

## Homework 1

1. (a) Find the quality (Q) factor of the cavity.

$$Q = \omega\tau = \frac{2\pi c}{\lambda_0}\tau$$

$$\tau = \frac{n}{\alpha c}$$

$$\alpha = \alpha_m + \alpha_i = \frac{1}{2L} \ln \frac{1}{R_1 R_2} + \alpha_i$$

Using  $\alpha_i = 10 \text{ cm}^{-1}$ ,  $R_1 = 1$ ,  $R_2 = 0.99$ , and  $L = 1\mu\text{m}$ , we find

$$\alpha_m = 50.25 \text{ cm}^{-1},$$

$$\alpha = 60.25 \text{ cm}^{-1}.$$

Plugging in, and using  $\lambda_0 = 1\mu\text{m}$  and  $n = 3.5$ , we find

$$Q = \frac{2\pi c}{\lambda_0} \frac{n}{\alpha} = 3650.$$

- (b) Find the threshold gain and quantum efficiency of the laser.

$$g_{th} = \frac{\alpha}{\Gamma} = 60.25 \text{ cm}^{-1}$$

using  $\alpha$  from before and  $\Gamma = 1$ .

$$\eta = \frac{\alpha_m}{\alpha_i + \alpha_m} = 0.834$$

using  $\alpha_m$  and  $\alpha_i$  from before.

2. (a) What is the total Q of the cavity?

Start with the definition of  $Q$ :

$$Q = \omega \frac{E}{P_{tot}} = \omega \frac{E}{P_r + P_m + P_s},$$

where  $P_{tot}$  is total power lost, and  $P_r$ ,  $P_m$ , and  $P_s$  are power lost in radiation, metal, and semiconductor respectively. We also have

$$P_m = \omega \frac{0.3E}{Q_m},$$

$$P_r = \omega \frac{E}{Q_r},$$

$$P_s = 0,$$

where  $Q_m = 50$  and  $Q_r = 500$ . Plugging these in and simplifying, we get

$$\begin{aligned} Q &= \omega \frac{E}{0.3\omega E/Q_m + \omega E/Q_r} \\ &= \frac{1}{0.3/Q_m + 1/Q_r} \\ &= 125. \end{aligned}$$

(b) What is the threshold gain and quantum efficiency of the laser?

$$\begin{aligned} g_{th} &= \frac{\alpha}{\Gamma} \\ &= \frac{2\pi n}{\lambda_0 Q \Gamma} \\ &= 2154 \text{ cm}^{-1}, \end{aligned}$$

using  $n = 3$ ,  $\lambda_0 = 1\mu\text{m}$ , and  $\Gamma = 0.7$ , since 70% of the energy is in the semiconductor, and only the semiconductor experiences gain.

$$\eta = \frac{P_r}{P_{tot}} = \frac{Q}{Q_r} = 0.25$$

3. (a) Use the energy reference below (i.e,  $E_V = 0$  and  $E_C = E_g$ , the bandgap energy), find  $E_1$  and  $E_2$  as functions of the photon energy,  $\hbar\omega$ .

$$\begin{aligned} E_1 &= -\frac{\hbar^2 k^2}{2m_h^*} \\ E_2 &= E_g + \frac{\hbar^2 k^2}{2m_e^*} \end{aligned}$$

Subtract  $E_2$  and  $E_1$  and equate to  $\hbar\omega$ :

$$\begin{aligned} E_2 - E_1 &= \hbar\omega \\ &= E_g + \frac{\hbar^2 k^2}{2} \left( \frac{1}{m_e^*} + \frac{1}{m_h^*} \right) \\ &= E_g + \frac{\hbar^2 k^2}{2m_r^*} \end{aligned}$$

Rewriting  $E_1$  and  $E_2$  in terms of  $m_r^*$  and using the above equation, we get:

$$\begin{aligned} E_1 &= -\frac{\hbar^2 k^2}{2m_r^*} \frac{m_r^*}{m_h^*} \\ &= -\frac{m_r^*}{m_h^*} (\hbar\omega - E_g) \end{aligned}$$

$$\begin{aligned}
E_2 &= E_g + \frac{\hbar^2 k^2}{2m_r^*} \frac{m_r^*}{m_e^*} \\
&= E_g + (\hbar\omega - E_g) \frac{m_r^*}{m_e^*}
\end{aligned}$$

(b) Derive  $f_C(E_2(\hbar\omega))$  as a function of  $\hbar\omega$ .

$$\begin{aligned}
f_C(E_2) &= \frac{1}{1 + e^{(E_2 - F_C)/kT}} \\
&= \frac{1}{1 + e^{(E_g + (\hbar\omega - E_g) \frac{m_r^*}{m_e^*} - F_C)/kT}},
\end{aligned}$$

plugging in for  $E_2$  from above.

