

UC Berkeley

Semiconductor Doping

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 48

48

UC Berkeley

Doping of Semiconductors

- Semiconductors are not intrinsically conductive
- To make them conductive, replace silicon atoms in the lattice with dopant atoms that have valence bands with fewer or more e⁻s than the 4 of Si
- If more e⁻s, then the dopant is a donor: P, As
 - The extra e⁻ is effectively released from the bonded atoms to join a cloud of free e⁻s, free to move like e⁻s in a metal

- The larger the # of donor atoms, the larger the # of free e⁻s → the higher the conductivity

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 49

49

UC Berkeley

Doping of Semiconductors (cont.)

- Conductivity Equation:**

$$\sigma = q\mu_n n + q\mu_p p$$

Labels for the equation:

 - σ : conductivity
 - q : charge magnitude on an electron
 - μ_n : electron mobility
 - n : electron density
 - μ_p : hole mobility
 - p : hole density
- If fewer e⁻s, then the dopant is an acceptor: B
 - Diagram showing the doping of silicon with boron (B). On the left, a silicon lattice is shown with four silicon atoms, each having four valence electrons. On the right, a boron atom (B) is introduced, labeled 'Dope'. The boron atom has three valence electrons. One of the silicon atoms is shown with a missing electron, labeled 'hole'. A red circle highlights the boron atom and its surrounding silicon atoms, with a note 'creates empty hole'.

- Lack of an e⁻ = hole = h⁺
- When e⁻s move into h⁺s, the h⁺s effectively move in the opposite direction → a h⁺ is a mobile (+) charge carrier

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 50

50

UC Berkeley

General Comments on Predeposition

- Higher doses only: $Q = 10^{13} - 10^{16} \text{ cm}^{-2}$ (I/I is $10^{11} - 10^{16}$)
- Dose not well controlled: $\pm 20\%$ (I/I can get $\pm 1\%$)
- Uniformity is not good
 - $\pm 10\%$ w/ gas source
 - $\pm 2\%$ w/ solid source
- Max. conc. possible limited by solid solubility
 - Limited to $\sim 10^{20} \text{ cm}^{-3}$
 - No limit for I/I → you force it in here!
- For these reasons, I/I is usually the preferred method for introduction of dopants in transistor devices
- But I/I is not necessarily the best choice for MEMS
 - I/I cannot dope the underside of a suspended beam
 - I/I yields one-sided doping → introduces unbalanced stress → warping of structures
 - I/I can do physical damage → problem if annealing is not permitted
- Thus, predeposition is often preferred when doping MEMS

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 65

65

Diffusion Modeling

UC Berkeley

Modeling $N(x)$

\Rightarrow Dopants from points of high conc. move to points of low conc. w/ flux J
 \Rightarrow Question: What's $N(x,t)$?
 ? fn of time

Fick's Law of Diffusion - (1st law)
 $J(x,t) = -D \frac{\partial N(x,t)}{\partial x}$ (1)
 Flux [$\#/cm^2 \cdot s$] \leftarrow Diffusion Coefficient

Continuity Equation for Particle Flux -
 General Form: $\frac{\partial N(x,t)}{\partial t} = -\nabla \cdot \vec{J}$
 rate of increase of conc. w/ time \leftarrow negative of the divergence of particle flux

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 66

66

Diffusion Modeling (cont.)

