bmatrix} -j\sqrt{\frac{L}{C}} & j\sqrt{\frac{L}{C}} \\ 1 & 1 \end{bmatrix}^{-1} \begin{bmatrix} 0 \\ \frac{1}{L} \end{bmatrix} \right) u(t) \quad (40)$$

$$= \begin{bmatrix} j\frac{1}{\sqrt{LC}} & 0 \\ 0 & -j\frac{1}{\sqrt{LC}} \end{bmatrix} \vec{y}(t) + \begin{bmatrix} \frac{1}{2LC} \\ \frac{1}{2LC} \end{bmatrix} u(t) \quad (41)$$

so our system of equations is

$$\frac{d}{dt} y_1(t) = j \frac{1}{\sqrt{LC}} y_1(t) + \frac{1}{2LC} u(t) \quad (42)$$

$$\frac{d}{dt}y_2(t) = -j\frac{1}{\sqrt{LC}}y_2(t) + \frac{1}{2LC}u(t) \quad (43)$$

$$(44)$$

Recall that $\vec{x}(0) = \vec{0}$, so $\vec{y}(t) = \vec{0}$ (where $\vec{0}$ is a vector of all zeros). Solving this differential equation now, we get

$$y_1(t) = \underbrace{y_1(0)}_0 e^{j\frac{1}{\sqrt{LC}}t} + \int_0^t e^{j\frac{1}{\sqrt{LC}}(t-\theta)} \left(\frac{1}{2LC}u(\theta) \right) d\theta \quad (45)$$

$$y_2(t) = \underbrace{y_2(0)}_0 e^{-j\frac{1}{\sqrt{LC}}t} + \int_0^t e^{-j\frac{1}{\sqrt{LC}}(t-\theta)} \left(\frac{1}{2LC}u(\theta) \right) d\theta \quad (46)$$

Simplifying and stacking the solutions in vector form,

$$\begin{bmatrix} v_C(t) \\ i_L(t) \end{bmatrix} = \vec{x}(t) = V \begin{bmatrix} \frac{1}{2LC}e^{j\frac{1}{\sqrt{LC}}t} \int_0^t e^{-j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta \\ \frac{1}{2LC}e^{-j\frac{1}{\sqrt{LC}}t} \int_0^t e^{j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta \end{bmatrix} \quad (47)$$

2.(b)iii:

We wish to show $\vec{x}(t)$ is unbounded, given some bounded input $u(t)$. When showing a vector is bounded, we can show that all of its individual, scalar entries are bounded. Alternatively, when showing a vector is unbounded, it is enough to show that one of its entries will be unbounded. Note that $i_L(t) = y_1(t) + y_2(t)$ (which we see by computing $\vec{x}(t) = V\vec{y}(t)$). We can show that this quantity is unbounded. Recall that

$$y_1(t) = \frac{e^{j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t e^{-j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta \quad (48)$$

$$y_2(t) = \frac{e^{-j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t e^{j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta \quad (49)$$

$$\implies i_L(t) = \frac{e^{j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t e^{-j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta + \frac{e^{-j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t e^{j\frac{1}{\sqrt{LC}}\theta} u(\theta) d\theta \quad (50)$$

Now, we have to make some choice of a bounded input $u(t)$ so the entire term is unbounded as $t \rightarrow \infty$. We can choose $u(t) = e^{-j\frac{1}{\sqrt{LC}}t} + e^{j\frac{1}{\sqrt{LC}}t} = 2 \cos\left(\frac{1}{\sqrt{LC}}t\right)$ which is a bounded sinusoidal function. We can first compute $i_L(t)$ with this input:

$$i_L(t) = \frac{e^{j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t e^{-j\frac{1}{\sqrt{LC}}\theta} \left(e^{-j\frac{1}{\sqrt{LC}}\theta} + e^{j\frac{1}{\sqrt{LC}}\theta} \right) d\theta + \frac{e^{-j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t e^{j\frac{1}{\sqrt{LC}}\theta} \left(e^{-j\frac{1}{\sqrt{LC}}\theta} + e^{j\frac{1}{\sqrt{LC}}\theta} \right) d\theta \quad (51)$$

$$= \frac{e^{j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t 1 + e^{-j\frac{2}{\sqrt{LC}}\theta} d\theta + \frac{e^{-j\frac{1}{\sqrt{LC}}t}}{2LC} \int_0^t 1 + e^{j\frac{2}{\sqrt{LC}}\theta} d\theta \quad (52)$$

$$= \frac{e^{j\frac{1}{\sqrt{LC}}t}}{2LC} \left(t + \frac{1 - e^{-j\frac{2}{\sqrt{LC}}t}}{j\frac{2}{\sqrt{LC}}} \right) + \frac{e^{-j\frac{1}{\sqrt{LC}}t}}{2LC} \left(t + \frac{e^{j\frac{2}{\sqrt{LC}}t} - 1}{j\frac{2}{\sqrt{LC}}} \right) \quad (53)$$

$$= \frac{t}{LC} \left(\frac{e^{j\frac{1}{\sqrt{LC}}t} + e^{-j\frac{1}{\sqrt{LC}}t}}{2} \right) + \frac{1}{\sqrt{LC}} \left(\frac{e^{j\frac{1}{\sqrt{LC}}t} - e^{-j\frac{1}{\sqrt{LC}}t}}{2j} \right) \quad (54)$$

$$= \frac{t}{LC} \cos\left(\frac{t}{\sqrt{LC}}\right) + \frac{1}{\sqrt{LC}} \sin\left(\frac{t}{\sqrt{LC}}\right) \quad (55)$$

Notice that the cos and sin terms are bounded, but the cos term is multiplied by a t , so as $t \rightarrow \infty$, $i_L(t) \rightarrow \infty$. Hence, the system is unstable. Generally, we say a system with eigenvalues having negative real part implies stability. Here, the real part of the eigenvalues is 0, so the system is unstable.

- (c) Thus far, we have dealt with continuous systems so it also makes sense to consider discrete systems. Consider the discrete system

$$x[i+1] = 2x[i] + u[i] \quad (56)$$

with $x[0] = 0$.

Is the system stable or unstable? If unstable, find a bounded input sequence $u[i]$ that causes the system to “blow up”.

Solution: Notice that, if we had the system

$$x[i+1] = 2x[i] \quad (57)$$

then we can write $x[i+1] = 2^i x[1]$. So, if we can somehow make $x[1]$ nonzero using a bounded input (e.g. equal to 1, for simplicity), then as $i \rightarrow \infty$, $x[i+1] \rightarrow \infty$. We know that $x[0] = 0$, and that $x[1] = 2x[0] + u[0] = u[0]$. Hence, we can set $u[0] = 1$ and then $x[1] = 1$. We have achieved what we wanted, i.e. to make $x[1]$ a nonzero value using the bounded input $u[0] = 1$. Now, for the other timesteps $i > 0$, we can set $u[i] = 0$ since that would leave us with the system in eq. (57). Written explicitly, our bounded input is

$$u[i] = \begin{cases} 1 & i = 0 \\ 0 & i > 0 \end{cases} \quad (58)$$

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