(c) Derive  $f_V(E_1(\hbar\omega))$  as a function of  $\hbar\omega$ .

$$\begin{aligned}
f_V(E_1) &= \frac{1}{1 + e^{(E_1 - F_V)/kT}} \\
&= \frac{1}{1 + e^{(-\frac{m_r^*}{m_h^*}(\hbar\omega - E_g) - F_V)/kT}}
\end{aligned}$$

(d) Calculate and plot the emission probability  $f_e(\hbar\omega) = f_C(\hbar\omega)(1 - f_V(\hbar\omega))$  for photon energies from 0.8 to 1.5 eV. Plot for two temperatures:  $T = 0$  and  $T = 300$  K.

Simply plug in  $f_C$  and  $f_V$  found above:

$$f_e(\hbar\omega) = \frac{1}{1 + e^{(E_g + (\hbar\omega - E_g) \frac{m_r^*}{m_e^*} - F_C)/kT}} \left( 1 - \frac{1}{1 + e^{(-\frac{m_r^*}{m_h^*}(\hbar\omega - E_g) - F_V)/kT}} \right)$$

To use this equation we must find  $F_C$  and  $F_V$ .

Since  $F_C - F_V > E_g$ ,  $F_C$  and/or  $F_V$  is degenerate. Since  $m_e^* < m_h^*$ , electrons are degenerate:  $F_C > E_g$ . Assume holes are also degenerate:  $F_V < 0$ . Then we have

$$\begin{aligned}
n &\propto N_C(F_C - E_g)^{3/2} \propto m_e^{*3/2}(F_C - E_g)^{3/2} \\
p &\propto N_V(-F_V)^{3/2} \propto m_h^{*3/2}(-F_V)^{3/2}
\end{aligned}$$

with the same proportionality factor for  $n$  and  $p$ . Since  $n = p$ , we have

$$\begin{aligned}
\frac{m_h^*}{m_e^*} &= \frac{F_C - E_g}{-F_V} \\
4 &= \frac{F_V + 1.2\text{eV} - 1\text{eV}}{-F_V} \\
-4F_V &= F_V + 0.2\text{eV} \\
F_V &= -0.04\text{eV} \\
F_C &= F_V + 1.2\text{eV} = 1.16\text{eV}.
\end{aligned}$$

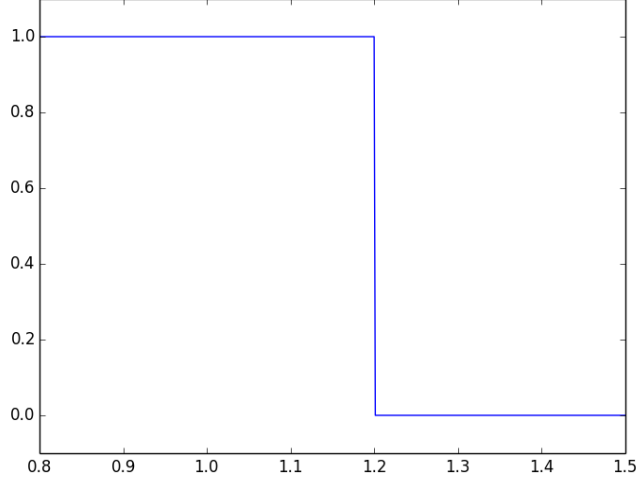


Figure 1:  $f_e$  vs  $\hbar\omega$  (eV) at  $T = 0K$ .

Note that  $F_V < 0$ , justifying our original assumption.

$f_e$  is plotted in Figs. 1 and 2.

- (e) Repeat part d) for the Fermi inversion factor  $f_g(\hbar\omega) = f_C(\hbar\omega) - f_V(\hbar\omega)$ .

$$f_g(\hbar\omega) = \frac{1}{1 + e^{\left(E_g + (\hbar\omega - E_g) \frac{m_r^*}{m_e^*} - F_C\right)/kT}} - \frac{1}{1 + e^{\left(-\frac{m_r^*}{m_h^*} (\hbar\omega - E_g) - F_V\right)/kT}}$$

$f_g$  is plotted in Figs. 3 and 4.

- (f) Plot the gain spectra for  $T = 0$  and  $T = 300$  K for the condition given in d).

$$g(\hbar\omega) = C_0 |\hat{\mathbf{e}} \cdot \mathbf{p}|^2 \rho_r(\hbar\omega - E_g) f_g(\hbar\omega)$$

$$C_0 = \frac{\pi e^2}{nc\epsilon_0 m_0^2 \omega}$$

$$|\hat{\mathbf{e}} \cdot \mathbf{p}|^2 = \frac{m_0}{6} E_p, \quad E_p = 25.7 \text{ eV}$$

$$\rho_r(\hbar\omega - E_g) = H(\hbar\omega - E_g) \frac{1}{2\pi^2} \left( \frac{2m_r^*}{\hbar^2} \right)^{3/2} \sqrt{\hbar\omega - E_g}$$

$$f_g(\hbar\omega) = f_C(\hbar\omega) - f_V(\hbar\omega)$$

Assuming  $n = 3.5$ ,  $g(\hbar\omega)$  is plotted in Figs. 5 and 6.