UC Berkeley

\Rightarrow We're interested for now in the one-dimensional form:

$$\frac{\partial N(x,t)}{\partial t} = -\frac{\partial J}{\partial x}$$
 (2)

$\left[\frac{\partial}{\partial x}(1) \text{ and substitute (2) in (1)} \right] \Rightarrow \frac{\partial N(x,t)}{\partial t} = D \frac{\partial^2 N(x,t)}{\partial x^2}$ [Fick's 2nd Law of Diffusion in 1-D]

Solutions: \rightarrow dependent upon boundary conditions
 \rightarrow use variable separation or Laplace Xform techniques

Case 1: Predeposition \rightarrow constant source diffusion: surface concentration stays the same during the diffusion

surface conc. stays constant $\rightarrow N_0$
 background conc. $\rightarrow N_B$
 t_1, t_2, t_3
 high T ($Dt_1 < Dt_2 < Dt_3$)
 complementary error function profile

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 67

67

Diffusion Modeling (Predeposition)

UC Berkeley

\Rightarrow if plotted on a linear scale, would look like this:

\Rightarrow **Boundary Condition:**
 (i) $N(0,t) = N_0$
 (ii) $N(\infty,t) = 0$

$$N(x,t) = N_0 \left[1 - \frac{1}{\sqrt{\pi}} \int_0^{\frac{x}{2\sqrt{Dt}}} e^{-y^2} dy \right]$$

$N(x,t) = N_0 \operatorname{erfc}\left(\frac{x}{2\sqrt{Dt}}\right)$ \Rightarrow again, complementary error function (read tables or graph)

Dose, $Q \cong$ total # of impurity atoms per unit area in the Si
 = area under the curve
 $Q = \int_0^{\infty} N(x,t) dx \Rightarrow Q(t) = N_0 \frac{2\sqrt{Dt}}{\sqrt{\pi}} \text{ cm}^{-2}$

$2\sqrt{Dt} \cong$ characteristic diffusion length

$N(x)$ \leftarrow linear scale
 area under this square is same as under the curve!
 $2\sqrt{Dt}$

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 68

68

Diffusion Modeling (Limited Source)

UC Berkeley

Case 2: Drive-in \rightarrow limited source diffusion, i.e., constant dose Q

$N_0(t_1)$
 $N_0(t_2)$
 $N_0(t_3)$
 N_B
 t_1, t_2, t_3
 x , distance from the surface

\Rightarrow **Boundary Condition:**
 (i) $N(\infty,t) = 0$
 (ii) $\frac{\partial N(x,t)}{\partial x} \Big|_{x=0} = 0$

Why? Constant Dose: $\int_0^{\infty} N(x,t) dx = Q \leftarrow \text{const.}$

This is equivalent to saying that there's no flux going out of the Si, i.e.,
 and that's what this says!

$J=0$

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 69

69

Diffusion Modeling (Limited Source)

UC Berkeley

Usually make delta fun. approx.: $N(x,0) = Q\delta(x)$
 \Rightarrow we can do this, because for sufficiently long diffusion times, no matter what the original shape of the dopant distribution, the diffused distribution will be the same

Get Gaussian Distribution: $N(x,t) = \frac{Q}{\sqrt{\pi Dt}} \exp\left[-\frac{x^2}{2Dt}\right]$ corresponds to a half Gaussian in this Equation

When the starting conc. profile is completely contained in the Si, then $Q = \frac{D_I}{2} = \text{half the implant dose}$

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 70

70

Two-Step Diffusion

UC Berkeley

- Two step diffusion procedure:
 - Step 1: predeposition (i.e., constant source diffusion)
 - Step 2: drive-in diffusion (i.e., limited source diffusion)
- For processes where there is both a predeposition and a drive-in diffusion, the final profile type (i.e., complementary error function or Gaussian) is determined by which has the much greater Dt product:
 - $(Dt)_{\text{predep}} \gg (Dt)_{\text{drive-in}} \Rightarrow$ impurity profile is complementary error function
 - $(Dt)_{\text{drive-in}} \gg (Dt)_{\text{predep}} \Rightarrow$ impurity profile is Gaussian (which is usually the case)

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 71

71

Successive Diffusions

UC Berkeley

- For actual processes, the junction/diffusion formation is only one of many high temperature steps, each of which contributes to the final junction profile
- Typical overall process:
 - Selective doping
 - Implant \rightarrow effective $(Dt)_1 = (\Delta R_p)^2/2$ (Gaussian)
 - Drive-in/activation $\rightarrow D_2 t_2$
 - Other high temperature steps
 - (eg., oxidation, reflow, deposition) $\rightarrow D_3 t_3, D_4 t_4, \dots$
 - Each has their own Dt product
 - Then, to find the final profile, use

$$(Dt)_{\text{tot}} = \sum_i D_i t_i$$
 in the Gaussian distribution expression.