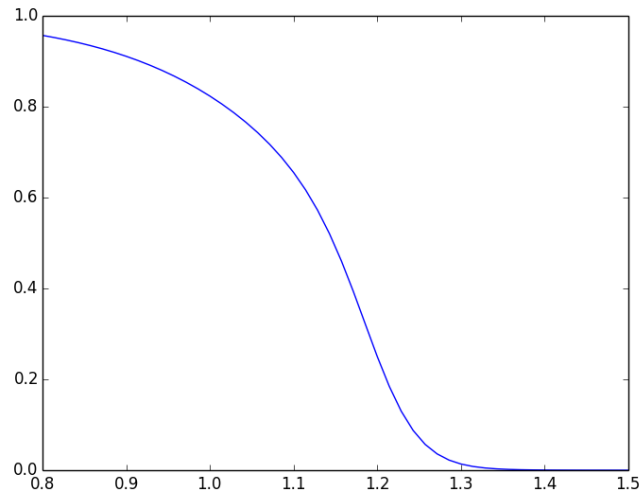


Figure 2:  $f_e$  vs  $\hbar\omega$  (eV) at  $T = 300K$ .

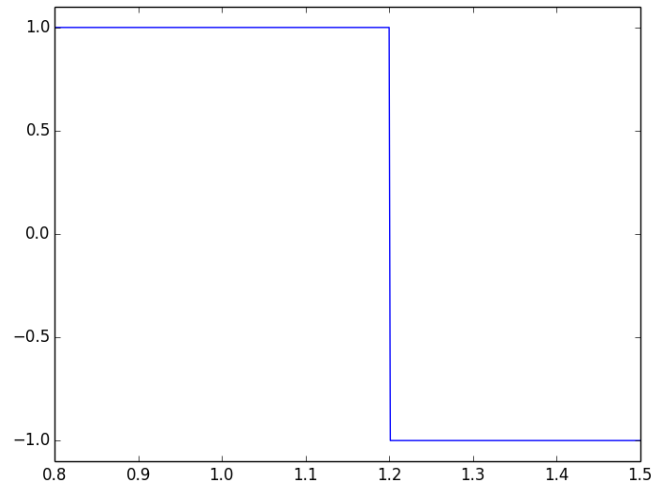


Figure 3:  $f_g$  vs  $\hbar\omega$  (eV) at  $T = 0K$ .

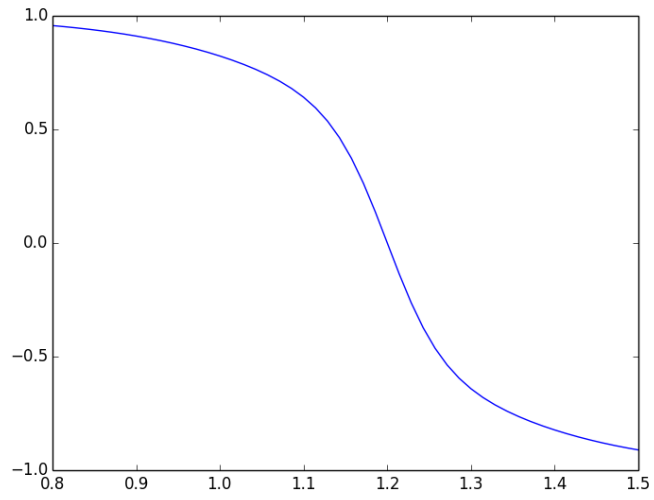


Figure 4:  $f_g$  vs  $\hbar\omega$  (eV) at  $T = 300K$ .

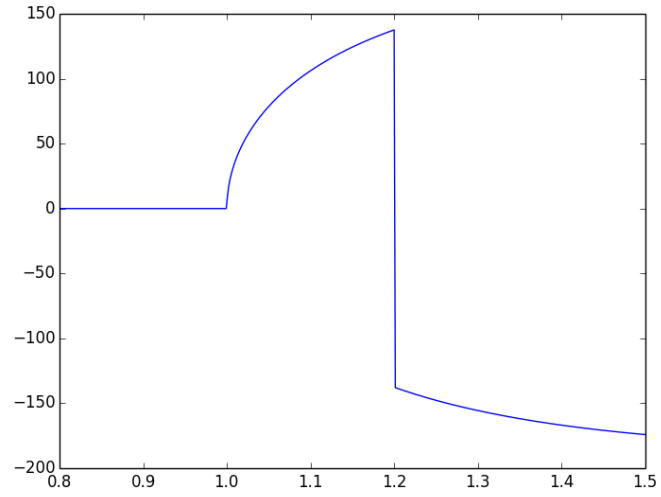


Figure 5:  $g$  ( $\text{cm}^{-1}$ ) vs  $\hbar\omega$  (eV) at  $T = 0K$ .