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 72

72

The Diffusion Coefficient

UC Berkeley

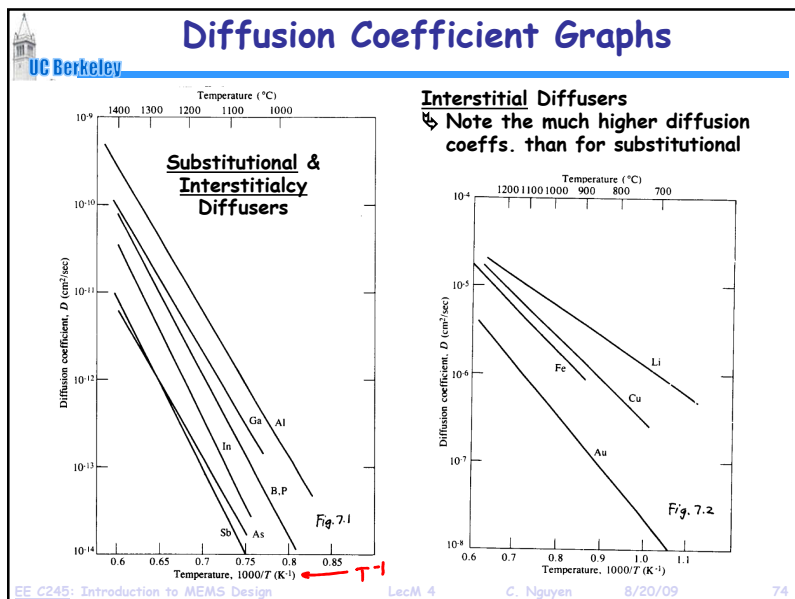
$$D = D_0 \exp\left(-\frac{E_A}{kT}\right) \quad (\text{as usual, an Arrhenius relationship})$$

Table 4.1 Typical Diffusion Coefficient Values for a Number of Impurities.

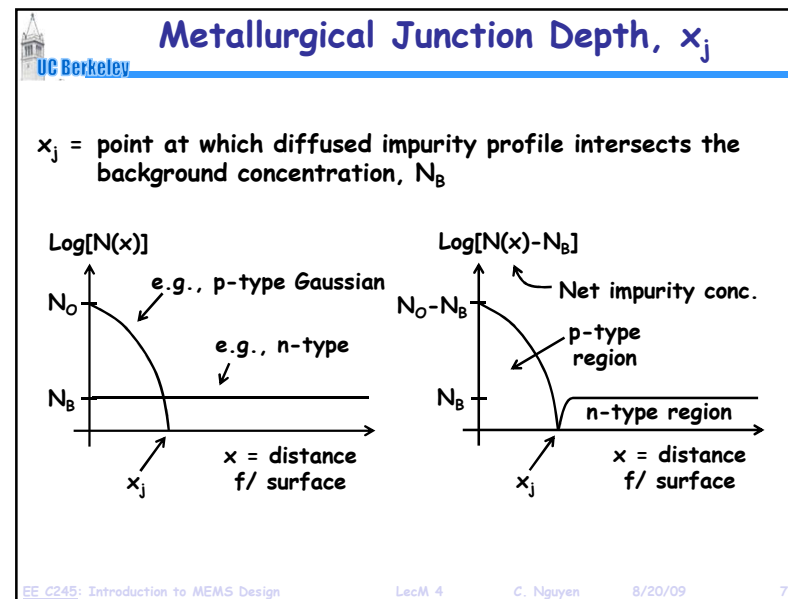
Element	$D_0(\text{cm}^2/\text{sec})$	$E_A(\text{eV})$
B	10.5	3.69
Al	8.00	3.47
Ga	3.60	3.51
In	16.5	3.90
P	10.5	3.69
As	0.32	3.56
Sb	5.60	3.95