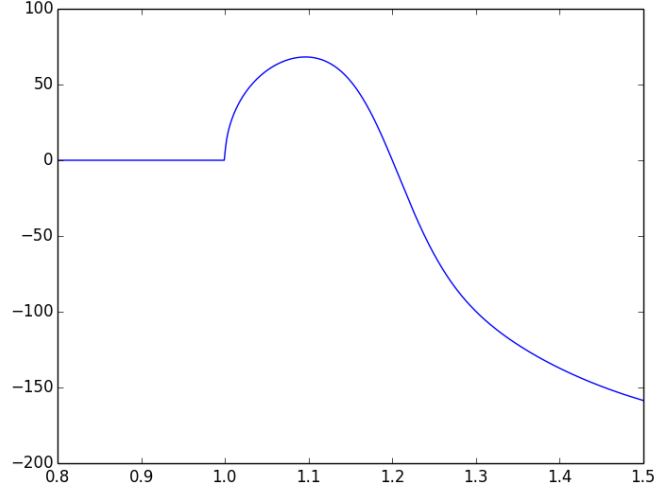


Figure 6:  $g$  (cm $^{-1}$ ) vs  $\hbar\omega$  (eV) at  $T = 300K$ .

- (g) Plot the spontaneous emission spectra for  $T = 0$  and  $T = 300$  K for the condition given in d).

$$r_{\text{spn}}(\hbar\omega) = \frac{8\pi n^2(\hbar\omega)^2}{h^3 c^2} \frac{f_e(\hbar\omega)}{f_g(\hbar\omega)} g(\hbar\omega)$$

$r_{\text{spn}}$  is plotted in Figs. 7 and 8.

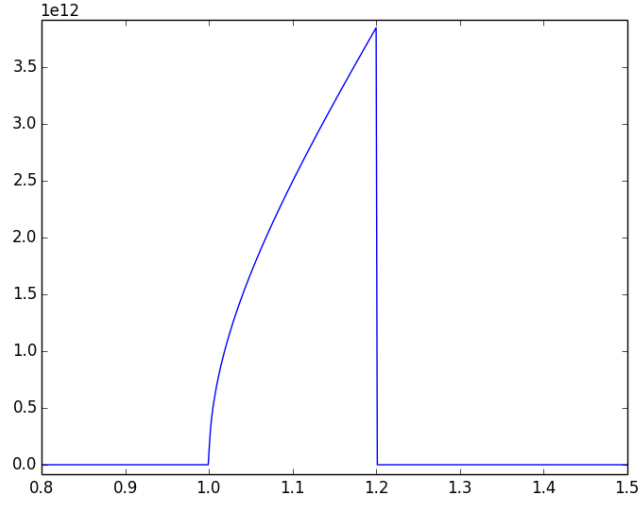


Figure 7:  $r_{\text{spon}}$  ( $\text{eV}^{-1}\text{um}^{-3}\text{s}^{-1}$ ) vs  $\hbar\omega$  (eV) at  $T = 0\text{K}$ .

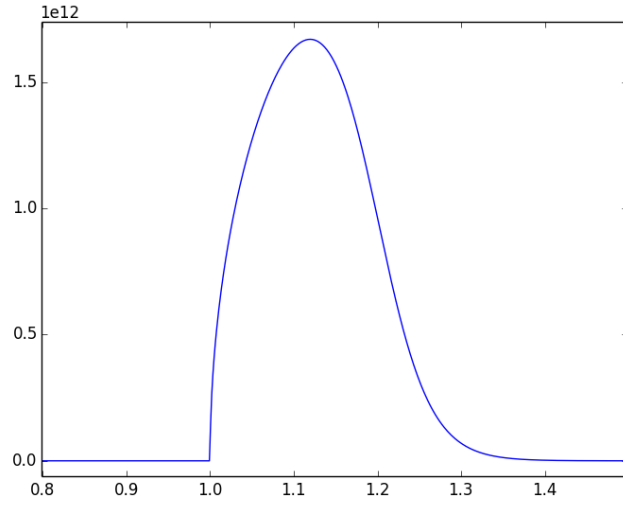


Figure 8:  $r_{\text{spon}}$  ( $\text{eV}^{-1}\text{um}^{-3}\text{s}^{-1}$ ) vs  $\hbar\omega$  (eV) at  $T = 300\text{K}$ .