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 73

73



74



75

Expressions for x_j

- Assuming a Gaussian dopant profile: (the most common case)

$$N(x_j, t) = N_0 \exp\left[-\left(\frac{x_j}{2\sqrt{Dt}}\right)^2\right] = N_B \rightarrow x_j = 2\sqrt{Dt \ln\left(\frac{N_0}{N_B}\right)}$$

- For a complementary error function profile:

$$N(x_j, t) = N_0 \operatorname{erfc}\left(\frac{x_j}{2\sqrt{Dt}}\right) = N_B \rightarrow x_j = 2\sqrt{Dt} \operatorname{erfc}^{-1}\left(\frac{N_B}{N_0}\right)$$

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 76

76

Sheet Resistance

- Sheet resistance provides a simple way to determine the resistance of a given conductive trace by merely counting the number of effective squares
- Definition:

$$R = \frac{\rho L}{A} = \left(\frac{\rho}{t}\right) \frac{L}{w} = R_s \left(\frac{L}{w}\right)$$

ohms per square
 ρ/t
sheet resistance
unit squares of material in the resistor

Uniformly doped material w/ resistivity $\rho = \frac{1}{\sigma}$
 $\sigma = \text{conductivity} = q(\mu_n n + \mu_p p)$

e.g.,
5 D's of material
: $R = R_s \times 5$

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 77

77

Squares From Non-Uniform Traces

UC Berkeley

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 78

78

Sheet Resistance of a Diffused Junction

UC Berkeley

• For diffused layers:

Majority carrier mobility

Net impurity concentration

Effective resistivity

Sheet resistance

$$R_s = \frac{\rho}{x_j} = \left[\int_0^{x_j} \sigma(x) dx \right]^{-1} = \left[\int_0^{x_j} q\mu N(x) dx \right]^{-1}$$

deplet. conc. decreases [extrinsic material]

• This expression neglects depletion of carriers near the junction, $x_j \rightarrow$ thus, this gives a slightly lower value of resistance than actual

• Above expression was evaluated by Irvin and is plotted in "Irvin's curves" on next few slides

↳ Illuminates the dependence of R_s on x_j , N_0 (the surface concentration), and N_B (the substrate background conc.)

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 79

79

Irvin's Curves (for n-type diffusion)

UC Berkeley

Example. p-type

Given:

- $N_B = 3 \times 10^{16} \text{ cm}^{-3}$
- $N_0 = 1.1 \times 10^{18} \text{ cm}^{-3}$ (n-type Gaussian)
- $x_j = 2.77 \text{ } \mu\text{m}$

Can determine these given known predep. and drive conditions

Determine the R_s .

Using Fig. 7.7:

$R_s x_j = 470 \text{ } \Omega \cdot \mu\text{m}$

$\therefore R_s = \frac{470}{2.77} = 170 \text{ } \Omega / \square$

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 80

80

Irvin's Curves (for p-type diffusion)

UC Berkeley

Example. n-type

Given:

- $N_B = 3 \times 10^{16} \text{ cm}^{-3}$
- $N_0 = 1.1 \times 10^{18} \text{ cm}^{-3}$ (p-type Gaussian)
- $x_j = 2.77 \text{ } \mu\text{m}$

Can determine these given known predep. and drive conditions

Determine the R_s .

Using Fig. 7.9:

$R_s x_j = 800 \text{ } \Omega \cdot \text{cm}$

$\therefore R_s = \frac{800}{2.77} = 289 \text{ } \Omega / \square$

EE C245: Introduction to MEMS Design LecM 4 C. Nguyen 8/20/09 81

81